

String and Chern-Simons Lie 3-Algebras

Urs Schreiber

with

Jim Stasheff

based in parts on work with

John Baez

Alissa Crans

David Roberts

Danny Stevenson

Konrad Waldorf

August 17, 2007

Plan

The Talk

- 1 Motivation
- 2 Connections with values in Lie n -Algebras
- 3 Conclusion

Further material

- Consequences
- Further topics

Thanks

to

- John Baez
- Bruce Bartlett
- Christoph Schweigert
- Jim Stasheff

for helpful comments on earlier versions of this talk.

Motivation

How can we understand quantum Chern-Simons theory as a 3-functorial Quantum Field Theory
– such that it allows us to derive the TFT construction of 2-dimensional CFT [Fuchs,Runkel,Schweigert] from first principles ?

Strategy

Proceed in two steps:

1. Understand classical Chern-Simons parallel transport.
2. Quantize.

This talk is about step 1.

Motivation

How can we understand quantum Chern-Simons theory as a 3-**functorial** Quantum Field Theory

– such that it allows us to derive the TFT construction of 2-dimensional CFT [Fuchs,Runkel,Schweigert] from first principles ?

Strategy

Proceed in two steps:

1. Understand classical Chern-Simons parallel transport.
2. Quantize.

This talk is about step 1.

Motivation

How can we understand quantum Chern-Simons theory as a **3-**functorial Quantum Field Theory

– such that it allows us to derive the TFT construction of 2-dimensional CFT [Fuchs,Runkel,Schweigert] from first principles ?

Strategy

Proceed in two steps:

1. Understand classical Chern-Simons parallel transport.
2. Quantize.

This talk is about step 1.

Motivation

How can we understand quantum Chern-Simons theory as a 3-functorial Quantum Field Theory
– such that it allows us to derive the TFT construction of 2-dimensional CFT [Fuchs,Runkel,Schweigert] from first principles ?

Strategy

Proceed in two steps:

1. Understand classical Chern-Simons parallel transport.
2. Quantize.

This talk is about step 1.

Motivation

How can we understand quantum Chern-Simons theory as a 3-functorial Quantum Field Theory

– such that it allows us to derive the TFT construction of 2-dimensional CFT [Fuchs,Runkel,Schweigert] from first principles ?

Strategy

Proceed in two steps:

- 1 Understand classical Chern-Simons parallel 3-transport.
- 2 Quantize.

This talk is about step 1.

Motivation

How can we understand quantum Chern-Simons theory as a 3-functorial Quantum Field Theory
– such that it allows us to derive the TFT construction of 2-dimensional CFT [Fuchs,Runkel,Schweigert] from first principles ?

Strategy

Proceed in two steps:

- 1 Understand classical Chern-Simons parallel 3-transport.
- 2 Quantize.

This talk is about step 1.

Motivation

How can we understand quantum Chern-Simons theory as a 3-functorial Quantum Field Theory
– such that it allows us to derive the TFT construction of 2-dimensional CFT [Fuchs,Runkel,Schweigert] from first principles ?

Strategy

Proceed in two steps:

- 1 Understand classical Chern-Simons parallel 3-transport.
- 2 Quantize.

This talk is about step 1.

Motivation

How can we understand quantum Chern-Simons theory as a 3-functorial Quantum Field Theory
– such that it allows us to derive the TFT construction of 2-dimensional CFT [Fuchs,Runkel,Schweigert] from first principles ?

Strategy

Proceed in two steps:

- 1 Understand classical Chern-Simons parallel 3-transport.
- 2 Quantize. (Or rather *3-quantize*.)

This talk is about step 1.

Motivation

How can we understand quantum Chern-Simons theory as a 3-functorial Quantum Field Theory
– such that it allows us to derive the TFT construction of 2-dimensional CFT [Fuchs,Runkel,Schweigert] from first principles ?

Strategy

Proceed in two steps:

- 1 Understand classical Chern-Simons parallel 3-transport.
- 2 Quantize.

This talk is about step 1.

A Quantum Field Theory is a Functor

- Atiyah and Segal have famously axiomatized d -dimensional QFTs
- as functors

$$Z : n\text{Cob}_S \rightarrow \text{Vect}$$

■

$$Z : \left(\partial_{\text{in}} \Sigma \xrightarrow{(\Sigma, g)} \partial_{\text{out}} \Sigma \right) \mapsto \left(H_{\text{in}} \xrightarrow{U(\Sigma, g)} H_{\text{out}} \right).$$

A Quantum Field Theory is a Functor

- Atiyah and Segal have famously axiomatized d -dimensional QFTs
- as functors

$$Z : n\text{Cob}_S \rightarrow \text{Vect}$$

■

$$Z : \left(\partial_{\text{in}} \Sigma \xrightarrow{(\Sigma, g)} \partial_{\text{out}} \Sigma \right) \mapsto \left(H_{\text{in}} \xrightarrow{U(\Sigma, g)} H_{\text{out}} \right).$$

A Quantum Field Theory is a Functor

- Atiyah and Segal have famously axiomatized d -dimensional QFTs
- as functors

$$Z : n\text{Cob}_S \rightarrow \text{Vect}$$

■

$$Z : \left(\partial_{\text{in}} \Sigma \xrightarrow{(\Sigma, g)} \partial_{\text{out}} \Sigma \right) \mapsto \left(H_{\text{in}} \xrightarrow{U(\Sigma, g)} H_{\text{out}} \right).$$

A Quantum Field Theory is a Functor

- Atiyah and Segal have famously axiomatized d -dimensional QFTs
- as functors

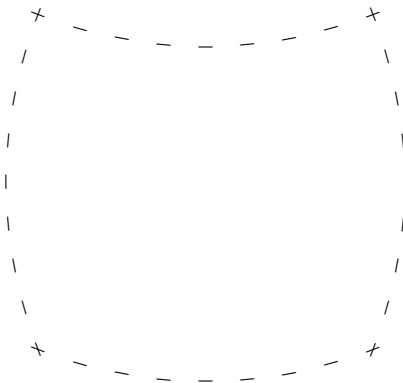
$$Z : n\text{Cob}_S \rightarrow \text{Vect}$$

-

$$Z : \left(\partial_{\text{in}} \Sigma \xrightarrow{(\Sigma, g)} \partial_{\text{out}} \Sigma \right) \mapsto \left(H_{\text{in}} \xrightarrow{U(\Sigma, g)} H_{\text{out}} \right) .$$

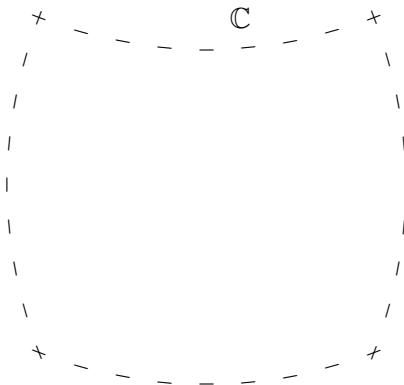
Cartoon of a 1-functorial QFT

$$\langle \phi | U(\Sigma) | \psi \rangle$$



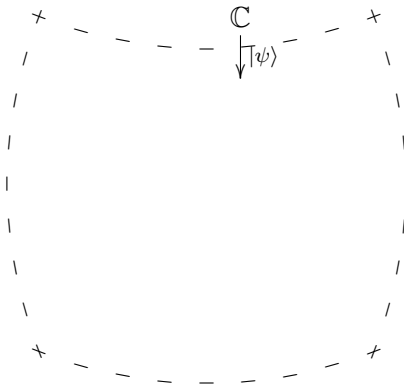
Cartoon of a 1-functorial QFT

$$\langle \phi | U(\Sigma) | \psi \rangle$$



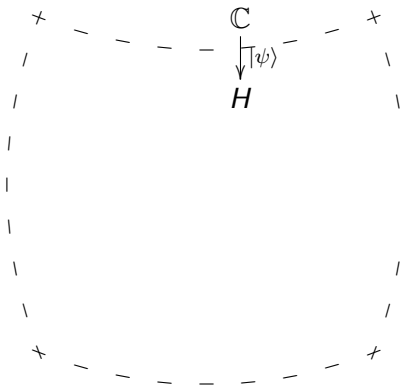
Cartoon of a 1-functorial QFT

$$\langle \phi | U(\Sigma) | \psi \rangle$$



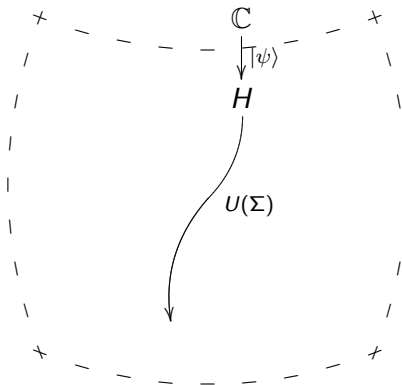
Cartoon of a 1-functorial QFT

$$\langle \phi | U(\Sigma) | \psi \rangle$$



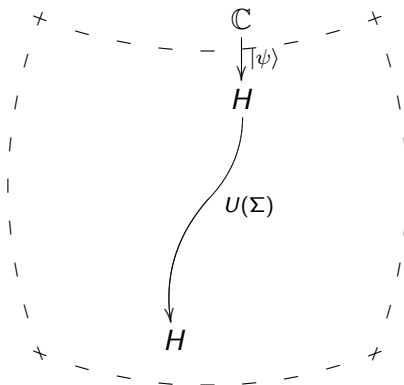
Cartoon of a 1-functorial QFT

$$\langle \phi | U(\Sigma) | \psi \rangle$$



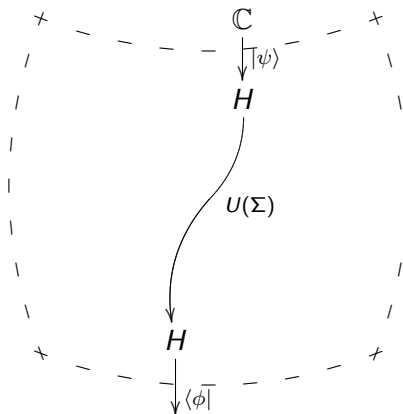
Cartoon of a 1-functorial QFT

$$\langle \phi | U(\Sigma) | \psi \rangle$$



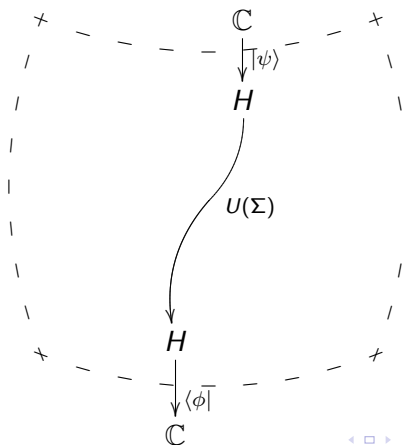
Cartoon of a 1-functorial QFT

$$\langle \phi | U(\Sigma) | \psi \rangle$$



Cartoon of a 1-functorial QFT

$$\langle \phi | U(\Sigma) | \psi \rangle$$



A Quantum Field Theory is an n -Functor

But later it was noticed that this is too imprecise if we want to be able to talk about

crucial requirements on QFT description

- *locality*
- *boundary conditions.*

Instead:

refined picture

An n -dimensional QFT should be an n -functor.
[Freed, Hopkins, Stolz, Teichner]

(remark on n -categories)

A Quantum Field Theory is an n -Functor

But later it was noticed that this is too imprecise if we want to be able to talk about

crucial requirements on QFT description

- *locality*
- *boundary conditions.*

Instead:

refined picture

An n -dimensional QFT should be an n -functor.
[Freed, Hopkins, Stolz, Teichner]

(remark on n -categories)

A Quantum Field Theory is an n -Functor

But later it was noticed that this is too imprecise if we want to be able to talk about

crucial requirements on QFT description

- *locality*
- *boundary conditions.*

Instead:

refined picture

An n -dimensional QFT should be an n -functor.
[Freed, Hopkins, Stolz, Teichner]

(remark on n -categories)

A Quantum Field Theory is an n -Functor

But later it was noticed that this is too imprecise if we want to be able to talk about

crucial requirements on QFT description

- *locality*
- *boundary conditions.*

Instead:

refined picture

An n -dimensional QFT should be an n -functor.
[Freed, Hopkins, Stolz, Teichner]

(remark on n -categories)

A Quantum Field Theory is an n -Functor

But later it was noticed that this is too imprecise if we want to be able to talk about

crucial requirements on QFT description

- *locality*
- *boundary conditions.*

Instead:

refined picture

An n -dimensional QFT should be an n -functor.
[Freed, Hopkins, Stolz, Teichner]

(remark on n -categories)

A Quantum Field Theory is an n -Functor

But later it was noticed that this is too imprecise if we want to be able to talk about

crucial requirements on QFT description

- *locality*
- *boundary conditions.*

Instead:

refined picture

An n -dimensional QFT should be an **n -functor**.

