

On the monoidal structure of matrix bi-factorisations

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We investigate tensor products of matrix factorisations. This is most naturally done by formulating matrix factorisations in terms of bimodules instead of modules. If the underlying ring is $\mathbb{C}[x_1, \dots, x_N]$ we show that bimodule matrix factorisations form a monoidal category.

This monoidal category has a physical interpretation in terms of defect lines in a two-dimensional Landau-Ginzburg model. There is a dual description via conformal field theory, which in the special case of $W = x^d$ is an $\mathcal{N} = 2$ minimal model, and which also gives rise to a monoidal category describing defect lines. We carry out a comparison of these two categories in certain subsectors by explicitly computing 6j-symbols.

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1. Introduction and summary

Matrix factorisations first appeared in mathematics in the study of maximal Cohen-Macaulay modules [Ei, Bu]. Later it was realised in the context of string theory that they also describe boundary conditions in two-dimensional $\mathcal{N} = 2$ supersymmetric Landau-Ginzburg models [Kt, KL1, BHLS, La]. Matrix factorisations form a category, and in the relation to Landau-Ginzburg models the objects of this category are interpreted as boundary conditions and the morphisms as boundary fields. The same structure appears in the study of two-dimensional conformal field theory. There, a boundary condition can be defined as a solution to a set of consistency conditions, the factorisation or sewing constraints [CL, Le, Kg, Fj], and again one obtains a category whose objects and morphisms correspond to boundary conditions and boundary fields, respectively.

Mathematically these two descriptions are very different, one is the category of matrix factorisations over a polynomial ring, the other is constructed from the representations of a vertex operator algebra. However, physically one expects that the infrared fixed point of the $\mathcal{N} = 2$ Landau-Ginzburg model is an $\mathcal{N} = 2$ supersymmetric conformal field theory, and that the supersymmetry preserving boundary conditions in these two theories should be related. This physical reasoning has been tested by matching subsectors of the two categories of boundary conditions, see e. g. [ADD, Ho, BG, ERR].

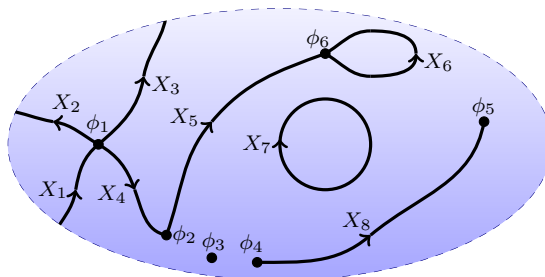
Apart from boundary conditions, there is another natural structure for two-dimensional field theories one can study, namely defect lines. A defect line is a line of inhomogeneity on the two-dimensional worldsheet, and the possible “defect conditions” again form a category. This is not surprising, because via the folding trick [WA] one can think of a defect line in a given theory as a boundary for the corresponding folded theory, and so the same mathematical framework applies. However, there is one crucial difference between boundaries and defects, and this is that several defect lines can meet in junction points. As we will outline below, defect junctions can be encoded by endowing the category of defects with a monoidal structure.

On the side of matrix factorisations, the relevant tensor product is known from [Yo, ADD, KhR, BR], but the full monoidal structure, including the associator and unit isomorphisms, has not yet been coherently given. We close this gap in the case of polynomial rings $R = \mathbb{C}[x_1, \dots, x_N]$ and prove that defects in any topologically B-twisted Landau-Ginzburg model indeed form a monoidal category. Some of the necessary ingredients are already contained in [KhR, BR]; we formulate them in the language of bimodule matrix factorisations – or *matrix bi-factorisations* for short – and show that they satisfy the necessary coherence conditions for a monoidal category (theorems 2.8 and 2.18). We also prove a corresponding statement for graded matrix bi-factorisations (theorems 2.17 and 2.18).

The infrared fixed point of the $\mathcal{N} = 2$ Landau-Ginzburg model with superpotential $W = x^d$ is the A-type $\mathcal{N} = 2$ super-Virasoro minimal model of central charge $c = 3 - \frac{6}{d}$ [KMS, Ma, VW, HW], and one can now try to compare the corresponding monoidal categories of defect conditions. Before describing the monoidal structure and the comparison calculations, let us motivate in more detail how the tensor product arises on the category of defect conditions.

Defects in two-dimensional quantum field theory

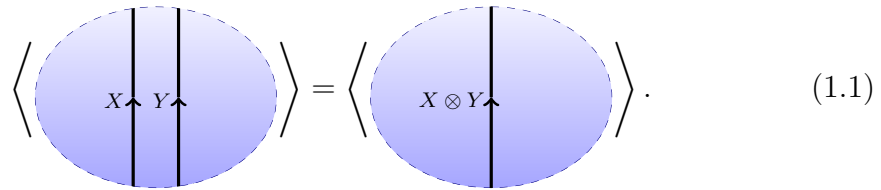
A defect line can either be an oriented closed loop, or an oriented interval joining two field insertion points. Each field insertion can be the start and end point of an arbitrary finite number of defect lines, resulting in n -valent defect junctions, $n = 0, 1, 2, \dots$. A typical piece of a worldsheet might look as follows:



Here we label the defect lines by defect conditions X_i and the field insertions by fields ϕ_i from the corresponding space of n -valent junction fields. The quantum

field theory assigns a value, the correlator, to each such worldsheet. These correlators are subject to consistency conditions arising from factorising a complicated worldsheet into simpler building blocks, see [RS, LMZ].

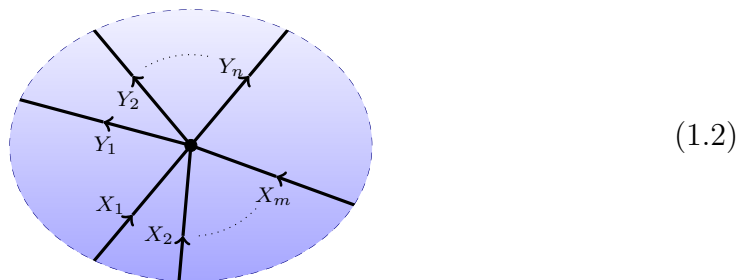
We are particularly interested in a special kind of defect condition, namely *topological defects*. These have the additional property that the path of the defect line can be deformed on the worldsheet without affecting the value of the correlator, as long as the defect line is not taken across field insertions or other defect lines. If we move two defect lines labelled by topological defect conditions X and Y arbitrarily close to each other, we obtain a new topological defect, the *fused defect*, and we write its defect condition as $X \otimes Y$:



$$\langle \text{Worldsheet with two lines } X \text{ and } Y \rangle = \langle \text{Worldsheet with one line } X \otimes Y \rangle. \quad (1.1)$$

In these pictures we show a fraction of the worldsheet, and it is understood that the two worldsheets differ only in the way indicated. The equality means that the correlators assigned to the two worldsheets coincide. We will denote the trivial defect – or identity defect – by $\mathbf{1}$. It acts as the identity under fusion: $\mathbf{1} \otimes X \cong X$ and $X \otimes \mathbf{1} \cong X$.

