# Representation categories of vertex algebras 

ZMP Seminar

Vertex operator algebras and topological field theories from twisted QFTs in 3d and 4d

## Overview

(1) Basic definitions
(2) The big picture of vertex algebra representation categories
(3) How to find representations

## Definition: Vertex algebra (VA)

## Data:

- vector space $V$
- vacuum vector $\Omega \in V$
- translation operator $T: V \rightarrow V$
- field $\operatorname{map} Y: V \otimes V \rightarrow V((z))$


## Axioms:

- vacuum axiom:

$$
\begin{aligned}
& Y(\Omega, z)=\operatorname{id}_{V} \text { and } \\
& Y(A, z) \Omega=A+z V[[z]], \forall A \in V
\end{aligned}
$$

- translation axiom:
$T \Omega=0$ and
$[T, Y(A, z)]=\partial_{z} Y(A, z)$
- locality: $\forall A, B \in V, \exists n \in \mathbb{N}$ $(z-w)^{n}[Y(A, z), Y(B, w)]=0$


## Consequence/Proposition

For all $A, B, C \in V$
$Y(A, z) Y(B, w) C \in V((z))((w))$
$Y(B, w) Y(A, z) C \in V((w))((z))$
$Y(Y(A, z-w) B, w) C \in V((w))((z-w))$

Expansion of same element in

$$
V[[z, w]]\left[z^{-1}, w^{-1},(z-w)^{-1}\right]
$$

The previous slide shows vertex algebras are close to associative commutative unital $\mathbb{C}$-algebras with a derivation. We also see that the fields $Y(A, z)$ are essentially an action of $V$ on itself.

## Definition: Vertex algebra module

Let $(V, \Omega, T, Y)$ be a vertex algebra. A $V$-module is a pair $\left(M, Y_{M}\right): M$ a vector space and $Y_{M}: V \otimes M \rightarrow M((z))$ such that

- $Y_{M}(\Omega, z)=\mathrm{id}_{M}$
- $Y_{M}(T A, z)=\partial_{z} Y_{M}(A, z)$
- For all $A, B \in V$ and $C \in M$ the expansions

$$
\begin{aligned}
& Y_{M}(A, z) Y_{M}(B, w) C \in M((z))((w)) \\
& Y_{M}(B, w) Y_{M}(A, z) C \in M((w))((z)) \\
& Y_{M}(Y(A, z-w) B, w) C \in M((w))((z-w))
\end{aligned}
$$

$$
\text { can be identified in } M[[z, w]]\left[z^{-1}, w^{-1},(z-w)^{-1}\right] .
$$

Many additional assumtions can be added to the above definition. E.g. bounded conformal weights, finite weight spaces, semi simplicity, etc.

Let $(V, \Omega, T, Y)$ be a vertex algebra and $\left(M_{1}, Y_{m_{1}}\right),\left(M_{2}, Y_{M_{2}}\right), \ldots$ be $V$-modules. In order use $V$ for conformal field theory, one needs to consider chiral correlation functions:
$\left\langle Y(A, z) \phi\left(m_{1}, x_{1}\right) \phi\left(m_{2}, x_{2}\right) \cdots\right\rangle$

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```
\langleY(A,z)\phi(m},\mp@subsup{m}{1}{},\mp@subsup{x}{1}{})\phi(\mp@subsup{m}{2}{},\mp@subsup{x}{2}{})\cdots
\langle\phi(Y(A,z-\mp@subsup{x}{1}{})\mp@subsup{m}{1}{},\mp@subsup{x}{1}{})\phi(\mp@subsup{m}{2}{},\mp@subsup{x}{2}{})\cdots\rangle
\langle\phi(m},\mp@subsup{m}{1}{})Y(A,z)\phi(m2,\mp@subsup{x}{2}{})\cdots
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& \left\langle\phi\left(m_{1}, x_{1}\right) Y(A, z) \phi\left(m_{2}, x_{2}\right) \cdots\right\rangle \\
& \left\langle\phi\left(m_{1}, x_{1}\right) \phi\left(Y\left(A, z-x_{2}\right) m_{2}, x_{2}\right) \cdots\right\rangle
\end{aligned}
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$\left\langle\phi\left(m_{1}, x_{1}\right) \phi\left(Y\left(A, z-x_{2}\right) m_{2}, x_{2}\right) \cdots\right\rangle$
These are essentially $V$-multilinear maps.