[Freed, Hopkins, Stolz, Teichner]

(remark on n -categories)

A Quantum Field Theory is an n -Functor

But later it was noticed that this is too imprecise if we want to be able to talk about

crucial requirements on QFT description

- *locality*
- *boundary conditions.*

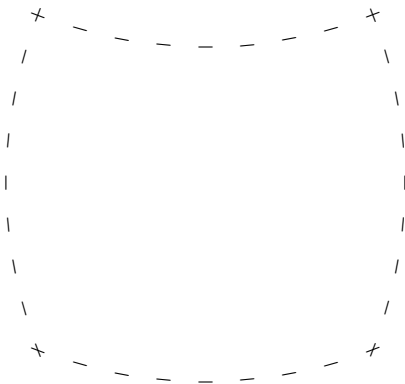
Instead:

refined picture

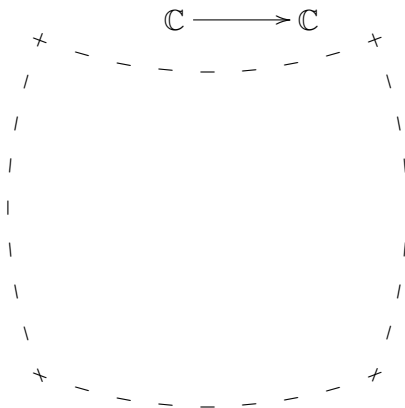
An n -dimensional QFT should be an n -functor.
[Freed, Hopkins, Stolz, Teichner]

(remark on n -categories)

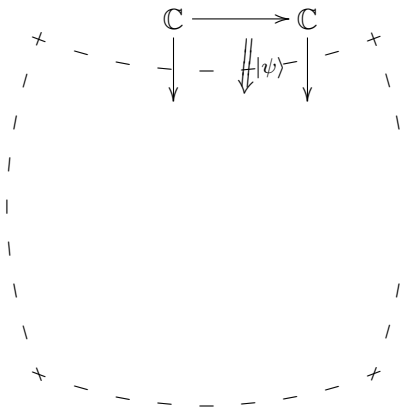
Cartoon of a 2-functorial QFT



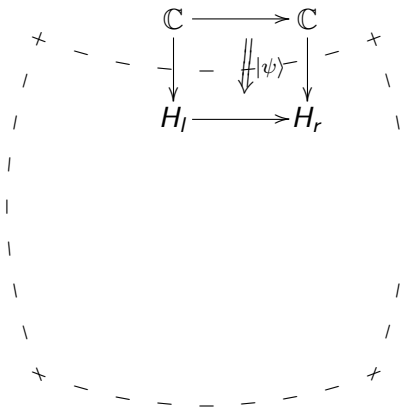
Cartoon of a 2-functorial QFT



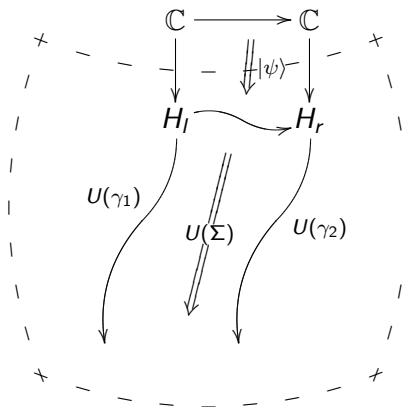
Cartoon of a 2-functorial QFT



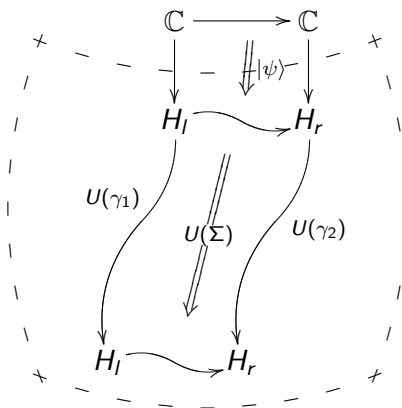
Cartoon of a 2-functorial QFT



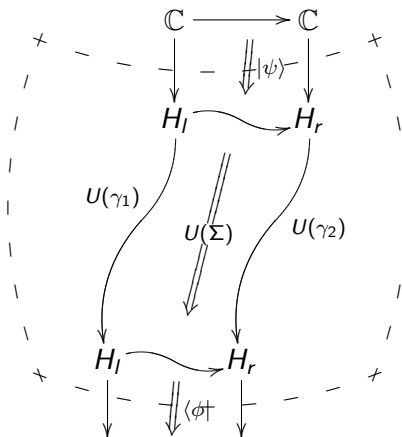
Cartoon of a 2-functorial QFT



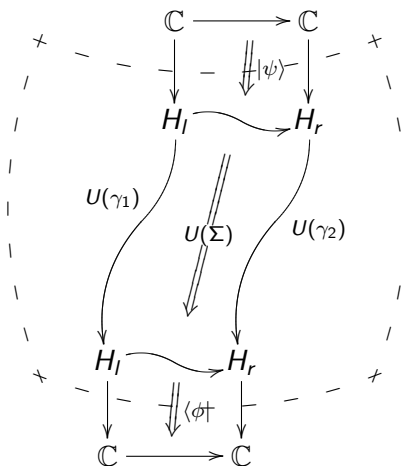
Cartoon of a 2-functorial QFT



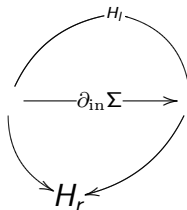
Cartoon of a 2-functorial QFT



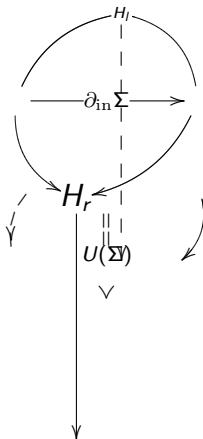
Cartoon of a 2-functorial QFT



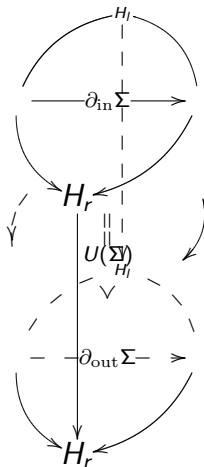
Cartoon of a 3-functorial QFT



Cartoon of a 3-functorial QFT



Cartoon of a 3-functorial QFT



n -Particles and $(n - 1)$ -Branes

It follows that the action of the n -particle. . .

n -Particle

- $n = 1$: the point particle
- $n = 2$: the string
- $n = 3$: the membrane
- n -particle $\simeq (n - 1)$ -brane

n -Particles and $(n - 1)$ -Branes

It follows that the action of the n -particle. . .

n -Particle

- $n = 1$: the point particle
- $n = 2$: the string
- $n = 3$: the membrane
- n -particle $\simeq (n - 1)$ -brane

n -Particles and $(n - 1)$ -Branes

It follows that the action of the n -particle. . .

n -Particle

- $n = 1$: the point particle
- $n = 2$: the string
- $n = 3$: the membrane
- n -particle $\simeq (n - 1)$ -brane

n -Particles and $(n - 1)$ -Branes

It follows that the action of the n -particle. . .

n -Particle

- $n = 1$: the point particle
- $n = 2$: the string
- $n = 3$: the membrane
- n -particle $\simeq (n - 1)$ -brane

n -Particles and $(n - 1)$ -Branes

It follows that the action of the n -particle. . .

n -Particle

- $n = 1$: the point particle
- $n = 2$: the string
- $n = 3$: the membrane
- n -particle $\simeq (n - 1)$ -brane

n -Bundles and $(n - 1)$ -Gerbes

It follows that the action of the n -particle charged under an n -bundle with connection...

n -background fields

- $n = 1$: the electromagnetic field
- $n = 2$: the Kalb-Ramond field
- $n = 3$: the supergravity 3-form field
- n -bundle $\simeq (n - 1)$ -gerbe

n -Bundles and $(n - 1)$ -Gerbes

It follows that the action of the n -particle charged under an n -bundle with connection...

n -background fields

- $n = 1$: the electromagnetic field
- $n = 2$: the Kalb-Ramond field
- $n = 3$: the supergravity 3-form field
- n -bundle $\simeq (n - 1)$ -gerbe

n -Bundles and $(n - 1)$ -Gerbes

It follows that the action of the n -particle charged under an n -bundle with connection...

n -background fields

- $n = 1$: the electromagnetic field
- $n = 2$: the Kalb-Ramond field
- $n = 3$: the supergravity 3-form field
- n -bundle $\simeq (n - 1)$ -gerbe

n -Bundles and $(n - 1)$ -Gerbes

It follows that the action of the n -particle charged under an n -bundle with connection...

n -background fields

- $n = 1$: the electromagnetic field
- $n = 2$: the Kalb-Ramond field
- $n = 3$: the supergravity 3-form field
- n -bundle $\simeq (n - 1)$ -gerbe

n -Bundles and $(n - 1)$ -Gerbes

It follows that the action of the n -particle charged under an n -bundle with connection...

n -background fields

- $n = 1$: the electromagnetic field
- $n = 2$: the Kalb-Ramond field
- $n = 3$: the supergravity 3-form field
- n -bundle $\simeq (n - 1)$ -gerbe

Parallel n -Transport

It follows that the action of the n -particle
charged under an n -bundle with connection
is itself an n -functor

$$\begin{aligned}
 \blacksquare \quad \text{tra}_1 : \left(x \xrightarrow{\gamma} y \right) &\mapsto \left(V_x \xrightarrow{P \exp \left(\int_{\gamma} A \right)} V_y \right) \\
 \blacksquare \quad \text{tra}_2 : \left(x \begin{array}{c} \xrightarrow{\gamma_1} \\ \Downarrow \Sigma \\ \xrightarrow{\gamma_2} \end{array} y \right) &\mapsto \left(V_x \begin{array}{c} \xrightarrow{P \exp \left(\int_{\gamma_1} A \right)} \\ \Downarrow P_A \exp \left(\int_{\Sigma} B \right) \\ \xrightarrow{P \exp \left(\int_{\gamma_2} A \right)} \end{array} V_y \right)
 \end{aligned}$$

Parallel n -Transport

It follows that the action of the n -particle charged under an n -bundle with connection is itself an n -functor

$$\begin{aligned}
 \blacksquare \quad \text{tra}_1 : \left(x \xrightarrow{\gamma} y \right) &\mapsto \left(V_x \xrightarrow{P \exp \left(\int_{\gamma} A \right)} V_y \right) \\
 \blacksquare \quad \text{tra}_2 : \left(x \begin{array}{c} \xrightarrow{\gamma_1} \\ \Downarrow \Sigma \\ \xrightarrow{\gamma_2} \end{array} y \right) &\mapsto \left(V_x \begin{array}{c} \xrightarrow{P \exp \left(\int_{\gamma_1} A \right)} \\ \Downarrow P_A \exp \left(\int_{\Sigma} B \right) \\ \xrightarrow{P \exp \left(\int_{\gamma_2} A \right)} \end{array} V_y \right)
 \end{aligned}$$

Parallel n -Transport

It follows that the action of the n -particle charged under an n -bundle with connection is itself an n -functor

$$\blacksquare \quad \text{tra}_1 : \left(x \xrightarrow{\gamma} y \right) \mapsto \left(V_x \xrightarrow{P \exp \left(\int_{\gamma} A \right)} V_y \right)$$

$$\blacksquare \quad \text{tra}_2 : \left(x \begin{array}{c} \xrightarrow{\gamma_1} \\ \Downarrow \Sigma \\ \xrightarrow{\gamma_2} \end{array} y \right) \mapsto \left(V_x \begin{array}{c} \xrightarrow{P \exp \left(\int_{\gamma_1} A \right)} \\ \Downarrow P_A \exp \left(\int_{\Sigma} B \right) \\ \xrightarrow{P \exp \left(\int_{\gamma_2} A \right)} \end{array} V_y \right)$$

Parallel n -Transport

It follows that the action of the n -particle charged under an n -bundle with connection is itself an n -functor

$$\begin{aligned}
 \blacksquare \quad \text{tra}_1 : \left(x \xrightarrow{\gamma} y \right) &\mapsto \left(V_x \xrightarrow{P \exp \left(\int_{\gamma} A \right)} V_y \right) \\
 \blacksquare \quad \text{tra}_2 : \left(x \begin{array}{c} \xrightarrow{\gamma_1} \\ \Downarrow \Sigma \\ \xrightarrow{\gamma_2} \end{array} y \right) &\mapsto \left(V_x \begin{array}{c} \xrightarrow{P \exp \left(\int_{\gamma_1} A \right)} \\ \text{\scriptsize $P_A \exp \left(\int_{\Sigma} B \right)$} \Downarrow \\ \xrightarrow{P \exp \left(\int_{\gamma_2} A \right)} \end{array} V_y \right)
 \end{aligned}$$

Parallel 3-Transport

It follows that the action of the *3-particle* charged under a *3-bundle with connection* is itself a 3-functor

$$\blacksquare \quad \text{tra}_3: \left(\begin{array}{c} \text{Diagram 1} \end{array} \right) \mapsto \left(\begin{array}{c} \text{Diagram 2} \end{array} \right)$$

Diagram 1 (Left): A circular diagram with two nodes, γ_1 and γ_2 . A solid arrow labeled x goes from γ_1 to γ_2 along the top arc. A solid arrow labeled y goes from γ_2 to γ_1 along the bottom arc. A dashed arrow labeled V goes from γ_1 to γ_2 through the center. The top arc is labeled Σ_1 and the bottom arc is labeled Σ_2 . A vertical double arrow labeled V connects Σ_1 and Σ_2 .

Diagram 2 (Right): A circular diagram with two nodes, V_x and V_y . A solid arrow labeled V_x goes from V_x to V_y along the top arc. A solid arrow labeled V_y goes from V_y to V_x along the bottom arc. A dashed arrow labeled $P_{A,B} \exp(\int_V C)$ goes from V_x to V_y through the center. The top arc is labeled $P_{A,B} \exp(\int_V C)$ and the bottom arc is labeled $-$. A vertical double arrow labeled $-$ connects the top and bottom arcs.

Parallel 3-Transport

It follows that the action of the *3-particle* charged under a *3-bundle with connection* is itself a 3-functor

$$\blacksquare \quad \text{tra}_3: \left(\begin{array}{c} \text{Diagram 1} \end{array} \right) \mapsto \left(\begin{array}{c} \text{Diagram 2} \end{array} \right)$$

Diagram 1 (Left): A circular diagram with two nodes, γ_1 and γ_2 , connected by two curved arrows. The top arrow is labeled Σ_1 and the bottom arrow is labeled Σ_2 . A dashed arrow labeled y points from γ_2 to γ_1 . A vertical double arrow labeled V points from Σ_1 to Σ_2 .

Diagram 2 (Right): A circular diagram with two nodes, V_x and V_y , connected by two curved arrows. The top arrow is labeled V_x and the bottom arrow is labeled V_y . A dashed arrow points from V_y to V_x . A vertical double arrow labeled $P_{A,B} \exp(\int_V C)$ points from the top arrow to the bottom arrow.

