

Exercises on the course 'Supersymmetric field theories'

Antoine Van Proeyen, Hamburg Sep. 12-14, 2012

The course is based on chapters of the book 'Supergravity', D. Freedman, and A. Van Proeyen, Cambridge Univ. Press, 2012

A link for a discount prize of 20% is

<http://www.cambridge.org/knowledge/discountpromotion?code=L2SUPE>

1 Scalar field theory and its symmetries

Ex. 1.5 Show that the action

$$S = \int d^D x \mathcal{L}(x) = -\frac{1}{2} \int d^D x [\eta^{\mu\nu} \partial_\mu \phi^i \partial_\nu \phi^i + m^2 \phi^i \phi^i]. \quad (1)$$

is invariant under the transformation

$$\phi^i(x) \xrightarrow{\Lambda} \phi^i(\Lambda x) \equiv \phi^i(\Lambda x). \quad (2)$$

Remember: $(\Lambda x)^\mu = \Lambda^\mu{}_\nu x^\nu$ and

$$\Lambda^\mu{}_\rho \eta_{\mu\nu} \Lambda^\nu{}_\sigma = \eta_{\rho\sigma}. \quad (3)$$

Only fields transform, not spacetime coordinates

Ex. 1.6 Compute the commutators $[L_{[\mu\nu]}, L_{[\rho\sigma]}]$ and show that they agree with that of

$$[m_{[\mu\nu]}, m_{[\rho\sigma]}] = \eta_{\nu\rho} m_{[\mu\sigma]} - \eta_{\mu\rho} m_{[\nu\sigma]} - \eta_{\nu\sigma} m_{[\mu\rho]} + \eta_{\mu\sigma} m_{[\nu\rho]} \quad (4)$$

for matrix generators. Show that to first order in $\lambda^{\rho\sigma}$

$$\phi^i(x^\mu) - \frac{1}{2} \lambda^{\rho\sigma} L_{[\rho\sigma]} \phi^i(x^\mu) = \phi^i(x^\mu + \lambda^{\mu\nu} x_\nu). \quad (5)$$

Remember:

$$L_{[\rho\sigma]} \equiv x_\rho \partial_\sigma - x_\sigma \partial_\rho. \quad (6)$$

2 The Dirac field

Ex. 2.9 Show, using only

$$\{\gamma^\mu, \gamma^\nu\} \equiv \gamma^\mu \gamma^\nu + \gamma^\nu \gamma^\mu = 2 \eta^{\mu\nu} \mathbb{1}, \quad (7)$$

that $[\Sigma^{\mu\nu}, \gamma^\rho] = 2\gamma^{[\mu} \eta^{\nu]\rho} = \gamma^\mu \eta^{\nu\rho} - \gamma^\nu \eta^{\mu\rho}$.

Prove the consistency of

$$\delta\Psi = -\frac{1}{2} \lambda^{\mu\nu} \Sigma_{\mu\nu} \Psi, \quad \delta\bar{\Psi} = \frac{1}{2} \lambda^{\mu\nu} \bar{\Psi} \Sigma_{\mu\nu}$$

Prove then the invariance of the action

$$S[\bar{\Psi}, \Psi] = - \int d^D x \bar{\Psi} [\gamma^\mu \partial_\mu - m] \Psi(x)$$

3 Clifford algebras and spinors

Ex 3.40 Rewrite

$$S[\Psi] = -\frac{1}{2} \int d^D x \bar{\Psi} [\gamma^\mu \partial_\mu - m] \Psi(x)$$

as

$$\begin{aligned} S[\psi] &= -\frac{1}{2} \int d^4 x [\bar{\Psi} \gamma^\mu \partial_\mu - m] (P_L + P_R) \Psi \\ &= - \int d^4 x [\bar{\Psi} \gamma^\mu \partial_\mu P_L \Psi - \frac{1}{2} m \bar{\Psi} P_L \Psi - \frac{1}{2} m \bar{\Psi} P_R \Psi] . \end{aligned}$$

and prove that the Euler-Lagrange equations are

$$\not{\partial} P_L \Psi = m P_R \Psi, \quad \not{\partial} P_R \Psi = m P_L \Psi. \quad (8)$$

Derive $\square P_{L,R} \Psi = m^2 P_{L,R} \Psi$ from the equations above.

Answers to questions asked

- Give an explicit construction of γ -matrices
- To construct a C that satisfies

$$(C\Gamma^{(r)})^T = -t_r C\Gamma^{(r)}, \quad t_r = \pm 1, \quad (9)$$

where $\Gamma^{(r)}$ is a matrix in the set

$$\{\Gamma^A = \mathbb{1}, \gamma^\mu, \gamma^{\mu_1\mu_2}, \gamma^{\mu_1\mu_2\mu_3}, \dots, \gamma^{\mu_1\cdots\mu_D}\} \quad (10)$$

of rank r , the following two properties are relevant

$$C^T = -t_0 C, \quad (C\gamma^\mu)^T = t_1 C\gamma^\mu. \quad (11)$$

or the last one is also

$$\gamma^{\mu T} = t_0 t_1 C\gamma^\mu C^{-1}$$

Indeed for all indices different

$$\begin{aligned} (C\gamma^{\mu_1\mu_2\cdots\mu_r})^T &= (C\gamma^{\mu_1}C^{-1}C\gamma^{\mu_2}C^{-1}\cdots C\gamma^{\mu_r}C^{-1}C)^T \\ &= -t_0(t_0t_1)^r C\gamma^{\mu_r\cdots\mu_1} \end{aligned} \quad (12)$$

See that this gives a number for $r \bmod 4$.

- See that there are two solutions C_+ and C_- for even dimension,
- To go to another representation

$$\gamma'^\mu = S\gamma^\mu S^{-1}, \quad C' = S^{-1T}C S^{-1}.$$