Representation categories of vertex algebras

ZMP Seminar

Vertex operator algebras and topological field theories from twisted QFTs in 3d and 4d

Overview

Basic definitions

2 The big picture of vertex algebra representation categories

3 How to find representations

Definition: Vertex algebra (VA) Data:

- vector space V
- vacuum vector $\Omega \in V$
- translation operator $T: V \rightarrow V$
- field map $Y: V \otimes V \rightarrow V((z))$

Axioms:

- vacuum axiom: $Y(\Omega, z) = \operatorname{id}_V$ and $Y(A, z)\Omega = A + zV[[z]], \ \forall A \in V$
- translation axiom: $T\Omega = 0$ and $[T, Y(A, z)] = \partial_z Y(A, z)$

• locality:
$$\forall A, B \in V, \exists n \in \mathbb{N}$$

 $(z - w)^n [Y(A, z), Y(B, w)] = 0$

Consequence/Proposition

For all $A, B, C \in V$ $Y(A, z)Y(B, w)C \in V((z))((w))$ $Y(B, w)Y(A, z)C \in V((w))((z))$ $Y(Y(A, z - w)B, w)C \in V((w))((z - w))$

Expansion of same element in $V[[z,w]][z^{-1},w^{-1},(z-w)^{-1}].$ The previous slide shows vertex algebras are close to associative commutative unital \mathbb{C} -algebras with a derivation. We also see that the fields Y(A, z) are essentially an action of *V* on itself.

Definition: Vertex algebra module

Let (V, Ω, T, Y) be a vertex algebra. A *V*-module is a pair (M, Y_M) : *M* a vector space and $Y_M : V \otimes M \to M((z))$ such that

- $Y_M(\Omega, z) = \mathrm{id}_M$
- $Y_M(TA, z) = \partial_z Y_M(A, z)$
- For all $A, B \in V$ and $C \in M$ the expansions $Y_M(A, z)Y_M(B, w)C \in M((z))((w))$ $Y_M(B, w)Y_M(A, z)C \in M((w))((z))$ $Y_M(Y(A, z - w)B, w)C \in M((w))((z - w))$ can be identified in $M[[z, w]][z^{-1}, w^{-1}, (z - w)^{-1}]$.

Many additional assumtions can be added to the above definition. E.g. bounded conformal weights, finite weight spaces, semi simplicity, etc.

Let (V, Ω, T, Y) be a vertex algebra and $(M_1, Y_{m_1}), (M_2, Y_{M_2}), \ldots$ be *V*-modules. In order use *V* for conformal field theory, one needs to consider chiral correlation functions:

 $\langle Y(A,z)\phi(m_1,x_1)\phi(m_2,x_2)\cdots\rangle$ $\langle \phi(Y(A,z-x_1)m_1,x_1)\phi(m_2,x_2)\cdots\rangle$ $\langle \phi(m_1,x_1)Y(A,z)\phi(m_2,x_2)\cdots\rangle$ $\langle \phi(m_1,x_1)\phi(Y(A,z-x_2)m_2,x_2)\cdots\rangle$ Let (V, Ω, T, Y) be a vertex algebra and $(M_1, Y_{m_1}), (M_2, Y_{M_2}), \dots$ be *V*-modules. In order use *V* for conformal field theory, one needs to consider chiral correlation functions:

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These are essentially V-multilinear maps.

Let (V, Ω, T, Y) be a vertex algebra and $(M_1, Y_{m_1}), (M_2, Y_{M_2}), (M_3, Y_{M_3})$ be *V*-modules. An intertwining operator of type $\binom{M_3}{M_1, M_2}$ is a map $\mathcal{Y} : M_1 \otimes M_2 \to M_3 x$ such that for all $m_i \in M_i$

- $\mathcal{Y}(m_1, z)m_2$ truncates below.
- $\mathcal{Y}(Tm_1, z) = \partial_z \mathcal{Y}(m_1, z).$
- The expansions

 $Y_{M_3}(A, z)\mathcal{Y}(m_1, x)m_2 \sim \mathcal{Y}(Y_{M_1}(A, z-x)m_1, x)m_2 \sim \mathcal{Y}(m_1, x)Y_{M_2}(A, z)m_2$ can be identified (via VA version of Jacobi identity).

Observations:

- The field map Y is an intertwining operator of type $\binom{V}{V,V}$.
- The action Y_M is an intertwining operator of type $\binom{M}{V,M}$.
- Intertwining operators are *V*-bilinear maps. All intertwining operators of a given type form a vector space. The field map *Y* and the action Y_M have a distinguished normalisation due to $Y_*(\omega, z) = id_*$.