## Definition: Intertwining operator

Let $(V, \Omega, T, Y)$ be a vertex algebra and $\left(M_{1}, Y_{m_{1}}\right),\left(M_{2}, Y_{M_{2}}\right),\left(M_{3}, Y_{M_{3}}\right)$ be $V$-modules. An intertwining operator of type $\binom{M_{1}}{M_{1}, M_{2}}$ is a map $\mathcal{Y}: M_{1} \otimes M_{2} \rightarrow M_{3} x$ such that for all $m_{i} \in M_{i}$

- $\mathcal{Y}\left(m_{1}, z\right) m_{2}$ truncates below.
- $\mathcal{Y}\left(m_{1}, z\right)=\partial_{z} \mathcal{Y}\left(m_{1}, z\right)$.
- The expansions

$$
Y_{M_{3}}(A, z) \mathcal{Y}\left(m_{1}, x\right) m_{2} \sim \mathcal{Y}\left(Y_{M_{1}}(A, z-x) m_{1}, x\right) m_{2} \sim \mathcal{Y}\left(m_{1}, x\right) Y_{M_{2}}(A, z) m_{2}
$$ can be identified (via VA version of Jacobi identity).

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- The expansions $Y_{M_{3}}(A, z) \mathcal{Y}\left(m_{1}, x\right) m_{2} \sim \mathcal{Y}\left(Y_{M_{1}}(A, z-x) m_{1}, x\right) m_{2} \sim \mathcal{Y}\left(m_{1}, x\right) Y_{M_{2}}(A, z) m_{2}$ can be identified (via VA version of Jacobi identity).


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- The expansions $Y_{M_{3}}(A, z) \mathcal{Y}\left(m_{1}, x\right) m_{2} \sim \mathcal{Y}\left(Y_{M_{1}}(A, z-x) m_{1}, x\right) m_{2} \sim \mathcal{Y}\left(m_{1}, x\right) Y_{M_{2}}(A, z) m_{2}$ can be identified (via VA version of Jacobi identity).


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$\mathcal{Y}: M_{1} \otimes M_{2} \rightarrow M_{3} x$ such that for all $m_{i} \in M_{i}$

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- The expansions $Y_{M_{3}}(A, z) \mathcal{Y}\left(m_{1}, x\right) m_{2} \sim \mathcal{Y}\left(Y_{M_{1}}(A, z-x) m_{1}, x\right) m_{2} \sim \mathcal{Y}\left(m_{1}, x\right) Y_{M_{2}}(A, z) m_{2}$ can be identified (via VA version of Jacobi identity).


## Observations:

- The field map $Y$ is an intertwining operator of type $\binom{V}{V, V}$.
- The action $Y_{M}$ is an intertwining operator of type $\left(\begin{array}{c} \\ V, M\end{array}\right)$.
- Intertwining operators are $V$-bilinear maps. All intertwining operators of a given type form a vector space. The field map $Y$ and the action $Y_{M}$ have a distinguished normalisation due to $Y_{*}(\omega, z)=\mathrm{id}_{*}$.

