

# Transgression of gauge group cocycles

Locally smooth 3-cocycles, gerbes, category of CAR representations

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# Motivation from gauge theory

$L^2$  condition on the curvature form of a Yang-Mills connection: The connection form at infinity in  $\mathbb{R}^n$  is a pure gauge mod terms of order  $1/r^{n/2+\epsilon}$ . Denote by  $\mathcal{G}_n$  the group of smooth based maps  $S^n \rightarrow G$ . Up to homotopy, the moduli space  $\mathcal{A}/\mathcal{G}_n$  is then parametrized by  $\text{Map}(S^{n-1}, G)$ . Up to homotopy, the bundle  $\mathcal{G}_n \rightarrow \mathcal{A} \rightarrow \mathcal{A}/\mathcal{G}_n$  is then the bundle

$$\mathcal{G}_n \rightarrow P \rightarrow \text{Map}(S^{n-1}, G)$$

where  $P$  is contractible and  $\mathcal{G}_0$  acts freely on  $P$ ; restricting everything to based maps we can take  $P$  as the group of paths  $f(t)$  in  $\mathcal{G}_{n-1}$  with  $f(0) = id$  and we get the fibration

$$\mathcal{G}_n \rightarrow P_n \rightarrow \mathcal{G}_{n-1}.$$

# Motivation from gauge theory

In particular, for  $n = 1$  we have  $\mathcal{G}_1 \rightarrow P_1 \rightarrow G$  a fibration over the finite dimensional group  $G$ , the fiber  $\mathcal{G}_1 = \Omega G$  the based loop group.

When  $G$  is simple compact Lie group  $\Omega G$  has up to isomorphism a unique central extension  $\hat{\Omega}_k G$  for each **level**  $k \in \mathbb{Z}$ . The extension can be given as a **locally smooth** 2-cocycle  $c_2 : \Omega G \times \Omega G \rightarrow \mathcal{S}^1$ . This cocycle is obtained from a class  $\omega_3 \in H^3(G, \mathbf{Z})$  which corresponds to a Lie algebra cohomology class in  $H^3(\mathfrak{g})$ .

So one can ask whether there is a corresponding cocycle in third group cohomology of  $G$ . The answer is yes if one considers again the locally smooth cohomology. About the meaning of the 3-cocycle later.....

# Relation to the BRS complex

Anomalies in quantized gauge theory can be computed from the BRS double complex. It starts from an even form  $\omega^{2n,0}$  which is a characteristic class of a vector bundle over the physical space-time  $M$ . Locally, we have  $\omega^{2n,0} = d\omega^{2n-1,0}$  where  $\omega^{2n-1,0}$  is a Chern-Simons form. One continues

$$\delta\omega^{2n-1,0} = d\omega^{2n-2,1}$$

where  $\delta$  is the coboundary operator in Lie algebra cohomology, here the Lie algebra is the algebra of infinitesimal gauge transformations. Next

$$\delta\omega^{2n-2,1} = d\omega^{2n-3,2}$$

and so on; the second index is the Lie algebra cohomology degree. In particular  $\omega^{2n-2,1}$  is the (infinitesimal) gauge anomaly and  $\omega^{2n-3,2}$  is the commutator anomaly (in space dimension  $2n - 3$ ).

Here we want to address the same problem on the level of locally smooth group cocycles.

# The 3-cocycle: categorical representation

$\mathcal{C}$  an abelian category,  $G$  a group

$g \in G$ ,  $F_g$  a functor in  $\mathcal{C}$

$i_{g,h} : F_g \circ F_h \rightarrow F_{gh}$  an isomorphism

$i_{g,hk} \circ i_{h,k}$  and  $i_{gh,k} \circ i_{g,h}$  isomorphisms  $F_g \circ F_h \circ F_k \rightarrow F_{ghk}$

They are not necessarily equal; one can have a *central extension*

$i_{g,hk} \circ i_{h,k} = \alpha(g, h, k) i_{gh,k} \circ i_{g,h}$  with  $\alpha(g, h, k) \in \mathbf{C}^\times$  a 3-cocycle

# 3-cocycles

Let  $\mathcal{B}$  be an associative algebra and  $G$  a group. Assume that we have a group homomorphism  $s : G \rightarrow \text{Out}(\mathcal{B})$  where  $\text{Out}(\mathcal{B})$  is the group of outer automorphisms of  $\mathcal{B}$ , that is,  $\text{Out}(\mathcal{B}) = \text{Aut}(\mathcal{B})/\text{In}(\mathcal{B})$ , all automorphisms modulo the normal subgroup of inner automorphisms.