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Tensor products pull multilinear algebra back to linear algebra!

Definition: Fusion product aka vertex algebra tensor product

Let (V, Ω, T, Y) be a vertex algebra and $(M_1, Y_{m_1}), (M_2, Y_{M_2})$ be *V*-modules. A fusion product is a triple $(M_1 \boxtimes M_2, Y_{M_1 \boxtimes M_2}, \mathcal{Y}_{M_1,M_2})$, where $(M_1 \boxtimes M_2, Y_{M_1 \boxtimes M_2})$ is a *V*-module and \mathcal{Y}_{M_1,M_2} is an intertwining operator of type $\binom{M_1 \boxtimes M_2}{M_1, M_2}$ such that the following universal property holds: For every *V*-module (X, Y_X) and intertwining operator \mathcal{Y}_X of type $\binom{X}{M_1, M_2}$



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Example: Heisenberg vertex algebra. Let F_{μ} , be the Fock space of weight $\mu \in \mathbb{C}$. Then $\dim {F_{\rho} \choose F_{\mu}, F_{\nu}} = \delta_{\rho, \mu+\nu}$ for all $\rho, \mu, \nu \in \mathbb{C}$. ${F_{\mu, F_{\nu}} \choose F_{\mu}, F_{\nu}}$ is spanned by

$$\begin{aligned} \mathcal{Y}_{F_{\mu},F_{\nu}}(p|\mu\rangle,z)q|\nu\rangle &= z^{\mu\nu}S_{\mu}\prod_{m\geq 1}\exp\left(\mu\frac{a_{-m}}{m}z^{m}\right)Y_{F_{\nu}}(p|0\rangle,z)\\ &\cdot\prod_{m\geq 1}\exp\left(-\mu\frac{a_{m}}{m}z^{-m}\right)q|\nu\rangle,\end{aligned}$$

where S_{μ} is the shift operator.

- For module homomorphimsm *f* : *X* → *Z*, *g* : *U* → *W*, *f* ⊠ *g* is uniquely characterised by
 (*f* ⊠ *g*)*Y*_{X⊠U} = *Y*_{Z⊠W} ∘ (*f* ⊗ *g*)
- *V* is the tensor identity and the unit isomorphisms are uniquely characterised by $\int_{M} (Q_{rac}(q, z)m) = V_{rac}(q, z)m \text{ and }$
 - $r_M(\mathcal{Y}_{M,V}(m,z)a) = e^{zT}Y_M(a,-z)m.$
- associativity isomorphisms (hardest part!) $A_{M_1,M_2,M_3}(\mathcal{Y}_{M_1,M_2\boxtimes M_3}(m_1,x_1)\mathcal{Y}_{M_2,M_3}(m_2,x_2)m_3) = \mathcal{Y}_{M_1\boxtimes M_2,M_3}(\mathcal{Y}_{M_1,M_2}(m_1,x_1)m_2,x_2)m_3$ All analytic details hidden.
- Braiding isomorphisms uniquely characterised by $c_{M_1,M_2}(\mathcal{Y}_{M_1,M_2}(m_1,x_1)m_2) = e^{zT}\mathcal{Y}_{M_2,M_1}(m_2,e^{i\pi}z)m_1$

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If the vertex algebra *V* is conformal (a vertex operator algebra) and the modules are chosen to be compatible with this conformal structure, then there is also a twist $\theta_M = e^{2\pi i L_0}|_M$, which satisfies the balancing equation

 $\theta_{M_1 \boxtimes M_2} = c_{M_1, M_2} \circ c_{M_2, M_1} \circ (\theta_{M_1} \boxtimes \theta_{M_2})$

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Theorem [Huang '04]: The Verlinde Conjecture

Let (V, Ω, ω, Y) be a vertex operator algebra and Adm V be the category of admissible *V*-modules. If

$$1 \dim V_0 = 1, \dim V_{-n} = 0, \dim V_n < \infty, \ n \in \mathbb{N},$$

2 V is simple as a module over itself,

3 $V \cong V'$, self-dual,

- $4 \dim V/c_2(V) < \infty,$
- **5** Adm(V) is semisimple,

then $\operatorname{Adm} V$ is a modular tensor category. Further the action of the modular group on the category (which determines Verlinde's formula) is equal (after a renormalisation) to the action of the modular group on module characters.

