How to get a spinning string from here to there

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joint work, in parts, with

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inspired, in parts, by

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 $\operatorname{String}(n)$ is some topological group, to be described in a moment. For various reasons it appears to be of interest to understand what a $\operatorname{String}(n)$ -bundle with connection would be. Here are some of these reasons.

1 The need for String(n)-Bundles

Killingback and others noticed [1, 2] that

 1
 super particles
 couple to
 Spin(n)-bundles
 with connection

 like

 2
 super strings
 couple to
 String(n)-bundles
 with (?)

Using the Atiyah-Segal observation that

1 quantum (super) particles are functors $1 \operatorname{Cob}_S \to \operatorname{Hilb}_S$

like

2 quantum (super) strings are functors $2Cob_S \rightarrow Hilb_S$

this should translate into a precise statement (about representations of cobordisms categories). Back then few people thought of categorification. But Stolz and Teichner later made two remarks.

First Remark. First, following Dan Freed, Segal's original viewpoint should be refined to

1	quantum (super) particles	are functors	$1\mathrm{Cob}_S \to \mathrm{Hilb}_S$
	like		
2	quantum (super) strings	are 2-functors	$\operatorname{Cob}_S^{\operatorname{ext}} \to 2\operatorname{Hilb}_S$

This is nowadays known as extended quantum field theory.

Second Remark. Moreover, it should be true that

 1
 Spin(n) bundles with connection are related to
 K-cohomology

 like

2 String(n)-bundles with connection are related to elliptic cohomology

All in all, this is supposed to be considerable reason to be interested in String(n)-bundles with connection.

2 What is String(n), anyway?

There is the classical definition of String(n), and there is a "revisionist" one. The latter is maybe intuitively more accessible.

Revisionist definition: String(n) as stringy Spin(n). Superstrings (in their RNS incarnation) are sometimes called *spinning strings*. Indeed, a superstring is much like a continuous line of spinors.

This suggests that the corresponding gauge group is the loop group

 Ω Spin(n)

or maybe its Kac-Moody central extension

 $\hat{\Omega}_k \operatorname{Spin}(n)$

or maybe the path group

PSpin(n).

Or maybe all of these. In fact, there are canonical group homomorphisms

$$\hat{\Omega}_k \operatorname{Spin}(n) \xrightarrow{t} P \operatorname{Spin}(n) \xrightarrow{\alpha} \operatorname{Aut}(\hat{\Omega}_k \operatorname{Spin}(n))$$

These satisfy two compatibility conditions which say that the groups here conspire to form a (strict Fréchet-Lie) 2-group

 $G_{(2)}$.

A 2-group is a category which behaves like a group.

Every topological category like this may be turned into a big ordinary topological space by taking its nerve and forming its geometric realization.

	0-simplices	1-simplices	2-simplices	etc.
nerve of category: (simplicial set)			g_3 $h_1 \swarrow h_2$	
(empreud cee)	g	$g_1 \xrightarrow{h} g_2$	$g_1 \xrightarrow{h_1 \cdot h_2} g_2$	
geometric	fill with standar	d simplices in \mathbb{R}^n	– glue along comn	non fa

If the category is a 2-group, the realization of its nerve is a topological group.

For $G_{(2)}$, this nerve is [8, 9]

realization:

$$|G_{(2)}| \simeq \operatorname{String}(n)$$
.

ON FIRST READING, SKIP TO 3, NOW

Classical definition. This is all that is needed about String(n) in the following. But for completeness, here is the classical definition.

Definition 1 The string group String_G of a simple, simply connected, compact topological group G is (a model for) the 3-connected topological group with the same homotopy groups as G, except

$$\pi_3(\operatorname{String}_G) = 0$$
,

which, furthermore, fits into the exact sequence

$$1 \longrightarrow (BU(1) \simeq K(\mathbb{Z}, 2)) \longrightarrow \operatorname{String}_G \longrightarrow G \longrightarrow 1$$

of topological groups.

