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Differential Equations II for Engineering Students

Homework sheet 2

Exercise 1:

Determine the solutions to the following initial value problems for $t \in \mathbb{R}^+$, $x \in \mathbb{R}$.

a) $u_t + 3u_x = 0$ with $u(x, 0) = u_0(x) = xe^{-x}$.

b)
$$2u_t + x^2 u_x = \frac{1}{u}$$
 with $u(x,0) = 2\sqrt{e^{-4x^2}}$.

Does there exist a solution for all $t \in \mathbb{R}^+$, $x \in \mathbb{R}$?

If not, can the solution be continuously extended in the definition gaps (to be defined in the whole domain)?

Solution:

a) $u_t + 3u_x = 0$ with $u(x, 0) = xe^{-x}$.

On a fixed characteristic (t, x(t)) we have:

$$\dot{x}(t) = 3 \implies x(t) = c + 3t, \ x(0) = c = x_0 = x - 3t.$$

 $\dot{u}(t) = 0 \implies u$ is constant on the characteristic!

So

$$\Rightarrow \begin{cases} u(x,t) = u(x_0,0) = u(x-3t,0) = u_0(x-3t) \\ u(x,0) = xe^{-x} \end{cases}$$
$$\Rightarrow u(x,t) = (x-3t)e^{-(x-3t)}$$

b) For x=0 one gets an ordinary differential equation $2u_t=\frac{1}{u}$ with the solution $u(0,t)=\sqrt{t+C}$. Using initial value we get C=4. For $x\neq 0$ we compute as follows

$$\frac{dx}{dt} = \frac{x^2}{2} \qquad \frac{dx}{x^2} = \frac{dt}{2} \qquad -\frac{1}{x} = \frac{t}{2} - C_1 \qquad C_1 = \frac{t}{2} + \frac{1}{x}$$

$$\frac{du}{dt} = \frac{1}{2u} \qquad 2u \cdot du = dt \qquad u^2 = t + C_2 \qquad C_2 = u^2 - t$$

$$C_2 = f(C_1) \iff u^2 - t = f(\frac{t}{2} + \frac{1}{x})$$

From the initial data follows

$$(u(x,0))^2 - 0 = 4e^{-4x^2} = f(\frac{1}{x}).$$

So

$$f(y) = 4e^{-4(1/y)^2} \implies u^2 = t + 4\exp\left(-4(\frac{t}{2} + \frac{1}{x})^{-2}\right)$$

$$u(x,t) = \sqrt{t + 4\exp\left(-4\left(\frac{(2x)^2}{(tx+2)^2}\right)} = \sqrt{t + 4e^{\frac{-16x^2}{(tx+2)^2}}}$$
.

The soltuion is not defined for x(t)=-2/t. For every fixed $t\in\mathbb{R}^+$ it holds $x^2=4/t^2\neq 0$ and

$$\lim_{x \to -2/t} \frac{-16x^2}{(tx+2)^2} = -\infty \Longrightarrow \lim_{x \to -2/t} e^{\frac{-16x^2}{(tx+2)^2}} = 0$$

and hence

$$\lim_{x \to -2/t} u(x,t) = \sqrt{t}.$$

Exercise 2:

A simple traffic flow model:

We consider a one-dimensional flow of vehicles along an infinitely long, single-lane road. In a so-called macroscopic model, one does not consider individual vehicles, but the total flow of vehicles. For this purpose, we introduce the following quantities:

u(x,t) = (length-)density of the vehicles at the point x at the time t

= vehicles/unit length at point x at the time t

v(x,t) = speed at the point x at the time t,

 $q(x,t) = u(x,t) \cdot v(x,t) = \text{flow}$

= amount of vehicles passing the point x at the time t per unit time

a) Assume that there are no entrances or exits, no vehicles are disappearing, and no new vehicles are appearing. Let $N(t, a, \Delta a) := \text{number of vehicles on a space interval}$ $[a, a + \Delta a]$ at the time t.