Similarly, we are interested in *topological junction fields*, which are junction fields that can be moved on the worldsheet without affecting the amplitude. Not all junction fields have this property, and whether or not there are any at all depends on the defect conditions placed on the defect lines joining at a junction. Suppose we have a defect junction with m defect lines X_1, \dots, X_m pointing towards the junction and n defect lines Y_1, \dots, Y_n pointing away from it:



$$\langle \text{Worldsheet with junction and lines } X_1, X_2, X_m, Y_1, Y_2, Y_n \rangle \quad (1.2)$$

We denote the \mathbb{C} -vector space of topological junction fields by¹

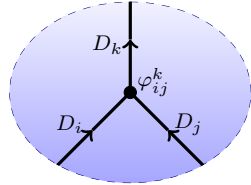
$$\text{Hom}_{\mathcal{D}}(X_1 \otimes \dots \otimes X_m, Y_1 \otimes \dots \otimes Y_n). \quad (1.3)$$

The defect category \mathcal{D} has as objects the different defect conditions, and as morphisms the topological junction fields between them. The fusion of defect lines should endow \mathcal{D} with the structure of a monoidal category, and indeed it does so in the theories we will consider.

Suppose that \mathcal{D} has a semi-simple subsector with finitely many simple objects that is closed under fusion. Denote this subsector by \mathcal{D}' and the simple objects by D_i , $i \in \mathcal{I}$ for some index set \mathcal{I} . The D_i label elementary topological defects, which means that up to scalars the identity field id_{D_i} on the defect D_i is the unique topological (2-valent) junction field, $\text{Hom}_{\mathcal{D}}(D_i, D_i) = \mathbb{C} \cdot \text{id}_{D_i}$. The fusion of two defects decomposes into a direct sum

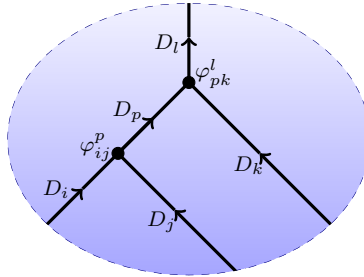
$$D_i \otimes D_j \cong \bigoplus_{k \in \mathcal{I}} N_{ij}^k D_k, \quad (1.4)$$

for some non-negative integers N_{ij}^k called fusion coefficients. To keep the notation at bay, let us assume that $N_{ij}^k \in \{0, 1\}$. This is the case in the examples we will study below. For a triple $i, j, k \in \mathcal{I}$ such that $N_{ij}^k = 1$, pick a non-zero topological junction field



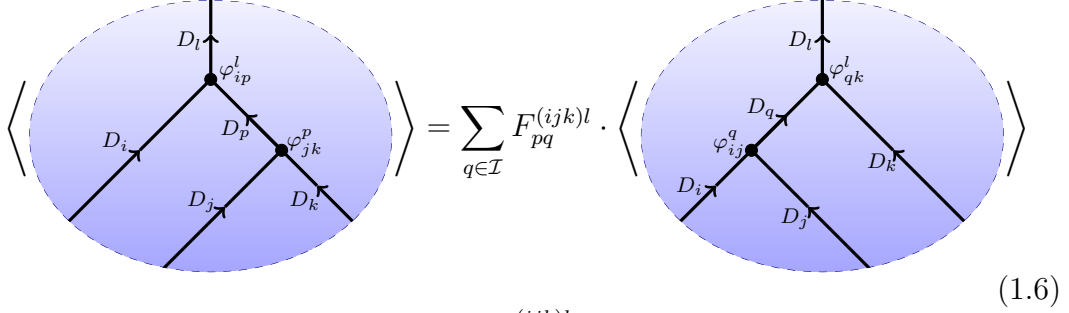
$$\varphi_{ij}^k \in \text{Hom}_{\mathcal{D}}(D_i \otimes D_j, D_k). \quad (1.5)$$

These can then be used to build junctions of higher valency. For example, one can obtain a basis of the space $\text{Hom}_{\mathcal{D}}(D_i \otimes D_j \otimes D_k, D_l)$ by first using the junction φ_{ij}^p , for all p , and then φ_{pk}^l :



¹Of course, one can have more general configurations of defect lines than those where the defect lines pointing towards and away from the junction are grouped together as in (1.2). Because the defect lines have an orientation, there is a conjugation on the defect conditions $X \mapsto X^\vee$, corresponding to orientation reversal, and the space of junction fields in (1.3) is then isomorphic to $\text{Hom}_{\mathcal{D}}(X_1 \otimes \dots \otimes X_m \otimes Y_n^\vee \otimes \dots \otimes Y_1^\vee, \mathbf{1})$. The topological junction fields for a more general arrangement of the orientations of the defect lines are obtained by placing the conjugations $(\cdot)^\vee$ accordingly.

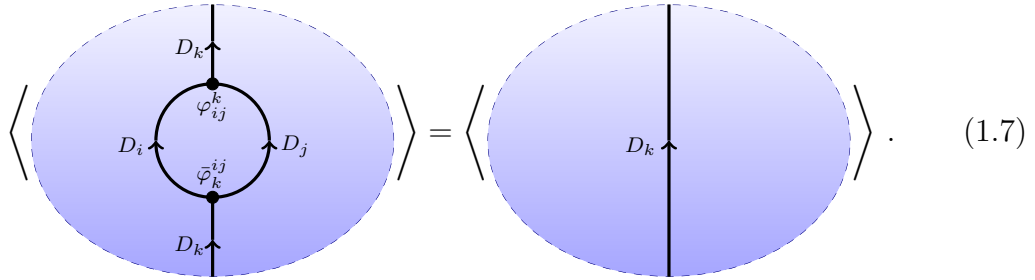
Similarly, one can build a basis by starting with φ_{jk}^q , for all q , and then use φ_{iq}^l . The two bases are related by a basis transformation, i.e. there have to be constants $F_{pq}^{(ijk)l} \in \mathbb{C}$ such that



$$\langle \left(\begin{array}{c} D_l \\ \varphi_{ip}^l \\ \begin{array}{ccc} D_i & & D_p \\ & \nearrow & \searrow \\ & \varphi_{jk}^p & \\ & \nwarrow & \nearrow \\ D_j & & D_k \end{array} \end{array} \right) \rangle = \sum_{q \in \mathcal{I}} F_{pq}^{(ijk)l} \cdot \langle \left(\begin{array}{c} D_l \\ \varphi_{qk}^l \\ \begin{array}{ccc} D_q & & D_k \\ & \nearrow & \searrow \\ & \varphi_{ij}^q & \\ & \nwarrow & \nearrow \\ D_i & & D_j \end{array} \end{array} \right) \rangle \quad (1.6)$$