Parallel n -Transport

A **parallel n -transport** is (locally) an n -functor from the path n -groupoid to the structure n -group.

$$\mathbf{tra}_n : \mathcal{P}_n(X) \rightarrow \Sigma G_{(n)}$$

$(n+1)$ -Curvature

Its $(n+1)$ -curvature is (locally) an $(n+1)$ -functor from the fundamental $(n+1)$ -groupoid to the inner automorphism $(n+1)$ -group of $G_{(n)}$.

$$d\mathbf{tra}_n := \mathbf{curv}_{(n+1)} : \Pi_{n+1}(X) \rightarrow \Sigma(\mathrm{INN} G_{(n)})$$

Parallel n -Transport

A parallel n -transport **is** (locally) an n -functor from the path n -groupoid to the structure n -group.

$$\mathrm{tra}_n : \mathcal{P}_n(X) \rightarrow \Sigma G_{(n)}$$

$(n+1)$ -Curvature

Its $(n+1)$ -curvature is (locally) an $(n+1)$ -functor from the fundamental $(n+1)$ -groupoid to the inner automorphism $(n+1)$ -group of $G_{(n)}$.

$$d\mathrm{tra}_n := \mathrm{curv}_{(n+1)} : \Pi_{n+1}(X) \rightarrow \Sigma(\mathrm{INN}G_{(n)})$$

Parallel n -Transport

A parallel n -transport is (locally) an n -functor from the path n -groupoid to the structure n -group.

$$\mathrm{tra}_n : \mathcal{P}_n(X) \rightarrow \Sigma G_{(n)}$$

$(n+1)$ -Curvature

Its $(n+1)$ -curvature is (locally) an $(n+1)$ -functor from the fundamental $(n+1)$ -groupoid to the inner automorphism $(n+1)$ -group of $G_{(n)}$.

$$d\mathrm{tra}_n := \mathrm{curv}_{(n+1)} : \Pi_{n+1}(X) \rightarrow \Sigma(\mathrm{INN} G_{(n)})$$

Parallel n -Transport

A parallel n -transport is (locally) an n -functor from the path n -groupoid to the structure n -group.

$$\mathrm{tra}_n : \mathcal{P}_n(X) \rightarrow \Sigma G_{(n)}$$

$(n+1)$ -Curvature

Its $(n+1)$ -curvature is (locally) an $(n+1)$ -functor from the fundamental $(n+1)$ -groupoid to the inner automorphism $(n+1)$ -group of $G_{(n)}$.

$$d\mathrm{tra}_n := \mathrm{curv}_{(n+1)} : \Pi_{n+1}(X) \rightarrow \Sigma(\mathrm{INN}G_{(n)})$$

Parallel n -Transport

A parallel n -transport is (locally) an n -functor from the path n -groupoid to the **structure n -group**.

$$\mathrm{tra}_n : \mathcal{P}_n(X) \rightarrow \Sigma G_{(n)}$$

$(n+1)$ -Curvature

Its $(n+1)$ -curvature is (locally) an $(n+1)$ -functor from the fundamental $(n+1)$ -groupoid to the inner automorphism $(n+1)$ -group of $G_{(n)}$.

$$d\mathrm{tra}_n := \mathrm{curv}_{(n+1)} : \Pi_{n+1}(X) \rightarrow \Sigma(\mathrm{INN} G_{(n)})$$

Parallel n -Transport

A parallel n -transport is (locally) an n -functor from the path n -groupoid to the structure n -group .

$$\mathrm{tra}_n : \mathcal{P}_n(X) \rightarrow \Sigma G_{(n)}$$

$(n+1)$ -Curvature

Its $(n+1)$ -curvature is (locally) an $(n+1)$ -functor from the fundamental $(n+1)$ -groupoid to the inner automorphism $(n+1)$ -group of $G_{(n)}$.

$$d\mathrm{tra}_n := \mathrm{curv}_{(n+1)} : \Pi_{n+1}(X) \rightarrow \Sigma(\mathrm{INN} G_{(n)})$$

Parallel n -Transport

A parallel n -transport is (locally) an n -functor from the path n -groupoid to the structure n -group .

$$\mathrm{tra}_n : \mathcal{P}_n(X) \rightarrow \Sigma G_{(n)}$$

$(n+1)$ -Curvature

Its $(n+1)$ -curvature is (locally) an $(n+1)$ -functor from the fundamental $(n+1)$ -groupoid to the inner automorphism $(n+1)$ -group of $G_{(n)}$.

$$d\mathrm{tra}_n := \mathrm{curv}_{(n+1)} : \Pi_{n+1}(X) \rightarrow \Sigma(\mathrm{INN} G_{(n)})$$

Parallel n -Transport

A parallel n -transport is (locally) an n -functor from the path n -groupoid to the structure n -group .

$$\mathrm{tra}_n : \mathcal{P}_n(X) \rightarrow \Sigma G_{(n)}$$

$(n+1)$ -Curvature

Its $(n+1)$ -curvature is (locally) **an $(n+1)$ -functor** from the fundamental $(n+1)$ -groupoid to the inner automorphism $(n+1)$ -group of $G_{(n)}$.

$$d\mathrm{tra}_n := \mathrm{curv}_{(n+1)} : \Pi_{n+1}(X) \rightarrow \Sigma(\mathrm{INN} G_{(n)})$$

Parallel n -Transport

A parallel n -transport is (locally) an n -functor from the path n -groupoid to the structure n -group .

$$\mathrm{tra}_n : \mathcal{P}_n(X) \rightarrow \Sigma G_{(n)}$$

$(n+1)$ -Curvature

Its $(n+1)$ -curvature is (locally) an $(n+1)$ -functor from the **fundamental $(n+1)$ -groupoid** to the inner automorphism $(n+1)$ -group of $G_{(n)}$.

$$d\mathrm{tra}_n := \mathrm{curv}_{(n+1)} : \Pi_{n+1}(X) \rightarrow \Sigma(\mathrm{INN} G_{(n)})$$

Parallel n -Transport

A parallel n -transport is (locally) an n -functor from the path n -groupoid to the structure n -group .

$$\mathrm{tra}_n : \mathcal{P}_n(X) \rightarrow \Sigma G_{(n)}$$

$(n+1)$ -Curvature

Its $(n+1)$ -curvature is (locally) an $(n+1)$ -functor from the fundamental $(n+1)$ -groupoid to the **inner automorphism $(n+1)$ -group** of $G_{(n)}$.

$$d\mathrm{tra}_n := \mathrm{curv}_{(n+1)} : \Pi_{n+1}(X) \rightarrow \Sigma(\mathrm{INN} G_{(n)})$$

Tangent Categories

Inner automorphism $(n+1)$ -Groups

- Every n -group $G_{(n)}$ has an $(n+1)$ -group $\text{AUT}(G_{(n)})$ of automorphisms.
- This sits inside an exact sequence

$$1 \rightarrow Z(G_{(n)}) \rightarrow \text{INN}(G_{(n)}) \rightarrow \text{AUT}(G_{(n)}) \rightarrow \text{OUT}(G_{(n)}) \rightarrow 1$$
- and INN_0 plays the role of the universal $G_{(n)}$ -bundle

$$G_{(n)} \rightarrow \text{INN}_0(G_{(n)}) \rightarrow \Sigma G_{(n)}$$

We will re-encounter these crucial facts in their Lie n -algebra incarnation shortly.

[U.S., David Roberts]

(on tangent categories) (on inner automorphisms)

Tangent Categories

Inner automorphism $(n+1)$ -Groups

- Every n -group $G_{(n)}$ has an $(n+1)$ -group $\text{AUT}(G_{(n)})$ of automorphisms.

- This sits inside an exact sequence

$$1 \rightarrow Z(G_{(n)}) \rightarrow \text{INN}(G_{(n)}) \rightarrow \text{AUT}(G_{(n)}) \rightarrow \text{OUT}(G_{(n)}) \rightarrow 1$$

- and INN_0 plays the role of the universal $G_{(n)}$ -bundle

$$G_{(n)} \rightarrow \text{INN}_0(G_{(n)}) \rightarrow \Sigma G_{(n)}$$

We will re-encounter these crucial facts in their Lie n -algebra incarnation shortly.

[U.S., David Roberts]

(on tangent categories) (on inner automorphisms)

Tangent Categories

Inner automorphism $(n+1)$ -Groups

- Every n -group $G_{(n)}$ has an $(n+1)$ -group $\text{AUT}(G_{(n)})$ of automorphisms.
- This sits inside an exact sequence

$$1 \rightarrow Z(G_{(n)}) \rightarrow \text{INN}(G_{(n)}) \rightarrow \text{AUT}(G_{(n)}) \rightarrow \text{OUT}(G_{(n)}) \rightarrow 1$$
- and INN_0 plays the role of the universal $G_{(n)}$ -bundle

$$G_{(n)} \rightarrow \text{INN}_0(G_{(n)}) \rightarrow \Sigma G_{(n)}$$

We will re-encounter these crucial facts in their Lie n -algebra incarnation shortly.

[U.S., David Roberts]

(on tangent categories) (on inner automorphisms)

Tangent Categories

Inner automorphism $(n+1)$ -Groups

- Every n -group $G_{(n)}$ has an $(n+1)$ -group $\text{AUT}(G_{(n)})$ of automorphisms.
- This sits inside an exact sequence

$$1 \rightarrow Z(G_{(n)}) \rightarrow \text{INN}(G_{(n)}) \rightarrow \text{AUT}(G_{(n)}) \rightarrow \text{OUT}(G_{(n)}) \rightarrow 1$$
- and INN_0 plays the role of the universal $G_{(n)}$ -bundle

$$G_{(n)} \rightarrow \text{INN}_0(G_{(n)}) \rightarrow \Sigma G_{(n)}$$

We will re-encounter these crucial facts in their Lie n -algebra incarnation shortly.

[U.S., David Roberts]

(on tangent categories) (on inner automorphisms)

Tangent Categories

Inner automorphism $(n+1)$ -Groups

- Every n -group $G_{(n)}$ has an $(n+1)$ -group $\text{AUT}(G_{(n)})$ of automorphisms.
- This sits inside an exact sequence

$$1 \rightarrow Z(G_{(n)}) \rightarrow \text{INN}(G_{(n)}) \rightarrow \text{AUT}(G_{(n)}) \rightarrow \text{OUT}(G_{(n)}) \rightarrow 1$$
- and INN_0 plays the role of the universal $G_{(n)}$ -bundle

$$G_{(n)} \rightarrow \text{INN}_0(G_{(n)}) \rightarrow \Sigma G_{(n)}$$

We will re-encounter these crucial facts in their Lie n -algebra incarnation shortly.

[U.S., David Roberts]

(on tangent categories) (on inner automorphisms)

Tangent Categories

Inner automorphism $(n+1)$ -Groups

- Every n -group $G_{(n)}$ has an $(n+1)$ -group $\text{AUT}(G_{(n)})$ of automorphisms.
- This sits inside an exact sequence

$$1 \rightarrow Z(G_{(n)}) \rightarrow \text{INN}(G_{(n)}) \rightarrow \text{AUT}(G_{(n)}) \rightarrow \text{OUT}(G_{(n)}) \rightarrow 1$$
- and INN_0 plays the role of the universal $G_{(n)}$ -bundle

$$G_{(n)} \rightarrow \text{INN}_0(G_{(n)}) \rightarrow \Sigma G_{(n)}$$

We will re-encounter these crucial facts in their Lie n -algebra incarnation shortly.

[U.S., David Roberts]

(on tangent categories) (on inner automorphisms)

Tangent Categories

Inner automorphism $(n+1)$ -Groups

- Every n -group $G_{(n)}$ has an $(n+1)$ -group $\text{AUT}(G_{(n)})$ of automorphisms.
- This sits inside an exact sequence

$$1 \rightarrow Z(G_{(n)}) \rightarrow \text{INN}(G_{(n)}) \rightarrow \text{AUT}(G_{(n)}) \rightarrow \text{OUT}(G_{(n)}) \rightarrow 1$$
- and INN_0 plays the role of the universal $G_{(n)}$ -bundle

$$G_{(n)} \rightarrow \text{INN}_0(G_{(n)}) \rightarrow \Sigma G_{(n)}$$

We will re-encounter these crucial facts in their Lie n -algebra incarnation shortly.

[U.S., David Roberts]

(on tangent categories) (on inner automorphisms)

Some structure n -Groups

Important structure (1-)Groups

electrically charged 1-particle: $G_{(1)} = U(1)$

spinning 1-particle: $G_{(1)} = \text{Spin}(n)$

Important structure (2-)Groups

electrically charged 2-particle: $G_{(2)} = U(1)$

spinning 2-particle: $G_{(2)} = \text{Spin}(n), \text{Spin}(n)$

Important Structure 3-Groups

Chern-Simons charged 3-particle: $G_{(3)} = \mathbb{Z}$

Tough question. Let's pass to the differential picture.

Some structure n -Groups

Important structure (1-)Groups

electrically charged 1-particle: $G_{(1)} = U(1)$

spinning 1-particle: $G_{(1)} = \text{Spin}(n)$

Important structure (2-)Groups

electrically charged 2-particle: $G_{(2)} = \text{SU}(2) \times \text{SU}(2)$

Important Structure 3-Groups

Chern-Simons charged 3-particle: $G_{(3)} = \mathbb{Z}$

Tough question. Let's pass to the differential picture.

Some structure n -Groups

Important structure (1-)Groups

electrically charged 1-particle: $G_{(1)} = U(1)$

spinning 1-particle: $G_{(1)} = \text{Spin}(n)$

Important structure (2-)Groups

electrically charged 2-particle: $G_{(2)} = \text{SU}(2) \ltimes \text{SU}(2)$

Important Structure 3-Groups

Chern-Simons charged 3-particle: $G_{(3)} = \mathbb{Z}$

Tough question. Let's pass to the differential picture.

Some structure n -Groups

Important structure (1-)Groups

electrically charged 1-particle: $G_{(1)} = U(1)$

spinning 1-particle: $G_{(1)} = \text{Spin}(n)$

Important structure (2-)Groups

Chern-Simons charged 2-particle: $G_{(2)} = \text{Spin}(n, \mathbb{R})$

Important Structure 3-Groups

Chern-Simons charged 3-particle: $G_{(3)} = ?$

Tough question. Let's pass to the differential picture.

Some structure n -Groups

Important structure (1-)Groups

electrically charged 1-particle: $G_{(1)} = U(1)$

spinning 1-particle: $G_{(1)} = \text{Spin}(n)$

Important structure (2-)Groups

Important Structure 3-Groups

Chern-Simons charged 3-particle: $G_{(3)} = \mathbb{Z}$

Tough question. Let's pass to the differential picture.

Some structure n -Groups

Important structure (1-)Groups

electrically charged 1-particle: $G_{(1)} = U(1)$

spinning 1-particle: $G_{(1)} = \text{Spin}(n)$

Important structure (2-)Groups

Kalb-Ramond charged 2-particle: $G_{(2)} = \Sigma U(1)$

spinning 2-particle: $G_{(2)} = \text{String}_k(\text{Spin}(n))$

[Bartels],[Baez,S],[S,Waldorf]

Important Structure 3-Groups

Tough question. Let's pass to the differential picture.