Then on the one hand it holds that

$$N(t, a, \Delta a) = \int_{a}^{a+\Delta a} u(x, t) dx$$

and on the other hand it also holds

$$N(t, a, \Delta a) - N(t_0, a, \Delta a) = \int_{t_0}^t q(a, \tau) - q(a + \Delta a, \tau) d\tau.$$

Derive from this the so-called conservation equation for the mass (number of vehicles)

$$u_t + q_x = 0.$$

Hints on how to proceed:

• Derive both formulas for N with respect to t. Please note that for the derivation of parameter-dependent integrals with sufficiently smooth f holds the **Leibniz** rule:

$$\frac{d}{dx} \int_{a(x)}^{b(x)} f(x,t) dt = \int_{a(x)}^{b(x)} \frac{d}{dx} f(x,t) dt + b'(x) f(x,b(x)) - a'(x) f(x,a(x))$$

- Divide by Δa .
- Consider the limit $\Delta a \to 0$.
- b) Additionally assume that the velocity depends only on the density: v = v(u). Show that in this case the equation

$$\frac{\partial u}{\partial t} + \frac{dq}{du} \cdot \frac{\partial u}{\partial x} = 0$$

describes the conservation of mass.

c) We now assume in a first simple model that the speed increases in inverse proportion to the density and that the density is positive.

$$v(x,t) = c + \frac{k}{u(x,t)}$$

What is the continuity equation (=conservation equation for the mass)?

Solution:

a) On the one hand, it holds $N(t) = \int_a^{a+\Delta a} u(x,t) dx$ and on the other hand $N(t) - N(t_0) = \int_{t_0}^t q(a,\tau) - q(a+\Delta a,\tau) d\tau.$

Differentiating with respect to t gives

$$\frac{\partial}{\partial t} N(t) = \frac{\partial}{\partial t} \int_{a}^{a+\Delta a} u(x,t) dx = q(a,t) - q(a+\Delta a,t)$$

Letting Δa to zero, and with sufficient smoothness of the functions, we have

$$\lim_{\Delta a \to 0} \frac{1}{\Delta a} \int_{a}^{a+\Delta a} \frac{\partial}{\partial t} u(x,t) dx = \lim_{\Delta a \to 0} -\frac{q(a+\Delta a,t) - q(a,t)}{\Delta a}$$

$$\implies \frac{\partial}{\partial t} u(a,t) = -\frac{\partial}{\partial a} q(a,t).$$

Since these considerations hold at every point, we have the continuity equation $u_t + q_x = 0$.

b) Actually is straightforward, since in this case we have $q(x,t) = u(x,t) \cdot v(u(x,t))$. The flow q is therefore a function of u(x,t). The assertion then follows from the chain rule.

In more details:

With $q(x,t) = u(x,t) \cdot v(u(x,t))$ we have

$$\frac{dq}{du} \cdot \frac{\partial u}{\partial x} = \frac{d}{du} (u \cdot v(u)) \cdot u_x = (v(u) + u \cdot v_u) \cdot u_x$$

and on the other hand it holds

$$\frac{\partial}{\partial x} q(x,t) = \frac{\partial}{\partial x} (u(x,t) \cdot v(u(x,t))) = u_x \cdot v(u) + u \cdot v_u \cdot u_x.$$

c)
$$v(x,t) = c + \frac{k}{u(x,t)} \quad q(x,t) = c \cdot u(x,t) + k$$

From the continuity equation from part b) we have

$$\frac{\partial u}{\partial t} + \frac{dq}{du} \cdot \frac{\partial u}{\partial x} = \frac{\partial u}{\partial t} + c \cdot \frac{\partial u}{\partial x} = 0$$

The linear transport equation is thus obtained.

For c=3 the equation is solved in Exercise 1a).

Note: This is a very simple, linearized model. For example, it allows for any density and any speed. A somewhat more realistic problem would already produce shock and rarefaction waves (see later exercises).

Submission: 24.- 28.04.23