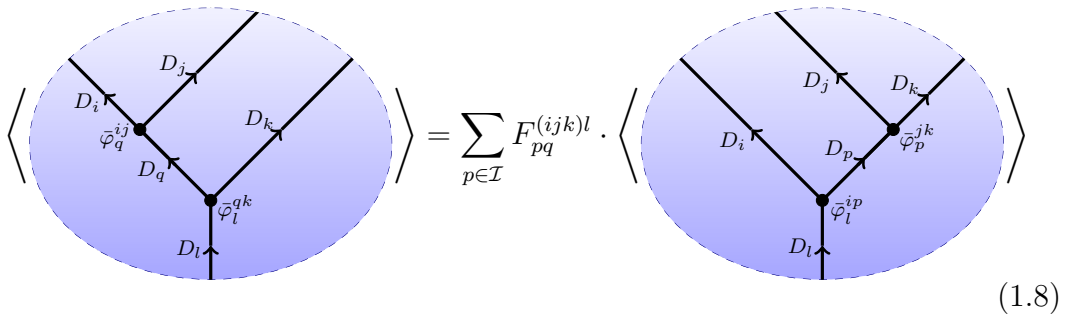
holds for all correlators. The constants $F_{pq}^{(ijk)l}$ are the entries of the so-called *fusing matrices*, and they depend on the choice of 3-valent junction fields in (1.5). When comparing fusing matrices of different realisations of the category \mathcal{D}' , one thus has to take care to compare basis-independent information.

The fusing matrices may just as well be extracted using the opposite orientation on the defect lines. To this end pick junction fields $\bar{\varphi}_k^{ij} \in \text{Hom}_{\mathcal{D}}(D_k, D_i \otimes D_j)$ such that



$$\langle \left(\begin{array}{c} D_k \\ \varphi_{ij}^k \\ \begin{array}{ccc} D_i & & D_j \\ & \nearrow & \searrow \\ & \varphi_k^{ij} & \\ & \nwarrow & \nearrow \\ D_k & & D_k \end{array} \end{array} \right) \rangle = \langle \left(\begin{array}{c} D_k \end{array} \right) \rangle \quad (1.7)$$

The $\bar{\varphi}_k^{ij}$ exist by semi-simplicity. Because the space $\text{Hom}_{\mathcal{D}}(D_k, D_i \otimes D_j)$ is one-dimensional (recall that $N_{ij}^k = 1$), this fixes $\bar{\varphi}_k^{ij}$ uniquely. Then



$$\langle \left(\begin{array}{c} D_l \\ \varphi_l^{qk} \\ \begin{array}{ccc} D_i & & D_k \\ & \nearrow & \searrow \\ & \varphi_q^{ij} & \\ & \nwarrow & \nearrow \\ D_j & & D_l \end{array} \end{array} \right) \rangle = \sum_{p \in \mathcal{I}} F_{pq}^{(ijk)l} \cdot \langle \left(\begin{array}{c} D_l \\ \varphi_l^{ip} \\ \begin{array}{ccc} D_j & & D_k \\ & \nearrow & \searrow \\ & \varphi_p^{jk} & \\ & \nwarrow & \nearrow \\ D_i & & D_l \end{array} \end{array} \right) \rangle \quad (1.8)$$

for the same constants $F_{pq}^{(ijk)l}$. To see this, compose (1.6) from below with an appropriate dual arrangement of defect junctions, and (1.8) from above, to find

in both cases

$$\langle \left(\text{Diamond with junctions } \varphi_{ip}, \varphi_{jk}, \varphi_{ij}, \varphi_{lk} \right) \rangle = F_{pq}^{(ijk)l} \cdot \langle \text{Vertical line } D_l \rangle \quad (1.9)$$

Let us now sketch how the defect category is described in the two models we consider.

Topological defects in rational conformal field theory

Topological defects in rational conformal field theory and their fusion have first been studied in [PZ] and defect junctions have been considered in [FRS4, Fr3]. In the formalism of [FRS1, FRS4, Fr3], one starts from the category \mathcal{C} of representations of a rational vertex operator algebra V (by “rational” we mean that it satisfies the finiteness conditions of [Hu]). In \mathcal{C} one chooses an algebra object A (a special symmetric Frobenius algebra) and this determines an oriented open/closed conformal field theory with left/right symmetry V . The topological defects which are compatible with the entire symmetry V are described by the category $\mathcal{D} = A\text{-mod-}A$ of A -bimodules in \mathcal{C} [FRS1, Fr3]. This is automatically a semi-simple monoidal category. Topological junction fields have left/right conformal weight $(0, 0)$. For a junction as in (1.3), such fields are precisely labelled by elements of the corresponding Hom-space in the bimodule category [FRS4, Fr3].

The case that will be of interest for us is when V is the bosonic subalgebra of the $\mathcal{N} = 2$ super-Virasoro vertex operator algebra with central charge $c = 3 - \frac{6}{d}$ for $d \in \mathbb{Z}_{\geq 1}$. Let us denote its representation category by $\mathcal{C}_d^{\mathcal{N}=2}$. The relevant algebra A and the computation of $A\text{-mod-}A$ are given in appendix A.2. Denote the resulting category of bimodules by $\mathcal{D}_d^{\mathcal{N}=2}$. One finds that in fact

$$\mathcal{D}_d^{\mathcal{N}=2} \cong \mathcal{C}_d^{\mathcal{N}=2} \quad \text{as monoidal categories.} \quad (1.10)$$

The simple objects in $\mathcal{D}_d^{\mathcal{N}=2}$, i. e. the elementary defects, are labelled by triples $[l, m, s]$, where l, m, s are integers subject to conditions (3.1) and (3.2) below. The comparison to the Landau-Ginzburg description will only be carried out in the subcategory of $\mathcal{D}_d^{\mathcal{N}=2}$ which is generated by objects labelled $[l, m, 0]$ (or, equivalently, $[d - 2 - l, m + d, 2]$, see (3.2)) via direct sums. This is a monoidal subcategory, and we denote it by $\mathcal{D}_{d,s=0}^{\mathcal{N}=2}$.

Topological defects in Landau-Ginzburg models

As mentioned above, $\mathcal{N} = 2$ supersymmetric Landau-Ginzburg models are believed to have $\mathcal{N} = 2$ superconformal field theories as their infrared fixed points

under renormalisation group flow. Since the topological twist of an $\mathcal{N} = 2$ superconformal field theory is equivalent to its subsector built on (anti-)chiral primaries (see e.g. [Zh]), topologically twisted Landau-Ginzburg models should be equivalent to the chiral ring of the associated conformal field theory.

Conformal boundary conditions of a topologically B-twisted Landau-Ginzburg model with superpotential W are the objects of the category $\text{MF}^{\text{R}}(W)$ of graded matrix factorisations of W , and its morphisms describe boundary fields. It was argued in [BR] that the same reasoning that leads to this description of the boundary sector equally well applies to defects: between theories with superpotentials W and W' , they are given by suitable matrix factorisations of $W - W'$, where the minus sign is due to the different relative orientation of the defect line as seen from the two theories. Furthermore, the fusion of defects is described by the tensor product of [Yo, ADD, KhR, BR].