The string group proper is obtained by setting G = Spin(n).

$$\operatorname{String}(n) := \operatorname{String}_{\operatorname{Spin}(n)}.$$

The way to see that such a group is a plausible candidate for something generalizing the Spin-group, which, recall, fits into the exact sequence

$$1 \to \mathbb{Z}_2 \to \operatorname{Spin}(n) \to SO(n) \to 1$$
,

is to note that the first few homotopy groups π_k of O(n) are

$$\frac{k}{\pi_k(O(n))} = \begin{bmatrix} 0 & 1 & 2 & 3 & 4 & 5 & 6 & 7 \\ \hline \pi_k(O(n)) = \begin{bmatrix} \mathbb{Z}_2 & \mathbb{Z}_2 & 0 & \mathbb{Z}_2 & 0 & 0 & 0 & \mathbb{Z} \end{bmatrix}$$

Starting with O(n), we can successively "kill" the lowest nonvanishing homotopy groups, thus obtaining first SO(n) (the connected component), then Spin(n) (the universal cover) and finally String(n) (the 3-connected cover). Notice that with π_3 vanishing, String(n) cannot be a Lie group – but it can be a Lie 2-group.

Usually (see [5]), the definition of String_G includes also a condition on the boundary map $\pi_3(G) \xrightarrow{\partial} \pi_2(K(\mathbb{Z},2))$. Our definition above is really geared towards the application where $G = \operatorname{Spin}(n)$, for which we find it more natural.

Namely, recall that every short exact sequence of topological groups

$$0 \to A \to B \to C \to 0 \,,$$

which happens to be a fibration, gives rise to a long exact sequence of homotopy groups:

$$\cdots \longrightarrow \pi_n(A) \longrightarrow \pi_n(B) \longrightarrow \pi_n(C) \xrightarrow{\partial} \pi_{n-1}(A) \longrightarrow \cdots$$

In our case this becomes

$$\cdots \longrightarrow \pi_n(K(\mathbb{Z},2)) \longrightarrow \pi_n(\operatorname{String}_G) \longrightarrow \pi_n(G) \xrightarrow{\partial} \pi_{n-1}(K(\mathbb{Z},2)) \longrightarrow \cdots$$

Demanding that $\pi_3(\operatorname{String}_G) = 0$ and assuming that also $\pi_2(\operatorname{String}_G) = 0$ (which we noticed above is the case for $G = \operatorname{Spin}(n)$) implies that we find inside this long exact sequence the short exact sequence

$$0 \longrightarrow (\pi_3(G) \simeq \mathbb{Z}) \xrightarrow{\partial} \mathbb{Z} \longrightarrow 0 .$$

But this implies that the boundary map ∂ here is an isomorphism, hence that it acts on \mathbb{Z} either by multiplication with k = 1 or k = -1. (This number is really the "level" governing this construction. If I find the time I will explain this later.)

In [5] this logic is applied the other way around. Instead of demanding that $\pi_3(\operatorname{String}_G) = 0$ it is demanded that the boundary map

$$\pi_3(G) \xrightarrow{\partial} \mathbb{Z}$$

is given by multiplication with the level, namely a specified element in $H^4(BG)$.

3 Connections, functorially: from here to there

Our statement

1 quantum (super) particles are functors $1 \operatorname{Cob}_S \to \operatorname{Hilb}_S$

like

2 quantum (super) strings are 2-functors $\operatorname{Cob}_S^{\operatorname{ext}} \to 2\operatorname{Hilb}_S$

should be thought of as the quantization of the statement

1	parallel transport of	particles	is a functor	$\mathcal{P}_1(X) \to \operatorname{Vect}$
	like			

2 parallel transport of strings is a 2-functor $\mathcal{P}_2(X) \to 2 \operatorname{Vect}$

Here $\mathcal{P}_n(X)$ is the *n*-groupoid of *n*-paths in the space X.



In fact, a smooth spinor bundle with connection is enirely encoded [17] in a smooth functor

$$\operatorname{tra}: \mathcal{P}_1(X) \to \operatorname{Vect}$$

which sends paths to the parallel transport along them

$$\operatorname{tra}: \left(\begin{array}{cc} x \xrightarrow{\gamma} y \end{array} \right) \mapsto \left(\begin{array}{cc} V_x \xrightarrow{\operatorname{tra}(\gamma)} & V_y \end{array} \right)$$

(fiber over x) $\xrightarrow{}$ identity along path γ (fiber over y)

and which factors locally through a functor with values in Spin(n):



4 2-Connections on String(n)-2-bundles

The above has a straightforward categorification. But first let's look at what Stolz and Teichner did.