Some structure n -Groups

Important structure (1-)Groups

electrically charged 1-particle: $G_{(1)} = U(1)$

spinning 1-particle: $G_{(1)} = \text{Spin}(n)$

Important structure (2-)Groups

Kalb-Ramond charged 2-particle: $G_{(2)} = \Sigma U(1)$

spinning 2-particle: $G_{(2)} = \text{String}_k(\text{Spin}(n))$

[Bartels],[Baez,S],[S,Waldorf]

Important Structure 3-Groups

Tough question. Let's pass to the differential picture.

Some structure n -Groups

Important structure (1-)Groups

electrically charged 1-particle: $G_{(1)} = U(1)$

spinning 1-particle: $G_{(1)} = \text{Spin}(n)$

Important structure (2-)Groups

Kalb-Ramond charged 2-particle: $G_{(2)} = \Sigma U(1)$

spinning 2-particle: $G_{(2)} = \text{String}_k(\text{Spin}(n))$

[Baez,Crans,S,Stevenson]

Important Structure 3-Groups

Tough question. Let's pass to the differential picture.

Some structure n -Groups

Important structure (1-)Groups

electrically charged 1-particle: $G_{(1)} = U(1)$

spinning 1-particle: $G_{(1)} = \text{Spin}(n)$

Important structure (2-)Groups

Kalb-Ramond charged 2-particle: $G_{(2)} = \Sigma U(1)$

spinning 2-particle: $G_{(2)} = \text{String}_k(\text{Spin}(n))$

[Baez,Crans,S,Stevenson]

Important Structure 3-Groups

Tough question. Let's pass to the differential picture.

Some structure n -Groups

Important structure (1-)Groups

electrically charged 1-particle: $G_{(1)} = U(1)$

spinning 1-particle: $G_{(1)} = \text{Spin}(n)$

Important structure (2-)Groups

Kalb-Ramond charged 2-particle: $G_{(2)} = \Sigma U(1)$

spinning 2-particle: $G_{(2)} = \text{String}_k(\text{Spin}(n))$

Important Structure 3-Groups

Chern-Simons charged 3-particle: $G_{(3)} = ?$

Tough question. Let's pass to the differential picture.

Some structure n -Groups

Important structure (1-)Groups

electrically charged 1-particle: $G_{(1)} = U(1)$

spinning 1-particle: $G_{(1)} = \text{Spin}(n)$

Important structure (2-)Groups

Kalb-Ramond charged 2-particle: $G_{(2)} = \Sigma U(1)$

spinning 2-particle: $G_{(2)} = \text{String}_k(\text{Spin}(n))$

Important Structure 3-Groups

Chern-Simons charged 3-particle: $G_{(3)} = ?$

Tough question. Let's pass to the differential picture.

Some structure n -Groups

Important structure (1-)Groups

electrically charged 1-particle: $G_{(1)} = U(1)$

spinning 1-particle: $G_{(1)} = \text{Spin}(n)$

Important structure (2-)Groups

Kalb-Ramond charged 2-particle: $G_{(2)} = \Sigma U(1)$

spinning 2-particle: $G_{(2)} = \text{String}_k(\text{Spin}(n))$

Important Structure 3-Groups

Chern-Simons charged 3-particle: $G_{(3)} = ?$

Tough question. Let's pass to the differential picture.

Some structure n -Groups

Important structure (1-)Groups

electrically charged 1-particle: $G_{(1)} = U(1)$

spinning 1-particle: $G_{(1)} = \text{Spin}(n)$

Important structure (2-)Groups

Kalb-Ramond charged 2-particle: $G_{(2)} = \Sigma U(1)$

spinning 2-particle: $G_{(2)} = \text{String}_k(\text{Spin}(n))$

Important Structure 3-Groups

Chern-Simons charged 3-particle: $G_{(3)} = ?$

Tough question.

Let's pass to the differential picture.

Finding the Chern-Simons Lie 3-algebra

Problem

Identify that class of 3-transport – given by its structure 3-group – which evaluates to the Chern-Simons functional on 3-dimensional morphisms.

Strategy

- Differentiate. Pass from Lie n -groups to Lie n -algebras.
- Find that Lie 3-algebra $\mathfrak{cs}_3(\mathfrak{g})$ with the property that connections taking values in it, $\text{Vect} \rightarrow \mathfrak{cs}_3(\mathfrak{g})$, correspond to triples (A, B, C) of forms such that $C = C_2(A) + \langle A, A \rangle$.

Finding the Chern-Simons Lie 3-algebra

Problem

Identify that class of 3-transport – given by its structure 3-group – which evaluates to the Chern-Simons functional on 3-dimensional morphisms.

Strategy

- Differentiate. Pass from Lie n -groups to Lie n -algebras.
- Find that Lie 3-algebra $\mathfrak{cs}_3(\mathfrak{g})$ with the property that connections taking values in it, restricted to $\mathfrak{cs}_3(\mathfrak{g})$, correspond to triples (A, B, C) of forms such that $C = C_2(A) + B \wedge A$.

Finding the Chern-Simons Lie 3-algebra

Problem

Identify that class of 3-transport – given by its structure 3-group –
 which evaluates to the Chern-Simons functional on 3-dimensional
 morphisms.

Strategy

- Differentiate. Pass from Lie groups to Lie algebras.
- Find that Lie 3-algebra $\mathfrak{cs}_3(\mathfrak{g})$ with the property that
 connections taking values in \mathfrak{g} evaluate to $\mathfrak{cs}_3(\mathfrak{g})$ correspond
 to triples (A, B, C) of forms such that $C = C_2(A)$.

Finding the Chern-Simons Lie 3-algebra

Problem

Identify that class of 3-transport – given by its structure 3-group – which evaluates to the Chern-Simons functional on 3-dimensional morphisms.

Strategy

- Differentiate. Pass from Lie n -groups to Lie n -algebras.
- Find that Lie 3-algebra $cs_k(\mathfrak{g})$ with the property that connections taking values in it, $\text{Vect} \rightarrow cs_k(\mathfrak{g})$, correspond to triples (A, B, C) of forms such that $C = CS_k(A) + dB$.

Finding the Chern-Simons Lie 3-algebra

Problem

Identify that class of 3-transport – given by its structure 3-group – which evaluates to the Chern-Simons functional on 3-dimensional morphisms.

Strategy

- Differentiate. Pass from Lie n -groups to Lie n -algebras.
- Find that Lie 3-algebra $\text{cs}_k(\mathfrak{g})$ with the property that connections taking values in it, $\text{Vect} \rightarrow \text{cs}_k(\mathfrak{g})$, correspond to triples (A, B, C) of forms such that $C = \text{CS}_k(A) + dB$.

Finding the Chern-Simons Lie 3-algebra

Problem

Identify that class of 3-transport – given by its structure 3-group – which evaluates to the Chern-Simons functional on 3-dimensional morphisms.

Strategy

- Differentiate. Pass from Lie n -groups to Lie n -algebras.
- Find that Lie 3-algebra $\text{cs}_k(\mathfrak{g})$ with the property that connections taking values in it, $\text{Vect} \rightarrow \text{cs}_k(\mathfrak{g})$, correspond to triples (A, B, C) of forms such that $C = \text{CS}_k(A) + dB$.

Finding the Chern-Simons Lie 3-algebra

Problem

Identify that class of 3-transport – given by its structure 3-group – which evaluates to the Chern-Simons functional on 3-dimensional morphisms.

Strategy

- Differentiate. Pass from Lie n -groups to Lie n -algebras.
- Find that Lie 3-algebra $\text{cs}_k(\mathfrak{g})$ with the property that connections taking values in it, $\text{Vect} \rightarrow \text{cs}_k(\mathfrak{g})$, correspond to triples (A, B, C) of forms such that $C = \text{CS}_k(A) + dB$.

Finding the Chern-Simons Lie 3-algebra

Problem

Identify that class of 3-transport – given by its structure 3-group – which evaluates to the Chern-Simons functional on 3-dimensional morphisms.

Strategy

- Differentiate. Pass from Lie n -groups to Lie n -algebras.
- Find that Lie 3-algebra $\text{cs}_k(\mathfrak{g})$ with the property that connections taking values in it, $\text{Vect} \rightarrow \text{cs}_k(\mathfrak{g})$, correspond to triples (A, B, C) of forms such that $C = \text{CS}_k(A) + dB$.

Finding the Chern-Simons Lie 3-algebra

Problem

Identify that class of 3-transport – given by its structure 3-group – which evaluates to the Chern-Simons functional on 3-dimensional morphisms.

Strategy

- Differentiate. Pass from Lie n -groups to Lie n -algebras.
- Find that Lie 3-algebra $\text{cs}_k(\mathfrak{g})$ with the property that connections taking values in it, $\text{Vect} \rightarrow \text{cs}_k(\mathfrak{g})$, correspond to triples (A, B, C) of forms such that $C = \text{CS}_k(A) + dB$.

Finding the Chern-Simons Lie 3-algebra

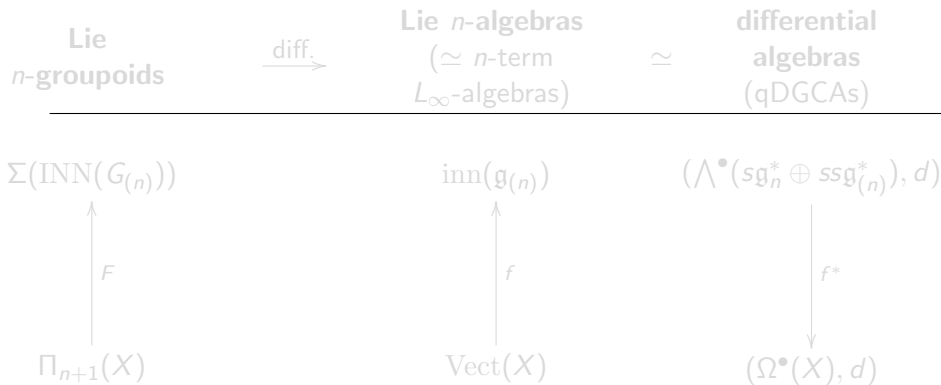
Problem

Identify that class of 3-transport – given by its structure 3-group – which evaluates to the Chern-Simons functional on 3-dimensional morphisms.

Strategy

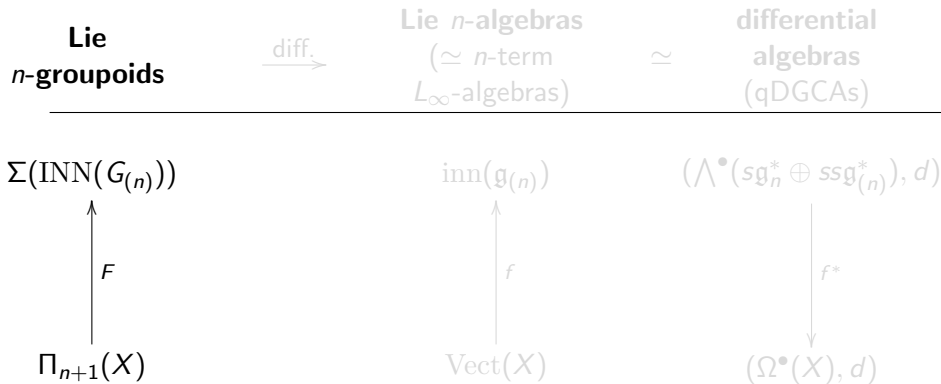
- Differentiate. Pass from Lie n -groups to Lie n -algebras.
- Find that Lie 3-algebra $\text{cs}_k(\mathfrak{g})$ with the property that connections taking values in it, $\text{Vect} \rightarrow \text{cs}_k(\mathfrak{g})$, correspond to triples (A, B, C) of forms such that $C = \text{CS}_k(A) + dB$.

From parallel n -transport to Lie n -algebra valued connections



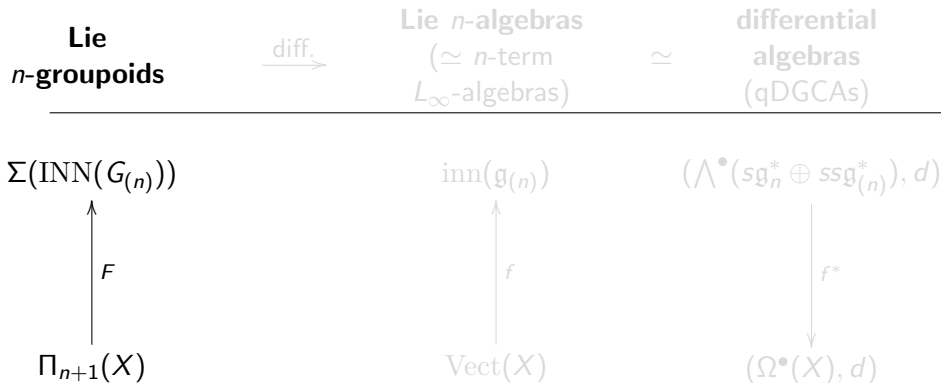
Parallel n -transport is a morphism of Lie $(n+1)$ -groupoids.

From parallel n -transport to Lie n -algebra valued connections



Parallel n -transport is a morphism of Lie $(n+1)$ -groupoids.

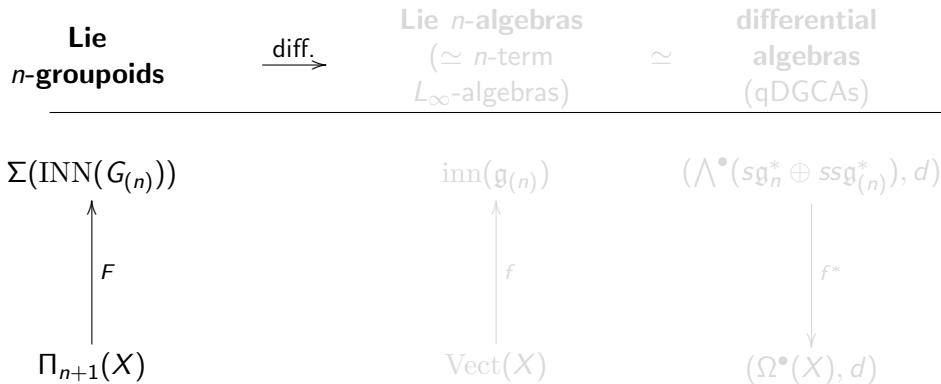
From parallel n -transport to Lie n -algebra valued connections



This morphism may be differentiated...