Another way to understand the minus sign is the folding trick: as matrix factorisations of W are the same as modules over the curved differential graded algebra $(R, 0, W)$ [KL2], defects between theories with potentials W and W' should be described by $(R, 0, W)$ - $(R', 0, W')$ -bimodules. The latter may also be viewed as modules over $(R, 0, W) \otimes (R, 0, W')^{\text{op}}$ where the opposite curved differential graded algebra is given by $(R, 0, W')^{\text{op}} = (R^{\text{op}}, 0, -W')$, see [KhR, KaR].

Hence we see again that defects are modelled by matrix factorisations of $W - W'$. As in the CFT case, bimodules feature prominently in the description of defects. We shall try to systematically exploit this fact, organising defects and their fields in a category $\text{MF}_{\text{bi}}^{\text{R}}(W, W')$ (to be defined in the next section) that is entirely treated in bimodule language.

Defects in a single Landau-Ginzburg model with superpotential $W = x^d$ are described by $\text{MF}_{\text{bi}}^{\text{R}}(x^d, x^d)$, or, using the equivalence between $\mathbb{C}[x]$ -bimodules and modules over $\mathbb{C}[x] \otimes_{\mathbb{C}} \mathbb{C}[x] \cong \mathbb{C}[x, y]$, by matrix factorisations of the potential $x^d - y^d$. The simplest examples are [ADD, BR]

$$P_S = \begin{pmatrix} 0 & \prod_{i \in S} (x - \eta^i y) \\ \prod_{i \in \{0, \dots, d-1\} \setminus S} (x - \eta^i y) & 0 \end{pmatrix} \quad (1.11)$$

where S is a subset of $\{0, \dots, d-1\}$ and $\eta = e^{2\pi i/d}$. Let us denote by $(\text{MF}_{\text{bi}}^{\text{R}}(W, W'))_0$ the subcategory of $\text{MF}_{\text{bi}}^{\text{R}}(W, W')$ whose morphism spaces are restricted to the R-charge zero subspaces (see definitions 2.12 and 2.14 for details). Let $(\mathcal{P}_d)_0$ be the full subcategory of $(\text{MF}_{\text{bi}}^{\text{R}}(x^d, x^d))_0$ generated via direct sums by all the P_S with subsets S consisting only of consecutive numbers (modulo d); $(\mathcal{P}_d)_0$ is also closed under taking tensor products [BR]. We can now formulate

Conjecture 1.1. $\mathcal{D}_{d,s=0}^{\mathcal{N}=2} \cong (\mathcal{P}_d)_0$ as monoidal categories.

On the level of objects this has been established in [BR]. We will test the conjecture by also comparing morphism spaces and in particular the action of the tensor product on morphisms in certain examples.

The evidence for conjecture 1.1 is gathered by computing the relevant structure on each side separately – in this case certain fusing matrices – and then comparing the results. It would be much preferable to have a more direct and conceptual argument relating the two sides. In fact, since the defects in $\mathcal{D}_{d,s=0}^{\mathcal{N}=2}$ are precisely those which are transparent to $G^\pm \bar{G}^\pm$,² we expect more generally that there is a (non-surjective) monoidal functor from the category of $G^\pm \bar{G}^\pm$ -preserving topological defects in an $\mathcal{N} = 2$ superconformal field theory to $(\text{MF}_{\text{bi}}^{\text{R}}(W, W))_0$ for an appropriate superpotential W .

This paper is organised as follows. In section 2 we define matrix bi-factorisations and recall the definition of a monoidal category. Then we show that matrix bi-factorisations of arbitrary polynomial superpotentials are endowed with a monoidal structure. In section 3 we restrict to the one-variable case and compare the monoidal structure of graded matrix bi-factorisations to that of the associated minimal superconformal field theory. Some details of the proofs of section 2 and the CFT computations of section 3 are postponed to the appendix.

2. Matrix bi-factorisations

In this section we define matrix bi-factorisations of arbitrary superpotentials and show that they form a monoidal category as expected from their physical interpretation as defects in Landau-Ginzburg models. After setting up a suitable bimodule language we discuss the case of superpotentials in only one variable in detail, treating both the situation with and without R-charge. Then we generalise these results to any finite number of variables.

2.1. Preliminaries

Let k be a field, let R_1 and R_2 be commutative k -algebras and let $W_1 \in R_1$, $W_2 \in R_2$ be two elements that we call superpotentials. Every R_1 - R_2 -bimodule M comes equipped with a left action $\rho_M^l : R_1 \times M \rightarrow M$ and a right action $\rho_M^r : M \times R_2 \rightarrow M$. We call an R_1 - R_2 -bimodule free if the corresponding $(R_1 \otimes_k R_2)$ -left module is free.

Definition 2.1. A *matrix bi-factorisation* of W_1 and W_2 is a tuple $D = (D_0, D_1, d_0, d_1)$ where D_i are free R_1 - R_2 -bimodules and $d_0 : D_0 \rightarrow D_1$ and $d_1 : D_1 \rightarrow D_0$ are bimodule maps such that

$$\begin{aligned} (d_1 \circ d_0)(m_0) &= \rho_{D_0}^l(W_1, m_0) - \rho_{D_0}^r(m_0, W_2), \\ (d_0 \circ d_1)(m_1) &= \rho_{D_1}^l(W_1, m_1) - \rho_{D_1}^r(m_1, W_2) \end{aligned}$$

for all $m_i \in D_i$.

² $G^\pm \bar{G}^\pm$ are the four (bosonic) combinations of the holomorphic and anti-holomorphic $\mathcal{N} = 2$ supercurrents.

We will often simply write $D = (D_0, D_1, d_0, d_1)$ in matrix form $\begin{pmatrix} 0 & d_1 \\ d_0 & 0 \end{pmatrix}$ if the source and target modules are clear from the context.

Definition 2.2. The category $\mathrm{DG}_{\mathrm{bi}}(W_1, W_2)$ has matrix bi-factorisations of W_1 and W_2 as objects, and its morphism spaces are \mathbb{Z}_2 -graded: even morphisms $\phi \equiv \begin{pmatrix} \phi_0 & 0 \\ 0 & \phi_1 \end{pmatrix}$ from D to D' are pairs of bimodule maps $\phi_0 : D_0 \rightarrow D'_0$, $\phi_1 : D_1 \rightarrow D'_1$, and odd morphisms $\psi \equiv \begin{pmatrix} 0 & \psi_1 \\ \psi_0 & 0 \end{pmatrix}$ are pairs of bimodule maps $\psi_0 : D_0 \rightarrow D'_1$, $\psi_1 : D_1 \rightarrow D'_0$. The composition in $\mathrm{DG}_{\mathrm{bi}}(W_1, W_2)$ is given by matrix multiplication.