Stolz-Teichner's definition of a connection on a String(n)-bundle. They observe the following important

Fact. To every String(n)-bundle is canonically associated a vonNeumann algebra bundle.

But algebras, vonNeumann or not, naturally live not in a 1-category, but in the 2-category

Bim.

So with the String(n)-associated algebra bundle given, Stolz and Teichner

consider parallel transport 2-functors



Why? Remember: String(n) is the realization of the nerve of a 2-group $G_{(2)}$.

To get a $G_{(2)}$ -associated 2-vector bundle,

$$\begin{array}{c|c} & \mathcal{P}_{2}(Y) & \xrightarrow{\pi} & \mathcal{P}_{2}(X) \\ & & \downarrow \\ (\text{local structure}) & \swarrow \\ & \Sigma G_{(2)} & \xrightarrow{\rho} \\ & & (\text{representation}) \end{array} & 2\text{Vect} \end{array}$$

we just need a 2-representation

$$\rho: \Sigma G_{(2)} \to 2 \text{Vect}.$$

5 The canonical 2-Representation

What is a 2-Vector space? Noticing that

$$\operatorname{Vect}_{\mathbb{C}} = {}_{\mathbb{C}}\operatorname{Mod}$$

and that Vect is again monoidal, we set

$$n$$
Vect = $_{(n-1)}$ Vect Mod .

General 2Vect is large and unwieldy. But we have a chain of canonical inclusions

$$\operatorname{Intertwin} \xrightarrow{} \operatorname{Bim} \xrightarrow{} \operatorname{Vect} \operatorname{Mod}$$

algebras homomorphisms intertwiners

(and, by the way $KV2Vect \longrightarrow Bim \longrightarrow Vect Mod$).

Canonical strict 2-Rep on Intertwin. Every strict 2-group $G_{(2)}$ comes from a crossed module of two ordinary groups

$$G_{(2)} = \left(H \xrightarrow{t} G \xrightarrow{\alpha} \operatorname{Aut}(H) \right).$$

For an ordinary representation

$$\rho_0: \Sigma H \to \text{Vect}$$

of H and with

 $A := \langle \operatorname{im}(\rho_0) \rangle$

the algebra generated from that we get a 2-functor

$$\rho: \Sigma G_{(2)} \to \text{Intertwin} \hookrightarrow \text{Bim} \hookrightarrow 2\text{Vect}$$

by



So:

For every strict 2-group $G_{(2)} = (t : H \to G)$ an ordinary representation ρ_0 of H induces a notion of $G_{(2)}$ -associated 2-vector bundles.

Their typical fiber is (the module category of) the algebra generated from the image of ρ_0 .

For the standard reps of $\hat{\Omega}_k \operatorname{Spin}(n)$ this algebra is a von Neumann type III factor.

6 Line 2-bundles versus String(n)-2-bundles

A simple example to keep in mind are rank-1 2-vector bundles – "line 2bundles". These are the 2-vector bundles canonically associated to the 2group

$$G_{(2)} = \Sigma U(1) = (U(1) \to 1)$$

by the standard representation of U(1) on \mathbb{C} :

- Let $G_{(2)} = \Sigma U(1)$,
- then $|G_{(2)}| \simeq PU(H);$
- local semi trivialization of ρ -associated $\Sigma U(1)$ -2-bundles are line bundle gerbes [18];
- indeed, these have the same classification as PU(H)-bundles, namely by classes in $H^3(X, \mathbb{Z})$.
- The typical fiber is (Morita!)-equivalent to \mathbb{C} , hence these are K(H)bundles (where K(H) is the algebra of finite rank operators.)