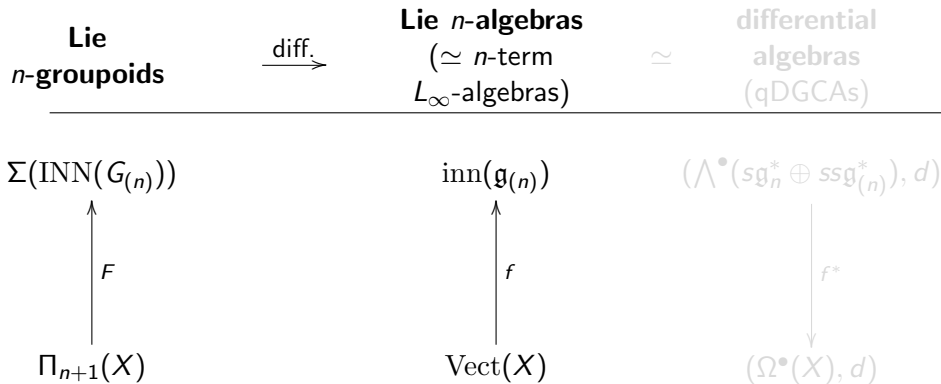
From parallel n -transport to Lie n -algebra valued connections



... to produce a morphism of Lie $(n+1)$ -algebroids.



From parallel n -transport to Lie n -algebra valued connections



These are best handled in terms of their dual maps,

From parallel n -transport to Lie n -algebra valued connections

$$\begin{array}{ccccc}
 \text{Lie } n\text{-groupoids} & \xrightarrow{\text{diff.}} & \text{Lie } n\text{-algebras} & \simeq & \text{differential algebras} \\
 & & (\simeq n\text{-term } L_\infty\text{-algebras}) & & (\text{qDGCAs}) \\
 \hline
 \Sigma(\text{INN}(G_{(n)})) & & \text{inn}(\mathfrak{g}_{(n)}) & & (\wedge^\bullet(\mathfrak{sg}_n^* \oplus \mathfrak{ssg}_{(n)}^*), d) \\
 \uparrow F & & \uparrow f & & \downarrow f^* \\
 \Pi_{n+1}(X) & & \text{Vect}(X) & & (\Omega^\bullet(X), d)
 \end{array}$$

which are morphisms of quasi-free differential-graded algebras.

The basic concepts

Lie 2-algebras

- Baez and Crans consider 2-vector spaces as categories internal to vector spaces.
- These are nothing but 2-term chain complexes – but interpreted suitably.
- They then define a Lie 2-algebra to be a Lie algebra internal to these 2-vector spaces.
- With strict skew symmetry but weak Jacobi identity.
- Notice one can also weaken the skew symmetry [Roytenberg].

The basic concepts

Lie 2-algebras

- Baez and Crans consider 2-vector spaces as categories internal to vector spaces.
- This are nothing but 2-term chain complexes – but interpreted suitably.
- They then define a Lie 2-algebra to be a Lie algebra internal to these 2-vector spaces.
- With strict skew symmetry but weak Jacobi identity.
- Notice one can also weaken the skew symmetry [Roytenberg].

The basic concepts

Lie 2-algebras

- Baez and Crans consider 2-vector spaces as categories internal to vector spaces.
- These are nothing but 2-term chain complexes – but interpreted suitably.
- They then define a Lie 2-algebra to be a Lie algebra internal to these 2-vector spaces.
- With strict skew symmetry but weak Jacobi identity.
- Notice one can also weaken the skew symmetry [Roytenberg].

The basic concepts

Lie 2-algebras

- Baez and Crans consider 2-vector spaces as categories internal to vector spaces.
- These are nothing but 2-term chain complexes – but interpreted suitably.
- They then define a Lie 2-algebra to be a Lie algebra internal to these 2-vector spaces.
 - With strict skew symmetry but weak Jacobi identity.
 - Notice one can also weaken the skew symmetry [Roytenberg].

The basic concepts

Lie 2-algebras

- Baez and Crans consider 2-vector spaces as categories internal to vector spaces.
- These are nothing but 2-term chain complexes – but interpreted suitably.
- They then define a Lie 2-algebra to be a Lie algebra internal to these 2-vector spaces.
- With strict skew symmetry but weak Jacobi identity.
- Notice one can also weaken the skew symmetry [Roytenberg].

The basic concepts

Lie 2-algebras

- Baez and Crans consider 2-vector spaces as categories internal to vector spaces.
- These are nothing but 2-term chain complexes – but interpreted suitably.
- They then define a Lie 2-algebra to be a Lie algebra internal to these 2-vector spaces.
- With strict skew symmetry but weak Jacobi identity.
- Notice one can also weaken the skew symmetry [Roytenberg].

The basic concepts

Lie n -algebras, L_∞ -algebras and qDGCAs

- Baez and Crans showed that their Lie 2-algebras are equivalent to 2-term L_∞ -algebras.
- In general, semistrict Lie n -algebras should be equivalent to n -term L_∞ -algebras.
- n -Term L_∞ -algebras are, in turn, equivalent to quasi-free differential graded algebras (qDGCAs) concentrated in the first n -degrees.
- The L_∞ - and qDGCA-formulation are good for computations. The categorical Lie n -algebra formulation is helpful conceptually.

The basic concepts

Lie n -algebras, L_∞ -algebras and qDGCAs

- Baez and Crans showed that their Lie 2-algebras are equivalent to 2-term L_∞ -algebras.
- In general, semistrict Lie n -algebras should be equivalent to n -term L_∞ -algebras.
- n -Term L_∞ -algebras are, in turn, equivalent to quasi-free differential graded algebras (qDGCAs) concentrated in the first n -degrees.
- The L_∞ - and qDGCA-formulation are good for computations. The categorical Lie n -algebra formulation is helpful conceptually.

The basic concepts

Lie n -algebras, L_∞ -algebras and qDGCA's

- Baez and Crans showed that their Lie 2-algebras are equivalent to 2-term L_∞ -algebras.
- In general, semistrict Lie n -algebras should be equivalent to n -term L_∞ -algebras.
- n -Term L_∞ -algebras are, in turn, equivalent to quasi-free differential graded algebras (qDGCA's) concentrated in the first n -degrees.
- The L_∞ - and qDGCA-formulation are good for computations. The categorical Lie n -algebra formulation is helpful conceptually.

The basic concepts

Lie n -algebras, L_∞ -algebras and qDGCA's

- Baez and Crans showed that their Lie 2-algebras are equivalent to 2-term L_∞ -algebras.
- In general, semistrict Lie n -algebras should be equivalent to n -term L_∞ -algebras.
- n -Term L_∞ -algebras are, in turn, equivalent to quasi-free differential graded algebras (qDGCA's) concentrated in the first n -degrees.
- The L_∞ - and qDGCA-formulation are good for computations. The categorical Lie n -algebra formulation is helpful conceptually.

The basic concepts

Lie n -algebras, L_∞ -algebras and qDGCA's

- Baez and Crans showed that their Lie 2-algebras are equivalent to 2-term L_∞ -algebras.
- In general, semistrict Lie n -algebras should be equivalent to n -term L_∞ -algebras.
- n -Term L_∞ -algebras are, in turn, equivalent to quasi-free differential graded algebras (qDGCA's) concentrated in the first n -degrees.
- The L_∞ - and qDGCA-formulation are good for computations. The categorical Lie n -algebra formulation is helpful conceptually.

Example: ordinary Lie algebras

For \mathfrak{g} an ordinary Lie algebra, and with $\{t_a\}$ a chosen basis with structure constants $\{C^a_{bc}\}$, the corresponding qDGCA is the graded-commutative exterior algebra

$$\wedge^\bullet(\mathfrak{sg}^*)$$

with \mathfrak{sg}^* denoting \mathfrak{g}^* in degree 1, on which

$$dt^a = -\frac{1}{2}C^a_{bc}t^b \wedge t^c$$

defines the differential.

It is a (sad) fact of life that qDGCA discussions are often unfeasible without choosing a basis

Example: ordinary Lie algebras

For \mathfrak{g} an ordinary Lie algebra, and with $\{t_a\}$ a chosen basis with structure constants $\{C^a_{bc}\}$, the corresponding qDGCA is the graded-commutative exterior algebra

$$\wedge^\bullet(\mathfrak{sg}^*)$$

with \mathfrak{sg}^* denoting \mathfrak{g}^* in degree 1, on which

$$dt^a = -\frac{1}{2}C^a_{bc}t^b \wedge t^c$$

defines the differential.

It is a (sad) fact of life that qDGCA discussions are often unfeasible without choosing a basis

Example: ordinary Lie algebras

For \mathfrak{g} an ordinary Lie algebra, and with $\{t_a\}$ a chosen basis with structure constants $\{C^a_{bc}\}$, the corresponding qDGCA is the graded-commutative exterior algebra

$$\wedge^\bullet(\mathfrak{sg}^*)$$

with \mathfrak{sg}^* denoting \mathfrak{g}^* in degree 1, on which

$$dt^a = -\frac{1}{2}C^a_{bc}t^b \wedge t^c$$

defines the differential.

It is a (sad) fact of life that qDGCA discussions are often
unfeasible without choosing a basis

Example: ordinary Lie algebras

For \mathfrak{g} an ordinary Lie algebra, and with $\{t_a\}$ a chosen basis with structure constants $\{C^a_{bc}\}$, the corresponding qDGCA is the graded-commutative exterior algebra

$$\wedge^\bullet(\mathfrak{sg}^*)$$

with \mathfrak{sg}^* denoting \mathfrak{g}^* in degree 1, on which

$$dt^a = -\frac{1}{2} C^a_{bc} t^b \wedge t^c$$

defines the differential.

It is a (sad) fact of life that qDGCA discussions are often unfeasible without choosing a basis

Example: ordinary Lie algebras

For \mathfrak{g} an ordinary Lie algebra, and with $\{t_a\}$ a chosen basis with structure constants $\{C^a_{bc}\}$, the corresponding qDGCA is the graded-commutative exterior algebra

$$\wedge^\bullet(\mathfrak{sg}^*)$$

with \mathfrak{sg}^* denoting \mathfrak{g}^* in degree 1, on which

$$dt^a = -\frac{1}{2} C^a_{bc} t^b \wedge t^c$$

defines the differential.

It is a (sad) fact of life that qDGCA discussions are often unfeasible without choosing a basis

Example: ordinary Lie algebras

For \mathfrak{g} an ordinary Lie algebra, and with $\{t_a\}$ a chosen basis with structure constants $\{C^a_{bc}\}$, the corresponding qDGCA is the graded-commutative exterior algebra

$$\wedge^\bullet(\mathfrak{sg}^*)$$

with \mathfrak{sg}^* denoting \mathfrak{g}^* in degree 1, on which

$$dt^a = -\frac{1}{2}C^a_{bc}t^b \wedge t^c$$

defines the differential.

It is a (sad) fact of life that qDGCA discussions are often unfeasible without choosing a basis

Example: ordinary Lie algebras

For \mathfrak{g} an ordinary Lie algebra, and with $\{t_a\}$ a chosen basis with structure constants $\{C^a_{bc}\}$, the corresponding qDGCA is the graded-commutative exterior algebra

$$\wedge^\bullet(\mathfrak{sg}^*)$$

with \mathfrak{sg}^* denoting \mathfrak{g}^* in degree 1, on which

$$dt^a = -\frac{1}{2}C^a_{bc}t^b \wedge t^c$$

defines the differential.

It is a (sad) fact of life that qDGCA discussions are often unfeasible without choosing a basis

The $\text{inn}(\cdot)$ -construction

Definition. (Inner derivation Lie $(n+1)$ -algebra)

$\text{inn}(\mathfrak{g}_{(n)})$ is the mapping cone of the identity on $\mathfrak{g}_{(n)}$

$$\text{inn}(\mathfrak{g}_{(n)}) \simeq (\wedge(\mathfrak{sg}_{(n)}) \oplus \mathfrak{ssg}_{(n)}), d')$$

Proposition

- There is a canonical injection $\mathfrak{g}_{(n)} \hookrightarrow \text{inn}(\mathfrak{g})$.
- $\text{inn}(\mathfrak{g}_{(n)})$ is *contractible*
- $(\wedge(\mathfrak{sg}_{(1)}) \oplus \mathfrak{ssg}_{(1)}), d')$ is the *Weil algebra* of $\mathfrak{g}_{(1)}$

Remark.

Hence $\text{inn}(\mathfrak{g}_{(1)})^*$ plays the role of differential forms on the universal G -bundle.

The $\text{inn}(\cdot)$ -construction

Definition. (Inner derivation Lie $(n+1)$ -algebra)

$\text{inn}(\mathfrak{g}_{(n)})$ is the mapping cone of the identity on $\mathfrak{g}_{(n)}$

$$\text{inn}(\mathfrak{g}_{(n)}) \simeq (\wedge(\mathfrak{sg}_{(n)}) \oplus \mathfrak{ssg}_{(n)}), d')$$

Proposition

- There is a canonical injection $\mathfrak{g}_{(n)} \hookrightarrow \text{inn}(\mathfrak{g})$.
- $\text{inn}(\mathfrak{g}_{(n)})$ is *contractible*
- $(\wedge(\mathfrak{sg}_{(1)}) \oplus \mathfrak{ssg}_{(1)}), d')$ is the *Weil algebra* of $\mathfrak{g}_{(1)}$

Remark.

Hence $\text{inn}(\mathfrak{g}_{(1)})^*$ plays the role of differential forms on the universal G -bundle.

The $\text{inn}(\cdot)$ -construction

Definition. (Inner derivation Lie $(n+1)$ -algebra)

$\text{inn}(\mathfrak{g}_{(n)})$ is the mapping cone of the identity on $\mathfrak{g}_{(n)}$

$$\text{inn}(\mathfrak{g}_{(n)}) \simeq (\wedge(\mathfrak{sg}_{(n)}) \oplus \mathfrak{ssg}_{(n)}), d')$$

Proposition

- There is a canonical injection $\mathfrak{g}_{(n)} \hookrightarrow \text{inn}(\mathfrak{g})$.
- $\text{inn}(\mathfrak{g}_{(n)})$ is *contractible*
- $(\wedge(\mathfrak{sg}_{(1)}) \oplus \mathfrak{ssg}_{(1)}), d')$ is the *Weil algebra* of $\mathfrak{g}_{(1)}$

Remark.