The category $\mathrm{DG}_{\mathrm{bi}}(W_1, W_2)$ is differential \mathbb{Z}_2 -graded with the differential given by

$$\begin{aligned} \delta : \mathrm{Hom}_{\mathrm{DG}_{\mathrm{bi}}(W_1, W_2)}(D, D') &\longrightarrow \mathrm{Hom}_{\mathrm{DG}_{\mathrm{bi}}(W_1, W_2)}(D, D'), \\ \phi &\longmapsto D'\phi - (-1)^{|\phi|}\phi D \end{aligned}$$

for homogeneous ϕ whose \mathbb{Z}_2 -degree $|\phi|$ is 0 or 1 depending on whether ϕ is even or odd as a supermatrix. In particular, one may consider the homotopy category $H_\delta^0(\mathrm{DG}_{\mathrm{bi}}(W_1, W_2))$ whose objects are the same as those of $\mathrm{DG}_{\mathrm{bi}}(W_1, W_2)$ and whose morphism spaces are the zeroth δ -cohomologies of the morphism spaces of $\mathrm{DG}_{\mathrm{bi}}(W_1, W_2)$.

Definition 2.3. The *category of matrix bi-factorisations of W_1 and W_2* is the homotopy category

$$\mathrm{MF}_{\mathrm{bi}}(W_1, W_2) = H_\delta^0(\mathrm{DG}_{\mathrm{bi}}(W_1, W_2)). \quad (2.1)$$

Concretely this means that morphisms from D to D' in $\mathrm{MF}_{\mathrm{bi}}(W_1, W_2)$ are even matrices ϕ satisfying $D'\phi = \phi D$ modulo matrices of the form $D'\psi + \psi D$ for any odd matrix ψ .

Given an R_1 - R_2 -bimodule M we can construct a new bimodule ${}_{\sigma_1}M_{\sigma_2}$ as follows. Let $\sigma_i \in \mathrm{Aut}(R_i)$ be automorphisms of R_i . Then as sets we have ${}_{\sigma_1}M_{\sigma_2} = M$, but the left and right actions are twisted via σ_1 and σ_2 , respectively:

$$\rho_{{}_{\sigma_1}M_{\sigma_2}}^l(r_1, m) = \rho_M^l(\sigma_1(r_1), m), \quad \rho_{{}_{\sigma_1}M_{\sigma_2}}^r(m, r_2) = \rho_M^r(m, \sigma_2(r_2)) \quad (2.2)$$

for all $r_i \in R_i$ and $m \in M$. We note that one can canonically identify

$$\mathrm{Hom}_{R_1\text{-mod-}R_2}(M, M') = \mathrm{Hom}_{R_1\text{-mod-}R_2}({}_{\sigma_1}M_{\sigma_2}, {}_{\sigma_1}M'_{\sigma_2}) \quad (2.3)$$

as the underlying maps of sets are the same. Hence for every matrix bi-factorisation D of W_1 and W_2 and automorphisms σ_i as above, it follows that

$${}_{\sigma_1}D_{\sigma_2} = ({}_{\sigma_1}(D_0)_{\sigma_2}, {}_{\sigma_1}(D_1)_{\sigma_2}, d_0, d_1) \quad (2.4)$$

is a matrix bi-factorisation of $\sigma_1(W_1)$ and $\sigma_2(W_2)$. In particular, if the automorphisms are symmetries of the superpotentials, i.e. $\sigma_1(W_1) = W_1$ and $\sigma_2(W_2) = W_2$, then ${}_{\sigma_1}D_{\sigma_2} \in \mathrm{MF}_{\mathrm{bi}}(W_1, W_2)$.

2.2. Tensor structure for one variable

We wish matrix bi-factorisations to describe defects in a given topological Landau-Ginzburg model with superpotential $W \in R := \mathbb{C}[x]$. Thus we restrict to the case $R_1 = R_2 = R$ with $W_1 = W_2 = W$ and consider matrix bi-factorisations of W and W , or simply of W for short. We also write $\text{MF}_{\text{bi}}(W) := \text{MF}_{\text{bi}}(W, W)$.

Under the identification of R -bimodules with $(R \otimes_{\mathbb{C}} R)$ -left modules, a bimodule describing a defect condition corresponds to a left module describing a boundary condition in the folded model, see e.g. [WA]. Because of the latter's geometric origin as superbundles [KL1, La] we also only consider free bimodules in the case of defects.

The main goal of this section is to prove that $\text{MF}_{\text{bi}}(W)$ is a monoidal category, hence we now recall the definition of the latter.

Definition 2.4. A *monoidal category* is a tuple $(\mathcal{M}, \otimes, \alpha, I, \lambda, \rho)$ such that

- \mathcal{M} is a category,
- $\otimes : \mathcal{M} \times \mathcal{M} \rightarrow \mathcal{M}$ is a bifunctor,
- $\alpha_{A,B,C} : (A \otimes B) \otimes C \rightarrow A \otimes (B \otimes C)$ are isomorphisms, natural in A, B, C ,
- $I \in \mathcal{M}$ is called the *unit object*,
- $\lambda_A : I \otimes A \rightarrow A$, $\rho_A : A \otimes I \rightarrow A$ are isomorphisms, natural in A ,
- the *triangle axiom* holds, i. e. the diagramme

$$\begin{array}{ccc}
 (A \otimes I) \otimes B & \xrightarrow{\alpha_{A,I,B}} & A \otimes (I \otimes B) \\
 \searrow \rho_A \otimes \text{id}_B & & \swarrow \text{id}_A \otimes \lambda_B \\
 & A \otimes B &
 \end{array} \tag{2.5}$$

commutes for all $A, B \in \mathcal{M}$,

- the *pentagon axiom* holds, i. e. the diagramme

$$\begin{array}{ccccc}
 & & (A \otimes (B \otimes C)) \otimes D & & \\
 & \nearrow \alpha_{A,B,C} \otimes \text{id}_D & & \searrow \alpha_{A,B \otimes C,D} & \\
 ((A \otimes B) \otimes C) \otimes D & & & & A \otimes ((B \otimes C) \otimes D) \\
 \downarrow \alpha_{A \otimes B,C,D} & & & & \downarrow \text{id}_A \otimes \alpha_{B,C,D} \\
 (A \otimes B) \otimes (C \otimes D) & \xrightarrow{\alpha_{A,B,C \otimes D}} & & & A \otimes (B \otimes (C \otimes D))
 \end{array} \tag{2.6}$$

commutes for all $A, B, C, D \in \mathcal{M}$.

The tensor product \otimes of $\text{MF}_{\text{bi}}(W)$ is given by [Yo, ADD, KhR]

$$D \otimes D' = \left(D_0 \otimes_R D'_0 \oplus D_1 \otimes_R D'_1, D_1 \otimes_R D'_0 \oplus D_0 \otimes_R D'_1, \right. \\ \left. \begin{pmatrix} d_0 \otimes_R \mathbb{1}_{D'_0} & -\mathbb{1}_{D_1} \otimes_R d'_1 \\ \mathbb{1}_{D_0} \otimes_R d'_0 & d_1 \otimes_R \mathbb{1}_{D'_1} \end{pmatrix}, \begin{pmatrix} d_1 \otimes_R \mathbb{1}_{D'_0} & \mathbb{1}_{D_0} \otimes_R d'_1 \\ -\mathbb{1}_{D_1} \otimes_R d'_0 & d_0 \otimes_R \mathbb{1}_{D'_1} \end{pmatrix} \right), \quad (2.7)$$

$$\phi \otimes \phi' = \begin{pmatrix} \phi_0 \otimes_R \phi'_0 & 0 & 0 & 0 \\ 0 & \phi_1 \otimes_R \phi'_1 & 0 & 0 \\ 0 & 0 & \phi_1 \otimes_R \phi'_0 & 0 \\ 0 & 0 & 0 & \phi_0 \otimes_R \phi'_1 \end{pmatrix}, \quad (2.8)$$

and it is also well-defined in $\text{DG}_{\text{bi}}^0(W)$. If we view matrix bi-factorisations as matrices then this definition of the tensor product is the natural one as it lets matrix bi-factorisations act as a derivation:

$$(D \otimes D')(m \otimes_R m') = D(m) \otimes_R m' + (-1)^{|m|} m \otimes_R D'(m') \quad (2.9)$$

for all homogeneous $m \in D_0 \oplus D_1$ and $m' \in D'_0 \oplus D'_1$.