We can essentially send this example from $G_{(2)} = \Sigma U(1)$ to the string 2-group by using three results (two of which we already discussed):

Three results that clarify the situation for String(n) :

• [9, 8]: There is a strict Fréchet Lie 2-group

 $G_{(2)} = \operatorname{Spin}(n)_k = (t : \hat{\Omega}_k \operatorname{Spin}(n) \to P \operatorname{Spin}(n))$

such that String(n) is the geometrical realization of its nerve

 $\operatorname{String}(n) \simeq |\operatorname{Spin}(n)_k|.$

- [6, 7]: $G_{(2)}$ -2-bundles have the same classification as $|G_{(2)}|$ -1-bundles.
- For every strict Lie 2-group there is a canonical 2-representation

$$G_{(2)} \xrightarrow{\rho}$$
Intertwin $\hookrightarrow Bim \hookrightarrow 2Vect$.

Caveat : Extending the canonical 2-rep from Lie to Fréchet is technically subtle. One has to pass from finite-dimensional vector spaces to Hilbert spaces and replace bimodule tensor products by Connes fusion.

	line 2-bundle	String 2-bundle
structure 2-group	$(U(1) \rightarrow 1)$	$(\hat{\Omega}\mathrm{Spin}(n) \to P\mathrm{Spin}(n))$
nerve of that	PU(H)	$\operatorname{String}(n)$

associated 2-vector bundle finite-rank operators von-Neumann algebras

Bottom line: 2-Vector bundles with 2-connection associated by the canonical 2-rep to $G_{(2)}$ -principal 2-bundles give a good definition for $\operatorname{String}(n) \simeq |G_{(2)}|$ connections.

This definition is, in many respects, very similar to that by Stolz-Teichner (for instance: fake flatness!).

Is it equivalent? I don't know.

But it is useful, for instance for the following considerations.

7 Cartan connection with values in Lie *n*algebras

When passing from the topological group $\operatorname{String}(n)$ to the 2-group $G_{(2)}$, we pass from the realm of pure topology to the realm of differential geometry. But it's still infinite-dimensional (Fréchet) differential geometry. We might want to go one step further and describe connections on $\operatorname{String}(n)$ -bundles in terms of plain old differential forms.

$\longleftarrow \text{topological}$	——Fréchet Lie——	$Lie \longrightarrow$
2-functors on String (n) -associated bundles	associated $G_{(2)}$ -2-vector transport	$\operatorname{Lie}(G_{(2)})$ -Cartan connection
Stolz-Teichner	BCSSW	Stasheff-S.

This can be done by passing from Lie *n*-groups to Lie *n*-algebras.

Differentiating parallel *n***-transport**. For differentiating a parallel transport 2-functor $\mathcal{D}_{n}(X) \to \Sigma C$

$$\mathcal{P}_2(X) \to \Sigma G_{(2)}$$

it is convenient to first pass to the corresponding curvature 3-transport [19, 18, 20].

	transport	curvature	Bianchi
	n	(n+1)	(n+2)
	tra: $\mathcal{P}_n(X) \to \Sigma G_{(n)}$	$\operatorname{curv}_{\operatorname{tra}}: \Pi_{n+1}(X) \to \Sigma \operatorname{INN}(G_{(n)})$	$\mathrm{curv}_{\mathrm{curv}_{\mathrm{tra}}}$
	arbitrary	flat (as $(n + 1)$ -transport) = Bianchi identity	trivial (as $(n+2)$ -transport)
n = 1	$A\in \Omega^1(Y,\mathfrak{g})$	$F_A \in \Omega^2(Y, \mathfrak{g})$	$d_A F_A \in \Omega^3(Y, \mathfrak{g})$
n = 2	$\left(\begin{smallmatrix} A,\\ B\end{smallmatrix}\right)\in \begin{smallmatrix} \Omega^1(Y,\mathfrak{g})\\ \times\Omega^2(Y,\mathfrak{h})\end{smallmatrix}$	$\begin{pmatrix} \beta := F_A + t_*B, \\ H := d_AB \end{pmatrix} \in \begin{array}{c} \Omega^2(Y, \mathfrak{g}) \\ \times \Omega^3(Y, \mathfrak{h}) \end{pmatrix}$	$\begin{pmatrix} d_A\beta - t_*H, \\ d_AH - \beta \wedge H \end{pmatrix} \in \begin{array}{c} \Omega^3(Y, \mathfrak{g}) \\ \times \Omega^4(Y, \mathfrak{g}) \end{pmatrix}$

For one, smooth (n+1)-functors on the fundamental (n+1)-groupoid $\Pi_{n+1}(Y)$ are related, differentially, to *pseudo*-1-functors on the pair groupoid.