Hence $\text{inn}(\mathfrak{g}_{(1)})^*$ plays the role of differential forms on the universal G -bundle.

The $\text{inn}(\cdot)$ -construction

Definition. (Inner derivation Lie $(n+1)$ -algebra)

$\text{inn}(\mathfrak{g}_{(n)})$ is the mapping cone of the identity on $\mathfrak{g}_{(n)}$

$$\text{inn}(\mathfrak{g}_{(n)}) \simeq (\wedge(\mathfrak{sg}_{(n)}) \oplus \mathfrak{ssg}_{(n)}), d')$$

Proposition

- There is a canonical injection $\mathfrak{g}_{(n)} \hookrightarrow \text{inn}(\mathfrak{g})$.
- $\text{inn}(\mathfrak{g}_{(n)})$ is *contractible*
- $(\wedge(\mathfrak{sg}_{(1)}) \oplus \mathfrak{ssg}_{(1)}), d')$ is the *Weil algebra* of $\mathfrak{g}_{(1)}$

Remark.

Hence $\text{inn}(\mathfrak{g}_{(1)})^*$ plays the role of differential forms on the universal G -bundle.

The $\text{inn}(\cdot)$ -construction

Definition. (Inner derivation Lie $(n+1)$ -algebra)

$\text{inn}(\mathfrak{g}_{(n)})$ is the mapping cone of the identity on $\mathfrak{g}_{(n)}$

$$\text{inn}(\mathfrak{g}_{(n)}) \simeq (\wedge(\mathfrak{sg}_{(n)}) \oplus \mathfrak{ssg}_{(n)}), d')$$

Proposition

- There is a canonical injection $\mathfrak{g}_{(n)} \hookrightarrow \text{inn}(\mathfrak{g})$.
- $\text{inn}(\mathfrak{g}_{(n)})$ is *contractible*
- $(\wedge(\mathfrak{sg}_{(1)}) \oplus \mathfrak{ssg}_{(1)}), d')$ is the *Weil algebra* of $\mathfrak{g}_{(1)}$

Remark.

Hence $\text{inn}(\mathfrak{g}_{(1)})^*$ plays the role of differential forms on the universal G -bundle.

The $\text{inn}(\cdot)$ -construction

The qDGCA of $\text{inn}(\mathfrak{g})$: the Weil algebra

$\text{inn}(\mathfrak{g}) \simeq (\bigwedge^\bullet (s\mathfrak{g}^* \oplus ss\mathfrak{g}^*), d)$ is spanned by generators $\{t^a\}$ in degree 1 and $\{r^a\}$ in degree 2, with differential

$$\begin{aligned} dt^a &= -\frac{1}{2} C^a_{bc} t^b \wedge t^c - r^a \\ dr^a &= -C^a_{bc} t^b \wedge r^c. \end{aligned}$$

We will now

- express the Lie algebra cohomology of \mathfrak{g} in terms of the cohomology of the qDGCA underlying $\text{inn}(\mathfrak{g})$.
- use the insight gained thereby to describe three families of Lie n -algebras: one for each cocycle, one for each invariant polynomial and one for each transgression element.
- then show that for the canonical 3-cocycle on a semisimple Lie algebra, connections with values in the Lie 3-algebra obtained this way describe the Chern-Simons parallel transport which we are after.

We will now

- express the Lie algebra cohomology of \mathfrak{g} in terms of the cohomology of the qDGCA underlying $\text{inn}(\mathfrak{g})$.
- use the insight gained thereby to describe three families of Lie n -algebras: one for each cocycle, one for each invariant polynomial and one for each transgression element.
- then show that for the canonical 3-cocycle on a semisimple Lie algebra, connections with values in the Lie 3-algebra obtained this way describe the Chern-Simons parallel transport which we are after.

We will now

- express the Lie algebra cohomology of \mathfrak{g} in terms of the cohomology of the qDGCA underlying $\text{inn}(\mathfrak{g})$.
- use the insight gained thereby to describe three families of Lie n -algebras: one for each cocycle, one for each invariant polynomial and one for each transgression element.
- then show that for the canonical 3-cocycle on a semisimple Lie algebra, connections with values in the Lie 3-algebra obtained this way describe the Chern-Simons parallel transport which we are after.

We will now

- express the Lie algebra cohomology of \mathfrak{g} in terms of the cohomology of the qDGCA underlying $\text{inn}(\mathfrak{g})$.
- use the insight gained thereby to describe three families of Lie n -algebras: one for each cocycle, one for each invariant polynomial and one for each transgression element.
- then show that for the canonical 3-cocycle on a semisimple Lie algebra, connections with values in the Lie 3-algebra obtained this way describe the Chern-Simons parallel transport which we are after.

Lie algebra cohomology in terms of $\text{inn}(\mathfrak{g})$

- A Lie algebra n -cocycle μ is

$$d|_{\wedge^{\bullet}(\mathfrak{sg}^*)}\mu = 0.$$

- An invariant degree n -polynomial k is

$$d|_{\wedge^{\bullet}(\mathfrak{ssg}^*)}k = 0.$$

- A transgression element cs is

$$\begin{aligned} \text{cs}|_{\wedge^{\bullet}\mathfrak{sg}^*} &= \mu \\ d\text{cs} &= k. \end{aligned}$$

Lie algebra cohomology in terms of $\text{inn}(\mathfrak{g})$

- A Lie algebra n -cocycle μ is

$$d|_{\wedge^{\bullet}(\mathfrak{sg}^*)}\mu = 0.$$

- An invariant degree n -polynomial k is

$$d|_{\wedge^{\bullet}(\mathfrak{ssg}^*)}k = 0.$$

- A transgression element cs is

$$\begin{aligned} cs|_{\wedge^{\bullet}\mathfrak{sg}^*} &= \mu \\ dcs &= k. \end{aligned}$$

Lie algebra cohomology in terms of $\text{inn}(\mathfrak{g})$

- A Lie algebra n -cocycle μ is

$$d| \wedge^{\bullet}_{(s\mathfrak{g}^*)} \mu = 0.$$

- An invariant degree n -polynomial k is

$$d| \wedge^{\bullet}_{(ss\mathfrak{g}^*)} k = 0.$$

- A transgression element cs is

$$\begin{aligned} cs| \wedge^{\bullet}_{s\mathfrak{g}^*} &= \mu \\ dcs &= k. \end{aligned}$$

Lie algebra cohomology in terms of $\text{inn}(\mathfrak{g})$

- A Lie algebra n -cocycle μ is

$$d| \wedge^{\bullet}_{(s\mathfrak{g}^*)} \mu = 0.$$

- An invariant degree n -polynomial k is

$$d| \wedge^{\bullet}_{(ss\mathfrak{g}^*)} k = 0.$$

- A transgression element cs is

$$\begin{aligned} cs| \wedge^{\bullet}_{s\mathfrak{g}^*} &= \mu \\ dcs &= k. \end{aligned}$$

Lie algebra cohomology in terms of $\text{inn}(\mathfrak{g})$

- A Lie algebra n -cocycle μ is

$$d| \wedge^{\bullet}_{(s\mathfrak{g}^*)} \mu = 0.$$

- An invariant degree n -polynomial k is

$$d| \wedge^{\bullet}_{(ss\mathfrak{g}^*)} k = 0.$$

- A **transgression element** cs is

$$\begin{aligned} \text{cs}| \wedge^{\bullet}_{s\mathfrak{g}^*} &= \mu \\ d\text{cs} &= k. \end{aligned}$$

Lie algebra cohomology in terms of $\text{inn}(\mathfrak{g})$

- A Lie algebra n -cocycle μ is

$$d| \bigwedge^{\bullet}_{(s\mathfrak{g}^*)} \mu = 0.$$

- An invariant degree n -polynomial k is

$$d| \bigwedge^{\bullet}_{(ss\mathfrak{g}^*)} k = 0.$$

- A transgression element cs is

$$\begin{aligned} \text{cs}| \bigwedge^{\bullet}_{s\mathfrak{g}^*} &= \mu \\ d\text{cs} &= k. \end{aligned}$$

The homotopy operator

- Recall that we said that $\text{inn}(\mathfrak{g}_{(n)})$ is trivializable.
- This means there is a homotopy

$$\begin{array}{ccc}
 & 0 & \\
 \text{inn}(\mathfrak{g}_{(n)}) & \begin{array}{c} \xrightarrow{\quad} \\ \Downarrow \tau \\ \xrightarrow{\quad} \end{array} & \text{inn}(\mathfrak{g}_{(n)}) \\
 & \text{Id}=[d,\tau] &
 \end{array}$$

- Since we have τ , we have an effective algorithm to always solve $k = dcs$ as

$$cs := \tau(k) + dq.$$

- The only nontrivial condition is hence $cs|_{\bigwedge^{\bullet} s\mathfrak{g}^*} = \mu.$

The homotopy operator

- Recall that we said that $\text{inn}(\mathfrak{g}_{(n)})$ is trivializable.
- This means there is a homotopy

$$\begin{array}{ccc}
 & 0 & \\
 \text{inn}(\mathfrak{g}_{(n)}) & \begin{array}{c} \xrightarrow{\quad} \\ \Downarrow \tau \\ \xrightarrow{\quad} \end{array} & \text{inn}(\mathfrak{g}_{(n)}) \\
 & \text{Id}=[d,\tau] &
 \end{array}$$

- Since we have τ , we have an effective algorithm to always solve $k = dcs$ as

$$cs := \tau(k) + dq.$$

- The only nontrivial condition is hence $cs|_{\bigwedge^{\bullet} \mathfrak{sg}^*} = \mu.$

The homotopy operator

- Recall that we said that $\text{inn}(\mathfrak{g}_{(n)})$ is trivializable.
- This means there is a homotopy

$$\begin{array}{ccc}
 & 0 & \\
 \text{inn}(\mathfrak{g}_{(n)}) & \begin{array}{c} \xrightarrow{\quad} \\ \Downarrow \tau \\ \xrightarrow{\quad} \end{array} & \text{inn}(\mathfrak{g}_{(n)}) \\
 & \text{Id}=[d,\tau] &
 \end{array}$$

- Since we have τ , we have an effective algorithm to always solve $k = d_{\text{CS}}$ as

$$\text{CS} := \tau(k) + dq.$$

- The only nontrivial condition is hence $\text{CS}|_{\bigwedge^{\bullet} \mathfrak{sg}^*} = \mu.$

The homotopy operator

- Recall that we said that $\text{inn}(\mathfrak{g}_{(n)})$ is trivializable.
- This means there is a homotopy

$$\begin{array}{ccc}
 & 0 & \\
 \text{inn}(\mathfrak{g}_{(n)}) & \begin{array}{c} \xrightarrow{\quad} \\ \Downarrow \tau \\ \xrightarrow{\quad} \end{array} & \text{inn}(\mathfrak{g}_{(n)}) \\
 & \text{Id}=[d, \tau] &
 \end{array}$$

- Since we have τ , we have an effective algorithm to always solve $k = d_{\text{CS}}$ as

$$_{\text{CS}} := \tau(k) + dq.$$

- The only nontrivial condition is hence $_{\text{CS}}|_{\bigwedge^{\bullet} \mathfrak{sg}^*} = \mu.$

A map of the cocycle situation

cocycle

Chern-Simons

inv. polynomial

$$(\wedge^\bullet(\mathfrak{sg}^*), d_{\mathfrak{g}}) \xleftarrow{i^*} (\wedge^\bullet(\mathfrak{sg}^* \oplus \mathfrak{ssg}^*), d_{\text{inn}(\mathfrak{g})}) \xleftarrow{p^*} (\wedge^\bullet(\mathfrak{ssg}^*)^*)$$

$$\begin{array}{ccccc}
 & & 0 & & \\
 & & \uparrow d_{\text{inn}(\mathfrak{g})} & & \\
 & & \text{---} & & \\
 & & p^* k & \xleftarrow{p^*} & k \\
 & & \uparrow d_{\text{inn}(\mathfrak{g})} & & \\
 \tau \swarrow & & & & \\
 & & \text{---} & & \\
 & & CS & & \\
 \mu \xleftarrow{i^*} & & & &
 \end{array}$$

Diagram illustrating the map of the cocycle situation. The diagram shows a commutative structure involving spaces of cocycles, Chern-Simons forms, and invariant polynomials. The top row shows the map i^* from $(\wedge^\bullet(\mathfrak{sg}^*), d_{\mathfrak{g}})$ to $(\wedge^\bullet(\mathfrak{sg}^* \oplus \mathfrak{ssg}^*), d_{\text{inn}(\mathfrak{g})})$, and the map p^* from $(\wedge^\bullet(\mathfrak{sg}^* \oplus \mathfrak{ssg}^*), d_{\text{inn}(\mathfrak{g})})$ to $(\wedge^\bullet(\mathfrak{ssg}^*)^*)$. The bottom part shows a diagram with nodes 0 , $p^* k$, k , CS , and μ , connected by arrows $d_{\text{inn}(\mathfrak{g})}$, p^* , i^* , and τ .

Baez-Crans Lie n -algebras from cocycles

Definition and proposition [Baez,Crans]

For every Lie algebra $(n+1)$ -cocycle μ of the Lie algebra \mathfrak{g} there is a skeletal Lie n -algebra

$$\mathfrak{g}_\mu.$$

Construction.

Set $\mathfrak{g}_\mu \simeq (\wedge^*(s\mathfrak{g}^* \oplus s^n \mathbb{R}^*), d)$ such that the differential is given by

$$\begin{aligned} dt^a &= -\frac{1}{2} C^a_{bc} t^b \wedge t^c \\ db &= -\mu \end{aligned}$$

Baez-Crans Lie n -algebras from cocycles

Definition and proposition [Baez,Crans]

For every Lie algebra $(n+1)$ -cocycle μ of the Lie algebra \mathfrak{g} there is a skeletal Lie n -algebra

$$\mathfrak{g}_\mu.$$

Construction.