Summary of what we've discussed so far

- Vertex algebras are almost commutative unital algebras with derivations.
- The conformal vector is a choice/structure: there can be 0, 1 or many.
- Vertex algebras admit modules. "Good choices" of module categories admit a tensor (aka fusion) product.
- With the exception of associators, the tensor structure morphisms follow from natural constructions and are easy to obtain.

• Nothing presented so far helps with actually finding modules.

- For commutative algebras, the regular module can be used to present any finitely generated module.
- This fails for vertex (operator) algebras. The vertex algebra as a module over itself is often assumed/required to be simple. In general almost all modules cannot be presented from sums of the vertex algebra.
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Definition/Proposition: Zhu's associative algebra, [Zhu '96]

Let (V, Ω, ω, Y) be a vertex operator algebra and consider the two binary operations

$$a \circ b = \operatorname{Res} Y(a, z)b \frac{(1+z)^{h_a}}{z^2}, \quad a * b = \operatorname{Res} Y(a, z)b \frac{(1+z)^{h_a}}{z}.$$

Let $O(V) = \operatorname{span}\{a \circ b | \forall a, b \in V\}.$ Then the following hold.

- **1** A(V) = V/O(V) is a unital associative algebra under the binary operation *.
- **2** The class of the vacuum vector $[\Omega] = \Omega + O(V)$ is the identity element.
- **3** The class of the Virasoro vector $[\omega] = \omega + O(V)$ lies in the centre.
- 4 Let *M* be a *V*-module with ground state space \overline{M} . On \overline{M} $(a * b)_0 = a_0 b_0$ for all $a, b \in V$.
- If M is a (left) module over A(V) then it can be induced to a V-module with M as the space of ground states.

Virasoro algebra example:

- For the universal virasoro vertex operator algebra (V_c, Ω, ω, Y), of central charge c ∈ C we have C[X] ≅ A(V_c), where the isormophism is given by X ↦ [ω].
- At minimal model central charges

 $c_{p,q} = 1 - 6 \frac{(p-q)^2}{pq}, p,q \ge 2, \operatorname{gcd}(p,q) = 1,$ there is a singular vector $\chi \in V_{c_{p,q}}$ at degree (p-1)(q-1). Under the above isomorphism $[\chi] \in A(V_{c_{p,q}})$ corresponds to some f(X) and $A(V_{c_{p,q}}/\langle \chi \rangle) \cong C[X]/\langle f(X) \rangle.$

• For the Yang-Lee minimal model p = 2, q = 5 $\chi = (L_{-2}^2 - \frac{3}{5}L_{-4})\Omega$ $Y(\chi; z) = :T(z)^2 : -\frac{3}{10}\partial^2 T(z)$ Affine example:

- Let g be a complex finite dimensional simple Lie algebra and let V_k(g) be the universal affine vertex operator algebra (conformal vector given by the Sugawara construction). Then U(g) ≅ A(V_k(g)), where isomorphism is given by x ↦ [x₋₁Ω].
- If $k \in \mathbb{Z}_{\geq 0}$, $(e_{-1}^{\theta})^{k+1} \Omega$ is singular and generates the maximal ideal. $A(V_k(\mathfrak{g})/\langle (e_{-1}^{\theta})^{k+1} \Omega \rangle) \cong U(\mathfrak{g})/\langle (e^{\theta})^{k+1} \rangle$ is finite dimensional and semi simple.
- \mathfrak{sl}_2 at $k = -\frac{4}{3}$. The singular vector $\chi = (h_{-3} + 3e_{-2}f_{-1} - 3e_{-1}f_{-2} + \frac{9}{2}h_{-1}e_{-1}f_{-1} + \frac{9}{8}h_{-1}^3 - \frac{9}{4}h_{-2}h_{-1})\Omega$ generates the unique non-trivial ideal.

Literature

- Vertex algebras as rings: Frenkel Ben-Zvi, Chapters 1-5, http://dx.doi.org/10.1090/surv/088
- Fusion/tensor product theory: Huang Lepowsky Zhang 8 Part Series: arXiv:1012.4193, arXiv:1012.4196, arXiv:1012.4197, arXiv:1012.4198, arXiv:1012.4199, arXiv:1012.4202, arXiv:1110.1929, arXiv:1110.1931
- Zhu algebra theory: Zhu '96 JAMS https://www.jstor.org/stable/2152847
- Zhu algebra specialised to affine VOAs DOI: 10.1215/S0012-7094-92-06604-X