Tensor products pull multilinear algebra back to linear algebra!
Definition: Fusion product aka vertex algebra tensor product Let $(V, \Omega, T, Y)$ be a vertex algebra and $\left(M_{1}, Y_{m_{1}}\right),\left(M_{2}, Y_{M_{2}}\right)$ be $V$-modules. A fusion product is a triple ( $M_{1} \boxtimes M_{2}, Y_{M_{1} \boxtimes M_{2}}, \mathcal{Y}_{M_{1}, M_{2}}$ ), where ( $M_{1} \boxtimes M_{2}, Y_{M_{1} \boxtimes M_{2}}$ ) is a $V$-module and $\mathcal{Y}_{M_{1}, M_{2}}$ is an intertwining operator of type $\binom{M_{1} \mathbb{\otimes} M_{2}}{M_{1}, M_{2}}$ such that the following universal property holds: For every $V$-module ( $X, Y_{X}$ ) and intertwining operator $\mathcal{Y}_{X}$ of type $\binom{X}{M_{1}, M_{2}}$

$$
M_{1} \otimes M_{2} \underbrace{\substack{\mathcal{Y}_{1}, M_{2}}}_{\substack{\mathcal{Y}_{X}}} M_{1} \boxtimes M_{2}\{z\}
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$$
M_{1} \otimes M_{2} \underbrace{\substack{y_{M_{1}, M_{2}}}}_{\substack{\nu_{x}}} M_{1} \boxtimes M_{2}\{z\}
$$

In contrast to linear algebra (or ring theory) constructing $M_{1} \boxtimes M_{2}$ and decomposing into a direct sum of indecomposable modules is extremely hard.

Example: Heisenberg vertex algebra. Let $F_{\mu}$, be the Fock space of weight $\mu \in \mathbb{C}$. Then $\operatorname{dim}\binom{F_{\rho}}{F_{\mu}, F_{\nu}}=\delta_{\rho, \mu+\nu}$ for all $\rho, \mu, \nu \in \mathbb{C}$. $\binom{F_{\mu}+\nu}{F_{\mu}, F_{\nu}}$ is spanned by

$$
\mathcal{Y}_{F_{\mu}, F_{\nu}}(p|\mu\rangle, z) q|\nu\rangle=z^{\mu \nu} S_{\mu} \prod_{m \geq 1} \exp \left(\mu \frac{a_{-m}}{m} z^{m}\right) Y_{F_{\nu}}(p|0\rangle, z)
$$

$$
\prod_{m \geq 1} \exp \left(-\mu \frac{a_{m}}{m} z^{-m}\right) q|\nu\rangle
$$

where $S_{\mu}$ is the shift operator.

Well chosen categories of modules are tensor categories with respect to $\boxtimes$ with the following structures.

- For module homomorphimsm $f: X \rightarrow Z, g: U \rightarrow W, f \boxtimes g$ is uniquely characterised by $(f \boxtimes g) \mathcal{Y}_{X \boxtimes U}=\mathcal{Y}_{Z \boxtimes W} \circ(f \otimes g)$

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- $V$ is the tensor identity and the unit isomorphisms are uniquely characterised by

$$
\begin{aligned}
& \ell_{M}\left(\mathcal{Y}_{V, M}(a, z) m\right)=Y_{M}(a, z) m \text { and } \\
& r_{M}\left(\mathcal{Y}_{M, V}(m, z) a\right)=e^{z T} Y_{M}(a,-z) m .
\end{aligned}
$$

Well chosen categories of modules are tensor categories with respect to $\boxtimes$ with the following structures.

- associativity isomorphisms (hardest part!)
$A_{M_{1}, M_{2}, M_{3}}\left(\mathcal{Y}_{M_{1}, M_{2} \boxtimes M_{3}}\left(m_{1}, x_{1}\right) \mathcal{Y}_{M_{2}, M_{3}}\left(m_{2}, x_{2}\right) m_{3}\right)=$ $\mathcal{Y}_{M_{1} \boxtimes M_{2}, M_{3}}\left(\mathcal{Y}_{M_{1}, M_{2}}\left(m_{1}, x_{1}\right) m_{2}, x_{2}\right) m_{3}$ All analytic details hidden.

Well chosen categories of modules are tensor categories with respect to $\boxtimes$ with the following structures.