Set $\mathfrak{g}_\mu \simeq (\wedge^*(s\mathfrak{g}^* \oplus s^n \mathbb{R}^*), d)$ such that the differential is given by

$$\begin{aligned} dt^a &= -\frac{1}{2} C^a_{bc} t^b \wedge t^c \\ db &= -\mu \end{aligned}$$

Baez-Crans Lie n -algebras from cocycles

Definition and proposition [Baez,Crans]

For every Lie algebra $(n+1)$ -cocycle μ of the Lie algebra \mathfrak{g} there is a skeletal Lie n -algebra

$$\mathfrak{g}_\mu.$$

Construction.

Set $\mathfrak{g}_\mu \simeq (\wedge^\bullet(s\mathfrak{g}^* \oplus s^n\mathbb{R}^*), d)$ such that the differential is given by

$$\begin{aligned} dt^a &= -\frac{1}{2} C^a_{bc} t^b \wedge t^c \\ db &= -\mu \end{aligned}$$

Baez-Crans Lie n -algebras from cocycles

Definition and proposition [Baez,Crans]

For every Lie algebra $(n+1)$ -cocycle μ of the Lie algebra \mathfrak{g} there is a skeletal Lie n -algebra

$$\mathfrak{g}_\mu.$$

Construction.

Set $\mathfrak{g}_\mu \simeq (\wedge^\bullet(\mathfrak{sg}^* \oplus s^n \mathbb{R}^*), d)$ such that with $\{t^a\}$ as basis for \mathfrak{sg}^* and with $\{b\}$ the canonical basis of \mathbb{R} the differential is given by

$$\begin{aligned} dt^a &= -\frac{1}{2} C^a_{bc} t^b \wedge t^c \\ db &= -\mu \end{aligned}$$

Baez-Crans Lie n -algebras from cocycles

Definition and proposition [Baez,Crans]

For every Lie algebra $(n+1)$ -cocycle μ of the Lie algebra \mathfrak{g} there is a skeletal Lie n -algebra

$$\mathfrak{g}_\mu.$$

Construction.

Set $\mathfrak{g}_\mu \simeq (\wedge^\bullet(s\mathfrak{g}^* \oplus s^n\mathbb{R}^*), d)$ such that the differential is given by

$$\begin{aligned} dt^a &= -\frac{1}{2} C^a_{bc} t^b \wedge t^c \\ db &= -\mu \end{aligned}$$

Chern Lie n -algebras from invariant polynomials

Definition and proposition

For every degree $(n + 1)$ Lie algebra invariant polynomial k of the Lie algebra \mathfrak{g} there is a Lie $(2n + 1)$ -algebra

$$\text{ch}_k(\mathfrak{g}).$$

Construction.

Set $\text{ch}_k(\mathfrak{g}) \simeq (\wedge^*(s\mathfrak{g}^* \oplus s\mathfrak{g}^* \oplus s^{(2n+1)}\mathbb{R}^*), d)$ such that we have

$$\begin{aligned} dt^a &= -\frac{1}{2} C^a_{bc} t^b \wedge t^c - r^a \\ dr^a &= -C^a_{bc} t^b \wedge t^c \\ dc &= k \end{aligned}$$

Chern Lie n -algebras from invariant polynomials

Definition and proposition

For every degree $(n + 1)$ Lie algebra invariant polynomial k of the Lie algebra \mathfrak{g} there is a Lie $(2n + 1)$ -algebra

$$\text{ch}_k(\mathfrak{g}).$$

Construction.

Set $\text{ch}_k(\mathfrak{g}) \simeq (\wedge^*(s\mathfrak{g}^* \oplus s\mathfrak{g}^* \oplus s^{(2n+1)}\mathbb{R}^*), d)$ such that we have

$$\begin{aligned} dt^a &= -\frac{1}{2} C^a_{bc} t^b \wedge t^c - r^a \\ dr^a &= -C^a_{bc} t^b \wedge t^c \\ dc &= k \end{aligned}$$

Chern Lie n -algebras from invariant polynomials

Definition and proposition

For every degree $(n+1)$ Lie algebra invariant polynomial k of the Lie algebra \mathfrak{g} there is a Lie $(2n+1)$ -algebra

$$\text{ch}_k(\mathfrak{g}).$$

Construction.

Set $\text{ch}_k(\mathfrak{g}) \simeq (\wedge^\bullet(\mathfrak{sg}^* \oplus \mathfrak{ssg}^* \oplus s^{(2n+1)}\mathbb{R}^*), d)$ such that we have

$$\begin{aligned} dt^a &= -\frac{1}{2} C^a_{bc} t^b \wedge t^c - r^a \\ dr^a &= -C^a_{bc} t^b \wedge t^c \\ dc &= k \end{aligned}$$

Chern Lie n -algebras from invariant polynomials

Definition and proposition

For every degree $(n + 1)$ Lie algebra invariant polynomial k of the Lie algebra \mathfrak{g} there is a Lie $(2n + 1)$ -algebra

$$\text{ch}_k(\mathfrak{g}).$$

Construction.

Set $\text{ch}_k(\mathfrak{g}) \simeq (\bigwedge^\bullet(\mathfrak{sg}^* \oplus \text{ssg}^* \oplus s^{(2n+1)}\mathbb{R}^*), d)$ such that with $\{t^a\}$ a basis for \mathfrak{sg}^* , $\{r^a\}$ the corresponding basis for ssg^* and with $\{c\}$ the canonical basis of $s^{(2n+1)}$ we have

$$\begin{aligned} dt^a &= -\frac{1}{2} C^a_{bc} t^b \wedge t^c - r^a \\ dr^a &= -C^a_{bc} t^b \wedge t^c \end{aligned}$$

Chern Lie n -algebras from invariant polynomials

Definition and proposition

For every degree $(n + 1)$ Lie algebra invariant polynomial k of the Lie algebra \mathfrak{g} there is a Lie $(2n + 1)$ -algebra

$$\text{ch}_k(\mathfrak{g}).$$

Construction.

Set $\text{ch}_k(\mathfrak{g}) \simeq (\wedge^\bullet(\mathfrak{sg}^* \oplus \mathfrak{s} \mathfrak{sg}^* \oplus s^{(2n+1)}\mathbb{R}^*), d)$ such that we have

$$\begin{aligned} dt^a &= -\frac{1}{2} C^a_{bc} t^b \wedge t^c - r^a \\ dr^a &= -C^a_{bc} t^b \wedge t^c \\ dc &= k \end{aligned}$$

Chern-Simons Lie n -algebras from transgression elements

Definition and proposition

For every transgression element q of degree $(2n + 1)$ there is a Lie $(2n + 1)$ -algebra

$$\text{cs}_k(\mathfrak{g}) .$$

Construction.

Set $\text{cs}_k(\mathfrak{g}) \simeq (\wedge^*(s\mathfrak{g}^* \oplus s s\mathfrak{g}^* \oplus \oplus s^{2n}\mathbb{R}^* \oplus s^{(2n+1)}\mathbb{R}^*), d)$ such that

$$dt^a = -\frac{1}{2} C^a_{bc} t^b \wedge t^c - r^a$$

$$dr^a = -C^a_{bc} t^b \wedge t^c$$

$$db = -\text{cs} + c$$

$$dc = k$$

Chern-Simons Lie n -algebras from transgression elements

Definition and proposition

For every transgression element q of degree $(2n + 1)$ there is a Lie $(2n + 1)$ -algebra

$$\text{CS}_k(\mathfrak{g}) .$$

Construction.

Set $\text{CS}_k(\mathfrak{g}) \simeq (\wedge^*(s\mathfrak{g}^* \oplus s s\mathfrak{g}^* \oplus \oplus s^{2n}\mathbb{R}^* \oplus s^{(2n+1)}\mathbb{R}^*), d)$ such that

$$dt^a = -\frac{1}{2} C^a_{bc} t^b \wedge t^c - r^a$$

$$dr^a = -C^a_{bc} t^b \wedge t^c$$

$$db = -\text{CS} + c$$

$$dc = k$$

Chern-Simons Lie n -algebras from transgression elements

Definition and proposition

For every transgression element q of degree $(2n + 1)$ there is a Lie $(2n + 1)$ -algebra

$$\text{cs}_k(\mathfrak{g}) .$$

Construction.

Set $\text{cs}_k(\mathfrak{g}) \simeq (\bigwedge^\bullet(s\mathfrak{g}^* \oplus s s\mathfrak{g}^* \oplus \oplus s^{2n}\mathbb{R}^* \oplus s^{(2n+1)}\mathbb{R}^*), d)$ such that

$$dt^a = -\frac{1}{2} C^a_{bc} t^b \wedge t^c - r^a$$

$$dr^a = -C^a_{bc} t^b \wedge t^c$$

$$db = -\text{cs} + c$$

$$dc = k$$

Chern-Simons Lie n -algebras from transgression elements

Definition and proposition

For every transgression element q of degree $(2n + 1)$ there is a Lie $(2n + 1)$ -algebra

$$\text{cs}_k(\mathfrak{g}) .$$

Construction.

Set $\text{cs}_k(\mathfrak{g}) \simeq (\bigwedge^\bullet(s\mathfrak{g}^* \oplus s s\mathfrak{g}^* \oplus \oplus s^{2n}\mathbb{R}^* \oplus s^{(2n+1)}\mathbb{R}^*), d)$ such that

$$dt^a = -\frac{1}{2}C^a_{bc}t^b \wedge t^c - r^a$$

$$dr^a = -C^a_{bc}t^b \wedge t^c$$

$$db = -\text{cs} + c$$

$$dc = k$$

Theorem

Whenever they exist, these Lie $(2n + 1)$ -algebras form a (weakly) short exact sequence:

$$0 \rightarrow \mathfrak{g}_{\mu_k} \rightarrow \mathrm{cs}_k(\mathfrak{g}) \rightarrow \mathrm{ch}_k(\mathfrak{g}) \rightarrow 0.$$

Theorem

Moreover, we have an isomorphism

$$\mathrm{cs}_k(\mathfrak{g}) \simeq \mathrm{inn}(\mathfrak{g}_{\mu_k}).$$

Theorem

Whenever they exist, these Lie $(2n + 1)$ -algebras form a (weakly) short exact sequence:

$$0 \rightarrow \mathfrak{g}_{\mu_k} \rightarrow \mathrm{cs}_k(\mathfrak{g}) \rightarrow \mathrm{ch}_k(\mathfrak{g}) \rightarrow 0.$$

Theorem

Moreover, we have an isomorphism

$$\mathrm{cs}_k(\mathfrak{g}) \simeq \mathrm{inn}(\mathfrak{g}_{\mu_k}).$$

Now we can study n -connections with values in these Lie n -algebras.

Definition

For our purposes, an n -connection with values in the Lie n -algebra $\mathfrak{g}_{(n)}$ on a space X is a morphism

$$A : \text{Vect}(X) \rightarrow \mathfrak{g}_{(n)},$$

which we conceive as the dual morphism

$$A^* : \mathfrak{g}_{(n)}^* \rightarrow \Omega^\bullet(X)$$

of differential algebras.

Now we can study n -connections with values in these Lie n -algebras.

Definition

For our purposes, an n -connection with values in the Lie n -algebra $\mathfrak{g}_{(n)}$ on a space X is a morphism

$$A : \text{Vect}(X) \rightarrow \mathfrak{g}_{(n)} ,$$

which we conceive as the dual morphism

$$A^* : \mathfrak{g}_{(n)}^* \rightarrow \Omega^\bullet(X)$$

of differential algebras.

Ordinary connection 1-forms

Ordinary connection 1-forms

$$n=1$$

$$\begin{array}{c} \mathfrak{g} \\ (A) \uparrow \\ F_A=0 \\ \text{Vect}(X) \end{array}$$

for $A \in \Omega^1(X, \mathfrak{g})$.

Morphisms into $\mathfrak{g}_{(1)}$ come from *flat* connection 1-forms.

Ordinary connection 1-forms

Ordinary connection 1-forms

$$\begin{array}{ccc}
 n=1 & & n=2 \\
 \\
 \begin{array}{c} \mathfrak{g} \\ \uparrow (A) \\ \text{Vect}(X) \end{array} & \xrightarrow{\quad} & \begin{array}{c} \text{inn}(\mathfrak{g}) \\ \uparrow (A) \\ \text{Vect}(X) \end{array} \\
 F_A=0 & &
 \end{array}$$

for $A \in \Omega^1(X, \mathfrak{g})$.

Morphisms into $\text{inn}(\mathfrak{g}_{(1)})$ come from *arbitrary* connection 1-forms.

General Chern-Simons-like connections

Theorem

For every degree $(2n + 1)$ Lie algebra transgressive element, $(2n + 1)$ -connections with values in $cs_k(\mathfrak{g})$ are in bijection with \mathfrak{g} -Chern-Simons forms.

This means...

General Chern-Simons-like connections

Theorem

For every degree $(2n + 1)$ Lie algebra transgressive element, $(2n + 1)$ -connections with values in $cs_k(\mathfrak{g})$ are in bijection with \mathfrak{g} -Chern-Simons forms.

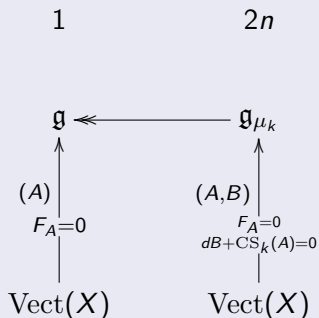
This means...