For three matrix bi-factorisations D, D', D'' of W we have an isomorphism

$$\alpha_{D,D',D''} : (D \otimes D') \otimes D'' \longrightarrow D \otimes (D' \otimes D'') \quad (2.10)$$

in $\text{MF}_{\text{bi}}(W)$ with

$$(\alpha_{D,D',D''})_0 = \begin{pmatrix} \mathbb{1}_{D_0 \otimes_R D'_0 \otimes_R D''_0} & 0 & 0 & 0 \\ 0 & 0 & 0 & \mathbb{1}_{D_0 \otimes_R D'_1 \otimes_R D''_1} \\ 0 & \mathbb{1}_{D_1 \otimes_R D'_1 \otimes_R D''_0} & 0 & 0 \\ 0 & 0 & \mathbb{1}_{D_1 \otimes_R D'_0 \otimes_R D''_1} & 0 \end{pmatrix}, \quad (2.11a)$$

$$(\alpha_{D,D',D''})_1 = \begin{pmatrix} \mathbb{1}_{D_1 \otimes_R D'_0 \otimes_R D''_0} & 0 & 0 & 0 \\ 0 & 0 & 0 & \mathbb{1}_{D_1 \otimes_R D'_1 \otimes_R D''_1} \\ 0 & \mathbb{1}_{D_0 \otimes_R D'_1 \otimes_R D''_0} & 0 & 0 \\ 0 & 0 & \mathbb{1}_{D_0 \otimes_R D'_0 \otimes_R D''_1} & 0 \end{pmatrix}. \quad (2.11b)$$

Lemma 2.5. The maps $\alpha_{D,D',D''}$, viewed as morphisms in $\text{DG}_{\text{bi}}(W)$, satisfy the pentagon axiom as the associator for \otimes .

Proof. Matrix multiplication. □

We also need to identify a unit object in $\text{MF}_{\text{bi}}(W)$. To do so, we first introduce some notation to calculate with free R -bimodules. Every free R -bimodule is isomorphic to $R \otimes_{\mathbb{C}} U \otimes_{\mathbb{C}} R$ for some \mathbb{C} -vector space U . For two vector spaces U, V consider a linear map $\phi \in \text{Hom}(U, V[a, b])$, where $V[a, b]$ are polynomials in

two formal variables a, b with coefficients in V . We can write $\phi = \sum_{m,n} \phi_{mn} a^m b^n$, where $\phi_{mn} \in \text{Hom}(U, V)$. The sum may be infinite, but for a given $u \in U$ only a finite number of $\phi_{mn}(u)$ will be non-zero. From ϕ we obtain an R -bimodule map by setting

$$\begin{aligned} \hat{\phi} : R \otimes_{\mathbb{C}} U \otimes_{\mathbb{C}} R &\longrightarrow R \otimes_{\mathbb{C}} V \otimes_{\mathbb{C}} R, \\ r \otimes_{\mathbb{C}} U \otimes_{\mathbb{C}} s &\longmapsto \sum_{m,n} r x^m \otimes_{\mathbb{C}} \phi_{mn}(u) \otimes_{\mathbb{C}} x^n s. \end{aligned} \quad (2.12)$$

This gives an isomorphism

$$\text{Hom}(U, V[a, b]) \xrightarrow{\sim} \text{Hom}_{R\text{-mod-}R}(R \otimes_{\mathbb{C}} U \otimes_{\mathbb{C}} R, R \otimes_{\mathbb{C}} V \otimes_{\mathbb{C}} R). \quad (2.13)$$

We denote its inverse by $(\check{})$, i. e. for a given R -bimodule map $\xi : R \otimes_{\mathbb{C}} U \otimes_{\mathbb{C}} R \rightarrow R \otimes_{\mathbb{C}} V \otimes_{\mathbb{C}} R$ we have $\xi = [\check{\xi}(a, b)]^{\wedge}$. For long expressions we write $\widehat{\text{long}} = [\text{long}]^{\wedge}$. For example, $[a - b]^{\wedge} : R \otimes_{\mathbb{C}} R \rightarrow R \otimes_{\mathbb{C}} R$ (here $U = V = \mathbb{C}$) is the map $r \otimes_{\mathbb{C}} s \mapsto xr \otimes_{\mathbb{C}} s - r \otimes_{\mathbb{C}} xs$, and $[a^i x^j b^k]^{\wedge}$ from $R \otimes_{\mathbb{C}} R \otimes_{\mathbb{C}} R$ to itself (here $U = V = R$) is the map $r \otimes_{\mathbb{C}} s \otimes_{\mathbb{C}} t \mapsto x^i r \otimes_{\mathbb{C}} x^j s \otimes_{\mathbb{C}} x^k t$. We define the composition of $\phi \in \text{Hom}(U, V[a, b])$ and $\psi \in \text{Hom}(T, U[a, b])$ to be

$$\phi \circ \psi = \sum_{k,l,m,n} \phi_{kl} \circ \psi_{mn} a^{k+m} b^{l+n} \in \text{Hom}(T, V[a, b]). \quad (2.14)$$

With this definition one has $\hat{\phi} \circ \hat{\psi} = [\phi \circ \psi]^{\wedge}$. For example, for $T = U = V = \mathbb{C}$, $\phi(a, b)$ and $\psi(a, b)$ are just polynomials, and composition amounts to multiplying these polynomials, $\hat{\phi} \circ \hat{\psi} = [\phi(a, b) \cdot \psi(a, b)]^{\wedge}$.

We are now in a position to define the matrix bi-factorisation

$$I = (R \otimes_{\mathbb{C}} R, R \otimes_{\mathbb{C}} R, \iota_0, \iota_1) \quad (2.15a)$$

where

$$\iota_0 = [(W(a) - W(b))/(a - b)]^{\wedge}, \quad \iota_1 = [a - b]^{\wedge} \quad (2.15b)$$

as in [KhR, KaR, BR]. We will see that I is the unit object in $\text{MF}_{\text{bi}}(W)$.