If one chooses any lift  $\tilde{s} : G \rightarrow \text{Aut}(\mathcal{B})$  then we can write

$$\tilde{s}(g)\tilde{s}(g') = \sigma(g, g') \cdot \tilde{s}(gg')$$

for some  $\sigma(g, g') \in \text{In}(\mathcal{B})$ . From the definition follows immediately the cocycle property

$$\sigma(g, g')\sigma(gg', g'') = [\tilde{s}(g)\sigma(g', g'')\tilde{s}(g)^{-1}]\sigma(g, g'g'')$$

# Prolongation by central extension

Let next  $H$  be any central extension of  $\text{In}(\mathcal{B})$  by an abelian group  $a$ . That is, we have an exact sequence of groups,

$$1 \rightarrow a \rightarrow H \rightarrow \text{In}(\mathcal{B}) \rightarrow 1.$$

Let  $\hat{\sigma}$  be a lift of the map  $\sigma : G \times G \rightarrow \text{In}(\mathcal{B})$  to a map  $\hat{\sigma} : G \times G \rightarrow H$  (by a choice of section  $\text{In}(\mathcal{B}) \rightarrow H$ ). We have then

$$\begin{aligned} \hat{\sigma}(g, g')\hat{\sigma}(gg', g'') &= [\tilde{s}(g)\hat{\sigma}(g', g'')\tilde{s}(g)^{-1}] \\ &\times \hat{\sigma}(g, g'g'') \cdot \alpha(g, g', g'') \text{ for all } g, g', g'' \in G \end{aligned}$$

where  $\alpha : G \times G \times G \rightarrow a$ .



# The 3-cocycle condition

Here the action of the outer automorphism  $s(g)$  on  $\hat{\sigma}(\ast)$  is defined by  $s(g)\hat{\sigma}(\ast)s(g)^{-1} =$  the lift of  $s(g)\sigma(\ast)s(g)^{-1} \in \text{In}(\mathcal{B})$  to an element in  $H$ . One can show that  $\alpha$  is a 3-cocycle

$$\begin{aligned} &\alpha(g_2, g_3, g_4)\alpha(g_1 g_2, g_3, g_4)^{-1}\alpha(g_1, g_2 g_3, g_4) \\ &\quad \times \alpha(g_1, g_2, g_3 g_4)^{-1}\alpha(g_1, g_2, g_3) = 1. \end{aligned}$$

**Remark** If we work in the category of topological groups (or Lie groups) the lifts above are in general discontinuous; normally, we can require continuity (or smoothness) only in an open neighborhood of the unit element.

Next we construct an example from quantum field theory. Let  $G$  be a compact simply connected Lie group and  $P$  the space of smooth paths  $f : [0, 1] \rightarrow G$  with initial point  $f(0) = e$ , the neutral element, and quasiperiodicity condition  $f^{-1}df$  a smooth function.

$P$  is a group under point-wise multiplication but it is also a principal  $\Omega G$  bundle over  $G$ . Here  $\Omega G \subset P$  is the loop group with  $f(0) = f(1) = e$  and  $\pi : P \rightarrow G$  is the projection to the end point  $f(1)$ . Fix an unitary representation  $\rho$  of  $G$  in  $\mathbf{C}^N$  and denote  $H = L^2(S^1, \mathbf{C}^N)$ .

# CAR representations

For each polarization  $H = H_- \oplus H_+$  we have a vacuum representation of the CAR algebra  $\mathcal{B}(H)$  in a Hilbert space  $\mathcal{F}(H_+)$ . Denote by  $\mathcal{C}$  the category of these representations. Denote by  $a(v)$ ,  $a^*(v)$  the generators of  $\mathcal{B}(H)$  corresponding to a vector  $v \in H$ ,

$$a^*(u)a(v) + a(v)a^*(u) = 2 \langle v, u \rangle$$

and all the other anticommutators equal to zero.

# Outer automorphisms

Any element  $f \in P$  defines a unique automorphism of  $\mathcal{B}(H)$  with  $\phi_f(a^*(v)) = a^*(f \cdot v)$ , where  $f \cdot v$  is the function on the circle defined by  $\rho(f(x))v(x)$ . These automorphisms are in general not inner except when  $f$  is periodic.

We have now a map  $s : G \rightarrow \text{Aut}(\mathcal{B})/\text{In}(\mathcal{B})$  given by  $g \mapsto F(g)$  where  $F(g)$  is an arbitrary smooth quasiperiodic function on  $[0, 1]$  such that  $F(g)(1) = g$ .

Any two such functions  $F(g), F'(g)$  differ by an element  $\sigma$  of  $\Omega G$ ,  $F(g)(x) = F'(g)(x)\sigma(x)$ . Now  $\sigma$  is an inner automorphism through a projective representation of the loop group  $\Omega G$  in  $\mathcal{F}(H_+)$ .

## 3-cocycle

In an open neighborhood  $U$  of the neutral element  $e$  in  $G$  we can fix in a smooth way for any  $g \in U$  a path  $F(g)$  with  $F(g)(0) = e$  and  $F(g)(1) = g$ .