- Braiding isomorphisms uniquely characterised by $c_{M_{1}, M_{2}}\left(\mathcal{Y}_{M_{1}, M_{2}}\left(m_{1}, x_{1}\right) m_{2}\right)=e^{z T} \mathcal{Y}_{M_{2}, M_{1}}\left(m_{2}, e^{i \pi} z\right) m_{1}$

Well chosen categories of modules are tensor categories with respect to $\boxtimes$ with the following structures.

- For module homomorphimsm $f: X \rightarrow Z, g: U \rightarrow W, f \boxtimes g$ is uniquely characterised by $(f \boxtimes g) \mathcal{Y}_{X \boxtimes U}=\mathcal{Y}_{Z \boxtimes W} \circ(f \otimes g)$
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\end{aligned}
$$

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$\mathcal{Y}_{M_{1} \boxtimes M_{2}, M_{3}}\left(\mathcal{Y}_{M_{1}, M_{2}}\left(m_{1}, x_{1}\right) m_{2}, x_{2}\right) m_{3}$
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- Braiding isomorphisms uniquely characterised by $c_{M_{1}, M_{2}}\left(\mathcal{Y}_{M_{1}, M_{2}}\left(m_{1}, x_{1}\right) m_{2}\right)=e^{z T} \mathcal{Y}_{M_{2}, M_{1}}\left(m_{2}, e^{i \pi} z\right) m_{1}$

If the vertex algebra $V$ is conformal (a vertex operator algebra) and the modules are chosen to be compatible with this conformal structure, then there is also a twist $\theta_{M}=\left.e^{2 \pi i L_{0}}\right|_{M}$, which satisfies the balancing equation

$$
\theta_{M_{1} \boxtimes M_{2}}=c_{M_{1}, M_{2}} \circ c_{M_{2}, M_{1}} \circ\left(\theta_{M_{1}} \boxtimes \theta_{M_{2}}\right)
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$$

Tensor categories of vertex operator algebra modules depend only very weakly on the conformal structure. Only the twist and the dual (not discussed) depend on the conformal structure.

## Theorem [Huang '04]: The Verlinde Conjecture

Let $(V, \Omega, \omega, Y)$ be a vertex operator algebra and Adm $V$ be the category of admissible $V$-modules. If
(1) $\operatorname{dim} V_{0}=1, \operatorname{dim} V_{-n}=0, \operatorname{dim} V_{n}<\infty, n \in \mathbb{N}$,
(2) $V$ is simple as a module over itself,
(3) $V \cong V^{\prime}$, self-dual,
(4) $\operatorname{dim} V / c_{2}(V)<\infty$,
(5) $\operatorname{Adm}(V)$ is semisimple,
then $\operatorname{Adm} V$ is a modular tensor category. Further the action of the modular group on the category (which determines Verlinde's formula) is equal (after a renormalisation) to the action of the modular group on module characters.

## Summary of what we've discussed so far

- Vertex algebras are almost commutative unital algebras with derivations.
- The conformal vector is a choice/structure: there can be 0,1 or many.
- Vertex algebras admit modules. "Good choices" of module categories admit a tensor (aka fusion) product.
- With the exception of associators, the tensor structure morphisms follow from natural constructions and are easy to obtain.


## Practical matters

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- Nothing presented so far helps with actually finding modules.
- For commutative algebras, the regular module can be used to present any finitely generated module.
- This fails for vertex (operator) algebras. The vertex algebra as a module over itself is often assumed/required to be simple. In general almost all modules cannot be presented from sums of the vertex algebra.
- For vertex operator algebras (so with conformal structure) we can use Zhu's algebra.