General Chern-Simons-like connections

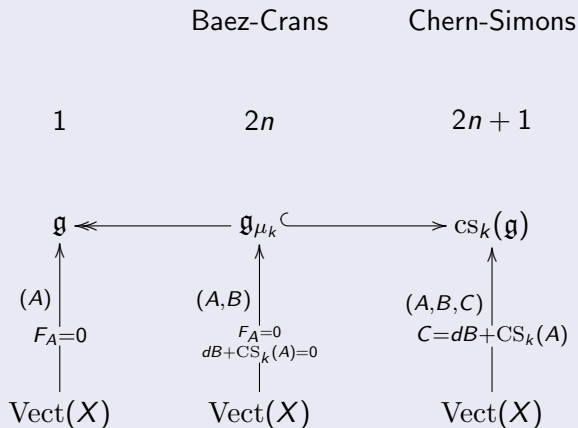
$$\begin{array}{c} 1 \\ \uparrow \\ \mathfrak{g} \\ \uparrow \\ (A) \mid \\ F_A=0 \\ \mid \\ \text{Vect}(X) \end{array}$$

General Chern-Simons-like connections

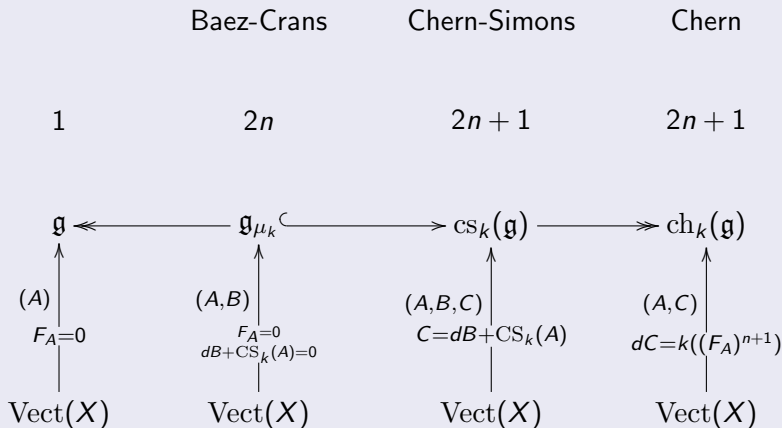
Baez-Crans



General Chern-Simons-like connections



General Chern-Simons-like connections



The standard Chern-Simons 3-connection

Finally: the case we wanted to understand

Let now \mathfrak{g} be semisimple and let

$$\mu = \langle \cdot, [\cdot, \cdot] \rangle$$

be the canonical 3-cocycle.

Theorem (Baez, Crans, S, Stevenson)

The corresponding Baez-Crans Lie 2-algebra \mathfrak{g}_μ is equivalent to that of the corresponding String 2-group

$$\mathfrak{g}_\mu \simeq \text{Lie}(\text{String}_k(G)).$$

The standard Chern-Simons 3-connection

Finally: the case we wanted to understand

Let now \mathfrak{g} be semisimple and let

$$\mu = \langle \cdot, [\cdot, \cdot] \rangle$$

be the canonical 3-cocycle.

Theorem (Baez, Crans, S, Stevenson)

The corresponding Baez-Crans Lie 2-algebra \mathfrak{g}_μ is equivalent to that of the corresponding String 2-group

$$\mathfrak{g}_\mu \simeq \text{Lie}(\text{String}_k(G)).$$

The standard Chern-Simons 3-connection

$$\begin{array}{c} \mathfrak{g} \\ \parallel \\ \mathfrak{g} \\ \uparrow (A) \\ F_A=0 \\ \text{Vect}(X) \end{array}$$

The standard Chern-Simons 3-connection

$$\begin{array}{ccc}
 \mathfrak{g} & \xleftarrow{\quad} & \text{string}_k(\mathfrak{g}) \\
 \parallel & & \parallel \sim \\
 \mathfrak{g} & \xleftarrow{\quad} & \mathfrak{g}_k \\
 (A) \uparrow & & (A,B) \uparrow \\
 F_A=0 & & F_A=0 \\
 \text{Vect}(X) & & dB+kCS(A)=0 \\
 & & \text{Vect}(X)
 \end{array}$$

The standard Chern-Simons 3-connection

$$\begin{array}{ccccc}
 \mathfrak{g} & \xleftarrow{\quad} & \text{string}_k(\mathfrak{g}) & \xhookrightarrow{\quad} & \text{inn}(\text{string}_k(\mathfrak{g})) \\
 \parallel & & \parallel \sim & & \parallel \sim \\
 \mathfrak{g} & \xleftarrow{\quad} & \mathfrak{g}_k & \xrightarrow{\quad} & \text{CS}_k(\mathfrak{g}) \\
 (A) \uparrow & & (A,B) \uparrow & & (A,B,C) \uparrow \\
 F_A=0 & & F_A=0 & & C=dB+k\text{CS}(A) \\
 \text{Vect}(X) & & \text{Vect}(X) & & \text{Vect}(X)
 \end{array}$$

$dB+k\text{CS}(A)=0$

The standard Chern-Simons 3-connection

$$\begin{array}{ccccccc}
 \mathfrak{g} & \xleftarrow{\quad} & \text{string}_k(\mathfrak{g}) & \hookrightarrow & \text{inn}(\text{string}_k(\mathfrak{g})) \\
 \parallel & & \parallel \sim & & \parallel \sim \\
 \mathfrak{g} & \xleftarrow{\quad} & \mathfrak{g}_k & \hookrightarrow & \text{CS}_k(\mathfrak{g}) & \twoheadrightarrow & \text{ch}_k(\mathfrak{g}) \\
 (A) \uparrow & & (A,B) \uparrow & & (A,B,C) \uparrow & & (A,C) \uparrow \\
 F_A=0 & & F_A=0 & & C=dB+k\text{CS}(A) & & dC=\langle F_A \wedge F_A \rangle \\
 \text{Vect}(X) & & \text{Vect}(X) & & \text{Vect}(X) & & \text{Vect}(X)
 \end{array}$$

$\text{Vect}(X)$ is the space of vector fields on X .
 $\text{Vect}(X)$ is the space of vector fields on X .
 $\text{Vect}(X)$ is the space of vector fields on X .
 $\text{Vect}(X)$ is the space of vector fields on X .

Conclusion

- 1 Every degree $(2n + 1)$ transgressive element in Lie algebra cohomology gives rise to an exact sequence

$$0 \rightarrow \mathfrak{g}_{\mu_k} \rightarrow \text{cs}_k(\mathfrak{g}) \rightarrow \text{ch}_k(\mathfrak{g}) \rightarrow 0$$

of Lie $(2n + 1)$ -algebras.

- 2 Connections with values in these Lie $(2n + 1)$ -algebras are precisely the corresponding Chern-Simons functionals, “localized” as a $(2n + 1)$ -transport.

Conclusion

- 1 Every degree $(2n + 1)$ transgressive element in Lie algebra cohomology gives rise to an exact sequence

$$0 \rightarrow \mathfrak{g}_{\mu_k} \rightarrow \text{cs}_k(\mathfrak{g}) \rightarrow \text{ch}_k(\mathfrak{g}) \rightarrow 0$$

of Lie $(2n + 1)$ -algebras.

- 2 Connections with values in these Lie $(2n + 1)$ -algebras are precisely the corresponding Chern-Simons functionals, “localized” as a $(2n + 1)$ -transport.

This is the end of the *talk*.

But for those reading this, here are a couple of further aspects and topics:

This is the end of the *talk*.

But for those reading this, here are a couple of further aspects and topics:

Further aspects and topics

- Which Lie 3-group does our Chern-Simons Lie 3-algebra $cs_k(\mathfrak{g})$ integrate to?
- How does the Chern-Simons 3-group help to understand the relation between Chern-Simons theory and Wess-Zumino-Witten theory?
- How does the Chern-Simons 3-transport help to understand String bundles with connection?
- How are the Chern Lie $(2n + 1)$ -algebras related to Roytenberg's Lie n -algebras with weak skew-symmetry?

The following slides contain brief comments on this.

Further aspects and topics

- Which Lie 3-group does our Chern-Simons Lie 3-algebra $cs_k(\mathfrak{g})$ integrate to?
- How does the Chern-Simons 3-group help to understand the relation between Chern-Simons theory and Wess-Zumino-Witten theory?
- How does the Chern-Simons 3-transport help to understand String bundles with connection?
- How are the Chern Lie $(2n + 1)$ -algebras related to Roytenberg's Lie n -algebras with weak skew-symmetry?

The following slides contain brief comments on this.

Further aspects and topics

- Which Lie 3-group does our Chern-Simons Lie 3-algebra $_{CS_k}(\mathfrak{g})$ integrate to?
- How does the Chern-Simons 3-group help to understand the relation between Chern-Simons theory and Wess-Zumino-Witten theory?
- How does the Chern-Simons 3-transport help to understand String bundles with connection?
- How are the Chern Lie $(2n + 1)$ -algebras related to Roytenberg's Lie n -algebras with weak skew-symmetry?

The following slides contain brief comments on this.

Further aspects and topics

- Which Lie 3-group does our Chern-Simons Lie 3-algebra $_{CS_k}(\mathfrak{g})$ integrate to?
- How does the Chern-Simons 3-group help to understand the relation between Chern-Simons theory and Wess-Zumino-Witten theory?
- How does the Chern-Simons 3-transport help to understand String bundles with connection?
- How are the Chern Lie $(2n + 1)$ -algebras related to Roytenberg's Lie n -algebras with weak skew-symmetry?

The following slides contain brief comments on this.

Integrating the Chern-Simons Lie 3-algebra

We have found the Lie 3-algebra

$$\mathrm{CS}_{\langle \cdot, \cdot \rangle}(\mathfrak{g})$$

which characterizes Chern-Simons 3-transp. Due to the isomorphism

$$\mathrm{CS}_{\langle \cdot, [\cdot, \cdot] \rangle} \simeq \mathrm{inn}(\mathfrak{g}_{\langle \cdot, [\cdot, \cdot] \rangle})$$

and the equivalence

$$\mathfrak{g}_{\langle \cdot, [\cdot, \cdot] \rangle} \simeq \mathrm{string}_{\langle \cdot, \cdot \rangle}(\mathfrak{g})$$

the corresponding Lie 3-group should be

$$\mathrm{INN}_0(\mathrm{String}_{\langle \cdot, \cdot \rangle})(G).$$

What does $\text{INN}_0(\text{String}_k(G))$ know about Wess-Zumino-Witten theory?

By the *n-functorial holographic principle* we expect that the WZW 2-transport arises as *transformations* of the Chern-Simons 3-transport.

Indeed, one can see that the 2-category of cylinders in the 3-group $\text{INN}_0(\text{String}_k(G))$ has the right properties for this.

(To see this, recall that, by [Baez,Crans,S,Schreiber], the 2-group $\text{String}_k(G)$ is, as a groupoid, precisely the canonical gerbe on G which enters the WZW theory.)

How does the Chern-Simons Lie 3-algebra help to understand String-bundles with connections?

Ask Google the question: "*How to get a spinning string from here to there?*"

Which role do weakly skew-symmetric Lie n -algebras play?

Baez and Crans considered Lie n -algebras which have a strictly skew-symmetric bracket, which however satisfies the Jacobi identity only weakly.

One can see that the Lie n -algebras \mathfrak{g}_μ for each $(n+1)$ -cocycle μ which we discussed are ordinary Lie algebras except that their $(n-2)$ nd coherence for the Jacobiator is nontrivial, and in fact precisely given by this cocycle.

Now, Roytenberg considered the case that skew-symmetry is weakened, too. It seems that one finds in his description that the higher coherences then are given by *symmetric* forms on the Lie algebra. In particular, for $n=2$ the Killing form may be regarded as the “skew-symmetrizer”.

This seems to suggest the following picture:

Which role do weakly skew-symmetric Lie n -algebras play?

This seems to suggest the following picture:

weak Jacobi

weak skew symmetry

$$\mathfrak{g}_{\mu_k} \longrightarrow \text{cs}_{\mu}(\mathfrak{g}) \longrightarrow \text{ch}_{\mu}(\mathfrak{g})$$

and might thus give an even more fundamental interpretation of the exact sequence which we found.

But this needs to be better understood.

More details

On the following slides are collected a couple of definitions and explanations which were omitted from the main part.

n -Categories

Our 3-Categories are Gray-Categories

Whenever details are relevant, we shall restrict to $n = 3$ and use 3-categories which are *Gray-categories*. In such 3-categories everything is strict, except possibly the exchange law for 2-morphisms.

Our 3-Groupoids are Gray-Groupoids

Given a notion of n -category, an n -groupoid is an n -category in which all k -morphisms are equivalences.

([back](#))

n -Groups

Definition

An n -group is an n -groupoid with a single object.

Notation

We write $G_{(n)}$ when we think of an n -group as a monoidal $(n-1)$ -groupoid. Then we write $\Sigma G_{(n)}$ for the corresponding one-object n -groupoid.

Example

For $G_{(2)}$ any strict 2-group, $\text{INN}_0(G_{(2)})$ is a 3-group coming from a Gray-groupoid.

([back](#))

Tangent n -Categories

Definition

For C any n -category, let $\mathbf{pt} := \{ \bullet \xrightarrow{\sim} \circ \}$ be the “fat point”, and let

$$TC \subset \mathrm{Hom}(\mathbf{pt}, C)$$

be that maximal sub- n -category which collapses to a 0-category when pulled back along the inclusion $\{\bullet\} \hookrightarrow \mathbf{pt}$.

We call TC the *tangent n -category* to C .

The general concept of tangent categories is important for the bigger picture underlying our discussion, but for the present purpose many of the details are not crucial.

Properties of the tangent n -category

Proposition

For C any n -category, we find that the n -bundle

$$TC \rightarrow \mathrm{Obj}(C)$$

is a “deformation retract” of $\mathrm{Obj}(C)$ in that

$$TC \simeq \mathrm{Obj}(C).$$

Properties of the tangent n -category

Proposition

For C any n -category, we find that the n -bundle

$$TC \rightarrow \mathrm{Obj}(C)$$

sits inside an exact sequence

$$\mathrm{Mor}(C) \rightarrow TC \rightarrow C.$$

When $C = \Sigma G_{(n)}$ comes from an n -groupoid, this “is” the universal $G_{(n)}$ -bundle.

Inner automorphism $(n+1)$ -Groups

Definition

Given any n -groupoid C , we say that

$$\text{INN}(C) := T_{\text{Id}_C}(\text{AUT}(C))$$

is the inner automorphism $(n+1)$ -group of C .

proposition

$\text{INN}(G_{(n)})$ sits inside the exact sequence

$$1 \rightarrow Z(G_{(n)}) \rightarrow \text{INN}(G_{(n)}) \rightarrow \text{AUT}(G_{(n)}) \rightarrow \text{OUT}(G_{(n)}) \rightarrow 1$$

of $(n+1)$ -groups. Here $Z(G_{(n)})$ is the (suspension of) the categorical center of $\Sigma G_{(n)}$.

Inner automorphism $(n + 1)$ -Groups

Proposition

There is a canonical inclusion

$$TC \hookrightarrow T_{\text{Id}_C}(\text{AUT}(C)).$$

Definition

The image under this inclusion we call

$$\text{INN}_0(C) \subset \text{INN}(C).$$

([back](#))