Of course, for a connected group  $G$  we can make this choice globally on  $G$  but then the dependence of the path  $F(g)$  would not be a continuous function of the end point. For a pair  $g_1, g_2 \in G$  we have

$$\sigma(g_1, g_2)F(g_1g_2) = F(g_1)F(g_2)$$

with  $\sigma(g_1, g_1) \in \Omega G$ .

# LG valued 2-cocycle

For a triple of elements  $g_1, g_2, g_3$  we have now

$$\begin{aligned}F(g_1)F(g_2)F(g_3) &= \sigma(g_1, g_2)F(g_1g_2)F(g_3) \\ &= \sigma(g_1, g_2)\sigma(g_1g_2, g_3)F(g_1g_2g_3).\end{aligned}$$

In the same way,

$$\begin{aligned}F(g_1)F(g_2)F(g_3) &= F(g_1)\sigma(g_2, g_3)F(g_2g_3) \\ &= [g_1\sigma(g_2, g_3)g_1^{-1}]F(g_1)F(g_2g_3) \\ &= [g_1\sigma(g_2, g_3)g_1^{-1}]\sigma(g_1, g_2g_3)F(g_1g_2g_3)\end{aligned}$$

which proves the 2-cocycle relation for  $\sigma$ .

## 3-cocycle $\alpha$ for $G$

Lifting the loop group elements  $\sigma$  to inner automorphisms  $\hat{\sigma}$  through a projective representation of  $\Omega G$  we can write

$$\hat{\sigma}(g_1, g_2)\hat{\sigma}(g_1 g_2, g_3) = \text{Aut}(g_1)[\hat{\sigma}(g_2, g_3)]\hat{\sigma}(g_1, g_2 g_3)\alpha(g_1, g_2, g_3),$$

where  $\alpha : G \times G \times G \rightarrow S^1$  is some phase function arising from the fact that the projective lift is not necessarily a group homomorphism.

Since (in the case of a Lie group) the function  $F(\cdot)$  is smooth only in a neighborhood of the neutral element, the same is true also for  $\sigma$  and finally for the 3-cocycle  $\alpha$ .

# The Lie algebra 3-cocycle

An equivalent point of view to the construction of the 3-cocycle  $\alpha$  is this: We are trying to construct a central extension  $\hat{P}$  of the group  $P$  of paths in  $G$  (with initial point  $e \in G$ ) as an extension of the central extension over the subgroup  $\Omega G$ . The failure of this central extension is measured by the cocycle  $\alpha$ , as an obstruction to associativity of  $\hat{P}$ .

On the Lie algebra level, we have a corresponding cocycle  $c_3 = d\alpha$  which is easily computed. The cocycle  $c$  of  $\Omega\mathfrak{g}$  extends to the path Lie algebra  $P\mathfrak{g}$  as

$$c(X, Y) = \frac{1}{4\pi i} \int_{[0, 2\pi]} \text{tr}(XdY - YdX).$$

This is an antisymmetric bilinear form on  $P\mathfrak{g}$  but it fails to be a Lie algebra 2-cocycle. The coboundary is given by



# The Lie algebra 3-cocycle

$$\begin{aligned}(\delta c)(X, Y, Z) &= c(X, [Y, Z]) + c(Y, [Z, X]) + c(Z, [X, Y]) \\ &= -\frac{1}{4\pi i} \text{tr } X[Y, Z]|_{2\pi} = d\alpha(X, Y, Z).\end{aligned}$$

Thus  $\delta c$  reduces to a 3-cocycle of the Lie algebra  $\mathfrak{g}$  of  $G$  on the boundary  $x = 2\pi$ . This cocycle defines by left translations on  $G$  the left-invariant de Rham form  $-\frac{1}{12\pi i} \text{tr} (g^{-1} dg)^3$ ; this is normalized as  $2\pi i$  times an integral 3-form on  $G$ .

Let  $\omega_3$  represent a class in the singular cohomology  $H^3(H, \mathbf{Z})$ . We shall now make the following assumptions: 1) The pull-back  $\pi^*(\omega_3) = d\theta_2$  is trivial on  $G$ . 2)  $H$  and  $G$  are simply connected and  $H_2(G, \mathbf{Z}) = H_2(H, \mathbf{Z}) = 0$ . Using the exact homotopy sequence from the fibration  $N \rightarrow G \rightarrow H$  we conclude that  $N$  is connected and  $\pi_1(N) = 0$  and thus also  $H_1(N, \mathbf{Z}) = 0$ . For each  $g \in G$  we select a path  $g(t)$  with end points  $g(0) = 1 \in G$  and  $g(1) = g$ . We can make the choice  $g \rightarrow g(t)$  in a locally smooth manner close to the neutral element  $1 \in G$ . In addition, since also  $N$  is connected, we may assume that  $g(t) \in N$  if  $g \in N$ . For a triple  $g, g_1, g_2 \in G$  we make a choice of a singular 2-simplex  $\Delta(g; g_1, g_2)$  such that its boundary is given by the union of the 1-simplices  $gg_1(t)$ ,  $gg_1(1)g_2(t)$  and  $g(g_1g_2)(1-t)$ . All this can be made in a locally smooth manner since locally the Lie groups are open contractible sets in a vector space.