Let $\mu : R \otimes_{\mathbb{C}} R \rightarrow R$ be the multiplication map, i. e. $\mu(r \otimes_{\mathbb{C}} s) = rs$ for all $r, s \in R$. Then it is easy to check that

$$\mu \circ \iota_0 = \partial W \mu, \quad \mu \circ \iota_1 = 0. \quad (2.16)$$

For example, $(\mu \circ [a - b]^{\wedge})(r \otimes_{\mathbb{C}} s) = \mu(rx \otimes_{\mathbb{C}} s - r \otimes_{\mathbb{C}} xs) = rxs - rxs = 0$, and similarly for the other identity.

For any $D \in \text{MF}_{\text{bi}}(W)$ we have maps

$$\lambda_D = \begin{pmatrix} \mu \otimes_R \mathbb{1}_{D_0} & 0 & 0 & 0 \\ 0 & 0 & 0 & \mu \otimes_R \mathbb{1}_{D_1} \end{pmatrix} : I \otimes D \longrightarrow D, \quad (2.17a)$$

$$\rho_D = \begin{pmatrix} \mathbb{1}_{D_0} \otimes_R \mu & 0 & 0 & 0 \\ 0 & 0 & \mathbb{1}_{D_1} \otimes_R \mu & 0 \end{pmatrix} : D \otimes I \longrightarrow D. \quad (2.17b)$$

Since we have $I_0 = I_1 = R \otimes_{\mathbb{C}} \mathbb{C} \otimes_{\mathbb{C}} R \equiv R \otimes_{\mathbb{C}} R$, the multiplication μ can act on I_0, I_1 . In a different guise and in another context such maps have already appeared in [KhR, props. 15 & 17].

To see that λ_D, ρ_D are indeed morphisms in $\text{MF}_{\text{bi}}(W)$, we have to check that

$$D\lambda_D = \lambda_D(I \otimes D), \quad D\rho_D = \rho_D(D \otimes I). \quad (2.18)$$

In the case of the left unit map λ_D we have according to (2.7) and (2.15):

$$I \otimes D = \begin{pmatrix} 0 & 0 & \iota_1 \otimes_R \mathbb{1}_{D_0} & \mathbb{1}_{I_0} \otimes_R d_1 \\ 0 & 0 & -\mathbb{1}_{I_1} \otimes_R d_0 & \iota_0 \otimes_R \mathbb{1}_{D_1} \\ \iota_0 \otimes_R \mathbb{1}_{D_0} & -\mathbb{1}_{I_1} \otimes_R d_1 & 0 & 0 \\ \mathbb{1}_{I_0} \otimes_R d_0 & \iota_1 \otimes_R \mathbb{1}_{D_1} & 0 & 0 \end{pmatrix} \quad (2.19)$$

and thus

$$\begin{aligned} D\lambda_D &= \begin{pmatrix} 0 & 0 & 0 & d_1 \circ (\mu \otimes_R \mathbb{1}_{D_1}) \\ d_0 \circ (\mu \otimes_R \mathbb{1}_{D_0}) & 0 & 0 & 0 \end{pmatrix}, \\ (\lambda_D(I \otimes D))_0 &= \left((\mu \otimes_R \mathbb{1}_{D_1}) \circ (\mathbb{1}_{I_0} \otimes_R d_0) \quad (\mu \otimes_R \mathbb{1}_{D_1}) \circ (\iota_1 \otimes_R \mathbb{1}_{D_1}) \right) \\ (\lambda_D(I \otimes D))_1 &= \left((\mu \otimes_R \mathbb{1}_{D_0}) \circ (\iota_1 \otimes_R \mathbb{1}_{D_0}) \quad (\mu \otimes_R \mathbb{1}_{D_0}) \circ (\mathbb{1}_{I_0} \otimes_R d_1) \right). \end{aligned}$$

Using the second relation of (2.16) on these expressions one immediately verifies that $D\lambda_D = \lambda_D(I \otimes D)$ and hence λ_D is a morphism. Similarly, one also finds that ρ_D is a morphism.

We will now show that (I, λ, ρ) provide $(\text{MF}_{\text{bi}}(W), \otimes, \alpha)$ with the remaining structure of a monoidal category. This means that we need to prove that λ_D and ρ_D are isomorphisms, that α, λ, ρ are natural isomorphisms, and that they satisfy the triangle axiom. The latter is straightforward:

Lemma 2.6. For all $D, D' \in \text{MF}_{\text{bi}}(W)$, the diagramme

$$\begin{array}{ccc} (D \otimes I) \otimes D' & \xrightarrow{\alpha_{D, I, D'}} & D \otimes (I \otimes D') \\ \rho_D \otimes \text{id}_{D'} \searrow & & \swarrow \text{id}_D \otimes \lambda_{D'} \\ & D \otimes D' & \end{array} \quad (2.20)$$

commutes (already on the DG level, i. e. in $\text{DG}_{\text{bi}}(W)$).

Proof. This follows from the definitions (2.8), (2.11), (2.17) and explicit matrix multiplication. \square

However, λ and ρ do not give isomorphisms in $\text{DG}_{\text{bi}}(W)$ and we have to pass to the cohomology level.

Lemma 2.7. In $\text{MF}_{\text{bi}}(W)$, λ_D and ρ_D are isomorphisms for all matrix bi-factorisations D , and we have $\lambda_I = \rho_I$.

The construction of inverses to λ and ρ in $\text{MF}_{\text{bi}}(W)$ is rather tedious and we relegate the proof of the above lemma to appendix A.1, where we also specify the associated homotopies.

Theorem 2.8. The category $\text{MF}_{\text{bi}}(W)$ together with the structure $(\otimes, \alpha, I, \lambda, \rho)$ is monoidal.

Proof. The naturality of α , λ and ρ and the functoriality of \otimes already hold in $\text{DG}_{\text{bi}}(W)$ as is easily checked from the definitions via direct computation. Then the theorem follows from lemmas 2.5, 2.6 and 2.7. \square

We remark that one may also consider the full subcategory $\text{MF}_{\text{bi}}(W)_{\text{f}}$ of $\text{MF}_{\text{bi}}(W)$ consisting of all matrix bi-factorisations $D \in \text{MF}_{\text{bi}}(W)$ which are isomorphic (in $\text{MF}_{\text{bi}}(W)$) to a matrix bi-factorisation D' for which the R -bimodules D'_0 and D'_1 have finite rank. As shown in [KhR, prop. 13] and [BR, sec. 4.2] this subcategory is closed under taking tensor products, and hence the above theorem also holds for $\text{MF}_{\text{bi}}(W)_{\text{f}}$.

Lemma 2.9 ([KhR, KaR, BR]). $\text{Hom}_{\text{MF}_{\text{bi}}(W)}(I, I) \cong R/(\partial W)$.