## Definition/Proposition: Zhu's associative algebra, [Zhu '96]

 Let $(V, \Omega, \omega, Y)$ be a vertex operator algebra and consider the two binary operations$a \circ b=\operatorname{Res} Y(a, z) b \frac{(1+z)^{h_{a}}}{z^{2}}, \quad a * b=\operatorname{Res} Y(a, z) b \frac{(1+z)^{h_{a}}}{z}$. Let $O(V)=\operatorname{span}\{a \circ b \mid \forall a, b \in V\}$. Then the following hold.
(1) $A(V)=V / O(V)$ is a unital associative algebra under the binary operation $*$.
(2) The class of the vacuum vector $[\Omega]=\Omega+O(V)$ is the identity element.
(3) The class of the Virasoro vector $[\omega]=\omega+O(V)$ lies in the centre.
4. Let $M$ be a $V$-module with ground state space $\bar{M}$. On $\bar{M}$ $(a * b)_{0}=a_{0} b_{0}$ for all $a, b \in V$.
5. If $\bar{M}$ is a (left) module over $A(V)$ then it can be induced to a $V$-module with $\bar{M}$ as the space of ground states.

Virasoro algebra example:

- For the universal virasoro vertex operator algebra $\left(V_{c}, \Omega, \omega, Y\right)$, of central charge $c \in \mathbb{C}$ we have $\mathbb{C}[X] \cong A\left(V_{c}\right)$, where the isormophism is given by $X \mapsto[\omega]$.
- At minimal model central charges
$c_{p, q}=1-6 \frac{(p-q)^{2}}{p q}, p, q \geq 2, \operatorname{gcd}(p, q)=1$,
there is a singular vector $\chi \in V_{c_{p, q}}$ at degree $(p-1)(q-1)$.
Under the above isomorphism $[\chi] \in A\left(V_{c_{p, q}}\right)$ corresponds to some $f(X)$ and $A\left(V_{c_{p, q}} /\langle\chi\rangle\right) \cong C[X] /\langle f(X)\rangle$.
- For the Yang-Lee minimal model $p=2, q=5$

$$
\begin{aligned}
& \chi=\left(L_{-2}^{2}-\frac{3}{5} L_{-4}\right) \Omega \\
& Y(\chi ; z)=: T(z)^{2}:-\frac{3}{10} \partial^{2} T(z)
\end{aligned}
$$

Affine example:

- Let $\mathfrak{g}$ be a complex finite dimensional simple Lie algebra and let $V_{k}(\mathfrak{g})$ be the universal affine vertex operator algebra (conformal vector given by the Sugawara construction). Then $U(\mathfrak{g}) \cong A\left(V_{k}(\mathfrak{g})\right)$, where isomorphism is given by $x \mapsto\left[x_{-1} \Omega\right]$.
- If $k \in \mathbb{Z}_{\geq 0},\left(e_{-1}^{\theta}\right)^{k+1} \Omega$ is singular and generates the maximal ideal. $A\left(V_{k}(\mathfrak{g}) /\left\langle\left(e_{-1}^{\theta}\right)^{k+1} \Omega\right\rangle\right) \cong U(\mathfrak{g}) /\left\langle\left(e^{\theta}\right)^{k+1}\right\rangle$ is finite dimensional and semi simple.
- $\mathfrak{s l}_{2}$ at $k=-\frac{4}{3}$. The singular vector
$\chi=\left(h_{-3}+3 e_{-2} f_{-1}-3 e_{-1} f_{-2}+\frac{9}{2} h_{-1} e_{-1} f_{-1}+\frac{9}{8} h_{-1}^{3}-\frac{9}{4} h_{-2} h_{-1}\right) \Omega$ generates the unique non-trivial ideal.


## Literature

- Vertex algebras as rings: Frenkel Ben-Zvi, Chapters 1-5, http://dx.doi.org/10.1090/surv/088
- Fusion/tensor product theory: Huang Lepowsky Zhang 8 Part Series: arXiv:1012.4193, arXiv:1012.4196, arXiv:1012.4197, arXiv:1012.4198, arXiv:1012.4199, arXiv:1012.4202, arXiv:1110.1929, arXiv:1110.1931
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