$$c_2(g; g_1, g_2) = \exp 2\pi i \langle \Delta(g; g_1, g_2), \theta_2 \rangle$$

using the duality pairing of singular 2-simplices and 2-cochains. This formula does not in general define a group cocycle for  $G$  but it gives a 2-cocycle for the group  $N$  with the right action of  $N$  on  $G$  and the corresponding action of  $N$  on  $A = \text{Map}(G, S^1)$ . To prove that indeed

$$\begin{aligned} (\delta c)_2(g; n_1, n_2, n_3) &= \\ c_2(g; n_1, n_2) c_2(g; n_1 n_2, n_3) c_2(g; n_1, n_2 n_3)^{-1} c_2(g n_1; n_2, n_3)^{-1} &= \\ &= 1 \end{aligned}$$

we just need to observe that the product is given through pairing the cochain  $\theta_2$  with the singular cycle defined as the union of the singular 2-simplices involved in the above formula. All these 2-simplices are in the same  $N$  orbit  $gN$  and since  $d\theta_2 = \pi^*\omega_3$  the cochain  $\theta_2$  is actually an integral cocycle on the  $N$  orbits and the pairing gives an integer  $k$  and  $\exp 2\pi i k = 1$ . For arbitrary  $g_i \in G$  the coboundary  $\delta c_2$  does not vanish but its value

$$(\delta c_2)(g; g_1, g_2, g_3) = \exp 2\pi i \langle \Delta(g; g_1, g_2, g_3), d\theta_2 \rangle$$

is given by pairing  $d\theta_2 = \pi^*\omega_3$  with the singular 3-simplex  $V$  with the boundary consisting of the sum of the faces  $\Delta(g; g_1, g_2), \Delta(g; g_1 g_2, g_3), \Delta(g; g_1, g_2 g_3), \Delta(g g_1; g_2, g_3)$ . But this is the same as  $\exp 2\pi i \langle \pi(V), \omega_3 \rangle$  and therefore it depends only on the projections  $\pi(g), \pi(g_i) \in H$ . Denote by  $c_3 = c_3(h; h_1, h_2, h_3)$  this locally smooth 3-cocycle on  $H$ . (This construction can be extended to higher cocycles under appropriate conditions on the homology groups of  $H$ .)

We may think of the cohomology class  $[c_3]$  as an obstruction to prolonging the principal  $N$  bundle  $G$  over  $H$  to a bundle  $\hat{G}$  with the structure group  $\hat{N}$ . Namely, if such a prolongation exists then there is a 2-cocycle  $c_2$  on  $G$  which when restricted to  $N$  orbits in  $G$  is equal to  $c_2(g; n_1, n_2)$ . If  $c'_2$  is another such a 2-cocycle then  $(\delta c'_2)(\delta c_2)^{-1}$  projects to a trivial 3-cocycle on  $H$ . Conversely, if  $c_3$  on  $H$  is a coboundary of some  $\xi_2$  then  $c'_2 = c_2(\pi^*\xi)^{-1}$  agrees with  $c_2$  on the  $N$  orbits and so the obstruction depends only on the cohomology class  $[c_3]$ .

Wagemann and Wockel defined a map from the locally smooth cohomology of a Lie group  $H$  to its Čech cohomology. There is also a map from the locally smooth group cohomology  $H_S^2(N, A)$  to the Čech cohomology  $\check{H}^2(H, A)$  by the formula

$$c_{ijk}(x) = \hat{\eta}_{ij}(x)\hat{\eta}_{jk}(x)\hat{\eta}_{ki}(x)$$

where  $\psi_i(x)\eta_{ij}(x) = \psi_j(x)$ ,  $\psi_i : U_i \rightarrow G$  are local smooth sections for an open good cover  $\{U_i\}$  of  $H$  and the  $\hat{\eta}_{ij}$ 's are lifts of the transition functions  $\eta_{ij} : U_i \cap U_j \rightarrow N$  to the extension  $\hat{N}$ ; the product on the right is determined by an element in  $H_S^2(N, A)$ . Although these Čech cocycles have values in  $A$  they correspond to a cocycle in  $H^3(H, \mathbf{Z})$  by the usual way, taking differences of logarithms  $\log c_{ijk}/2\pi i$  on intersections  $U_{ijkl}$  which must be integer constants for a good cover.