(illustration goes here)

So upon differentiation, our parallel transport n-functors turn into morphisms from the tangent algebroid

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\operatorname{Vect}(Y)
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```
to the Lie (n + 1)-algebra
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 $\operatorname{inn}(\mathfrak{g}_{(n)})$

of inner derivations of the Lie n-algebra

$$\mathfrak{g}_{(n)} := \operatorname{Lie}(G_{(n)}) \,.$$

	Lie <i>n</i> -groupoids	differentiation	Lie <i>n</i> -algebras (\simeq <i>n</i> -term L_{∞} -algebras)	$\begin{array}{rl} {\color{black} \textbf{quasi free}} \\ {\color{black} \textbf{differential}} \\ \simeq & {\color{black} \textbf{graded commutative}} \\ & {\color{black} \textbf{algebras}} \\ & {\color{black} (qfDGCAs)} \end{array}$
morphism	$\Sigma(\text{INN}(G_{(n)}))$ $ \downarrow^{F}_{\Pi_{n+1}(P)} $		$\operatorname{inn}(\mathfrak{g}_{(n)})$	$(\bigwedge^{\bullet} s\mathfrak{g}^{*}_{(n)}, d_{\operatorname{inn}(\mathfrak{g}_{(n)})}) \\ \downarrow^{f^{*}} \\ (\Omega^{\bullet}(P), d)$
description	smooth pseudofunctor from pair groupoid of X to inner automorphisms of structure Lie <i>n</i> -group $G_{(n)}$		morphism of Lie <i>n</i> -algebroids $\simeq n$ -term L_{∞} -algebras from tangent algebroid of X to inner derivation Lie $(n + 1)$ -algebra $\mathfrak{g}_{(n)} := \operatorname{Lie}(G_{(n)})$	dual morphism of qfDGCAs

The *n*-connections which factor through $\mathfrak{g}_{(n)}$ itself, however, are the *flat n*-connections: all their curvature *k*-forms vanish:



8 String and Chern-Simons connections

In fact, in terms of Lie 2-algebras, that big scary topological group String(n) becomes a small, handy Lie 2-algebra [8, 9]

$$\operatorname{string}(\mathfrak{g}) := \mathfrak{g}_{\mu},$$

which is essentially just a simple Lie algebra \mathfrak{g} together with a multiple of the canonical 3-cocycle μ .

This is just a first example of a general pattern [11, 13]:

• For every Lie (n + 1)-cocycle μ on \mathfrak{g} there is a Lie *n*-algebra

 \mathfrak{g}_{μ} .

• For every invariant degree (n + 1)-polynomial k on \mathfrak{g} there is a Lie (2n + 1)-algebra

 $\operatorname{ch}_k(\mathfrak{g})$.

• For every transgressive element (Chern-Simons form) there is a short (weakly) exact sequence

$$0 \longrightarrow \mathfrak{g}_{\mu_k} \longrightarrow \operatorname{cs}_k(\mathfrak{g}) \longrightarrow \operatorname{ch}_k(\mathfrak{g}) \longrightarrow 0.$$





Chern-Simons

Chern

The Lie (2n+1)-algebra

 $\operatorname{cs}_k(\mathfrak{g})$

is essentially defined by the fact that a connection form with values in it is the corresponding Chern-Simons form $CS_k(A)$:



For the special case of interest here, String(n), where everything is controlled by the canonical 3-cocycle

$$\mu = \langle \cdot, [\cdot, \cdot] \rangle$$

and the corresponding characteristic class

$$k(F_A) = \langle F_A \wedge F_A \rangle$$

we have



,

for

$$(A, B, C) \in \Omega^1(X, \mathfrak{g}) \times \Omega^2(X) \times \Omega^3(X)$$

This should be one incarnation of the statement found in [5]:

String connections trivialize Chern-Simons theory.

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