Proof. Let ϕ be a representative of an element in $\text{Hom}_{\text{MF}_{\text{bi}}(W)}(I, I)$. We can write $\phi = \begin{pmatrix} \hat{\phi}_0 & 0 \\ 0 & \hat{\phi}_1 \end{pmatrix}$ where according to the relation (2.12), $\phi_0, \phi_1 \in \text{Hom}(\mathbb{C}, \mathbb{C}[a, b])$ are polynomials in the formal variables a, b . From the closedness condition of ϕ , i. e. $I\phi = \phi I$, we find that $(a-b)\phi_1 = \phi_0(a-b)$ and hence we must have $\phi_0 = \phi_1$. The condition $\iota_1 \circ \hat{\phi}_1 = \hat{\phi}_0 \circ \iota_1$ is then automatic, so that the space of closed elements in $\text{Hom}_{\text{DG}_{\text{bi}}^0(W)}(I, I)$ is isomorphic to $\mathbb{C}[a, b]$.

Exact elements are of the form $I\psi + \psi I$ for arbitrary degree-one morphisms $\psi = \begin{pmatrix} 0 & \hat{\psi}_1 \\ \hat{\psi}_0 & 0 \end{pmatrix} \in \text{Hom}_{\text{DG}_{\text{bi}}^1(W)}(I, I)$, and this is equal to

$$\begin{pmatrix} \hat{\phi}'_0 & 0 \\ 0 & \hat{\phi}'_0 \end{pmatrix}, \quad \phi'_0 = (a-b)\psi_0 + \psi_1 \frac{W(a) - W(b)}{a-b}. \quad (2.21)$$

So we see that the space of exact elements in $\text{Hom}_{\text{DG}_{\text{bi}}^0(W)}(I, I)$ is isomorphic to the ideal $(a-b, (W(a) - W(b))/(a-b))$, and hence $\text{Hom}_{\text{MF}_{\text{bi}}(W)}(I, I) \cong R/(\partial W)$. \square

Remark. The physical interpretation of this is that the space of defect fields on the identity defect is the same as the space of bulk fields in the full CFT, and that this in particular also holds for the ring of chiral primaries. The latter is equivalent to the closed sector of the associated topologically twisted Landau-Ginzburg model and hence given by the Jacobi ring $R/(\partial W)$ [Va].

We close this subsection with another look at automorphisms which allows us to introduce the matrix bi-factorisations that will be of interest in section 3. Any automorphism of $R = \mathbb{C}[x]$ is fixed by its action on x , so we have a family $\{\sigma_\alpha\}_{\alpha \in \mathbb{C}} \subset \text{Aut}(R)$ given by

$$\sigma_\alpha(x) = e^{iq_x \alpha} x \quad (2.22)$$

where the constant q_x is conventional. The infrared fixed point of a Landau-Ginzburg model with superpotential $W = x^d + \sum_{i>d} c_i x^i$ is the same as that of a model with superpotential x^d . Accordingly, the categories $\text{MF}_{\text{bi}}(W)$ and $\text{MF}_{\text{bi}}(x^d)$ are equivalent, and we may restrict to superpotentials of the form $W = x^d$. In this case we set

$$q_x = \frac{2}{d} \quad (2.23)$$

such that $\sigma_\alpha(W) = e^{2i\alpha} W$.

For a subset $S \subset \{0, \dots, d-1\}$ we define

$$p_S(a, b) = \prod_{i \in S} (a - \eta^i b), \quad \eta = e^{2\pi i/d}. \quad (2.24)$$

Then we get a matrix bi-factorisation

$$P_S = (R \otimes_{\mathbb{C}} R, R \otimes_{\mathbb{C}} R, \hat{p}_{\{0, \dots, d-1\} \setminus S}, \hat{p}_S) \equiv \begin{pmatrix} 0 & \hat{p}_S \\ \hat{p}_{\{0, \dots, d-1\} \setminus S} & 0 \end{pmatrix} \quad (2.25)$$

as in [ADD, BR].

In the above formulation the unit object I is given by $P_{\{0\}}$. Furthermore, we note that $\sigma_{\pi m}(W) = W$ and $\sigma_{\pi m} = \sigma_{\pi(m+d)}$ for all $m \in \mathbb{Z}$.

Lemma 2.10. $\sigma_{\pi m} I_{\sigma_{\pi n}} \cong P_{\{m-n\}}$ in $\text{DG}_{\text{bi}}(x^d)$.

Proof. Recalling the notation introduced at the end of subsection 2.1, this can be seen by direct computation: An isomorphism $\phi \equiv \begin{pmatrix} \phi_0 & 0 \\ 0 & \phi_1 \end{pmatrix} : \sigma_{\pi m} I_{\sigma_{\pi n}} \rightarrow P_{\{m-n\}}$ is provided by $\phi_0(r \otimes_{\mathbb{C}} s) = \sigma_{-\pi m}(r) \otimes_{\mathbb{C}} \sigma_{-\pi n}(s)$ and $\phi_1(r \otimes_{\mathbb{C}} s) = \eta^{-m} \sigma_{-\pi m}(r) \otimes_{\mathbb{C}} \sigma_{-\pi n}(s)$. \square

2.3. R-charge

To compare the defect category $\text{MF}_{\text{bi}}(W)$ with the CFT description it is crucial to have additional information on the R-charges of morphisms. We will now formulate the notion of R-charge in bimodule language and indicate its relation to the standard formulation. In the following we abbreviate ${}_\alpha M_\alpha \equiv {}_{\sigma_\alpha} M_{\sigma_\alpha}$.

Definition 2.11. A *bimodule with $\mathfrak{u}(1)$ -action* is a pair (M, φ) where M is an R -bimodule and φ assigns to each $\alpha \in \mathbb{C}$ a bimodule isomorphism $\varphi_\alpha : M \rightarrow {}_\alpha M_\alpha$ such that $\varphi_{\alpha+\beta} = \varphi_\beta \circ \varphi_\alpha$.

Here, φ_β above is understood as an element of $\text{Hom}({}_\alpha M_\alpha, {}_{\alpha+\beta} M_{\alpha+\beta}) = \text{Hom}(M, {}_\beta M_\beta)$ via the identification (2.3). Similarly, in the next definition the map F is viewed as an element of both $\text{Hom}(M, M')$ and $\text{Hom}({}_\alpha M_\alpha, {}_\alpha M'_\alpha)$.

Definition 2.12. Let (M, φ) and (M', φ') be bimodules with $\mathfrak{u}(1)$ -action and let $F : M \rightarrow M'$ be a bimodule map. We define

$$F_\alpha = (\varphi'_\alpha)^{-1} \circ F \circ \varphi_\alpha, \quad (2.26)$$

and we say that F has *R-charge* q if $F_\alpha = e^{-i\alpha q} F$ for all $\alpha \in \mathbb{C}$.

The minus sign in $e^{-i\alpha q}$ is used to match the standard conventions in the literature as we will see below.

Definition 2.13. A *graded matrix bi-factorisation* is a tuple

$$D = (D_0, D_1, d_0, d_1, \varphi^{D_0}, \varphi^{D_1}) \quad (2.27)$$

such that

- (D_0, D_1, d_0, d_1) is a matrix bi-factorisation,
- (D_0, φ^{D_0}) and (D_1, φ^{D_1}) are bimodules with $\mathfrak{u}(1)$ -action,
- d_0 and d_1 have R-charge 1.

To make contact with the standard formulation, we first define for any \mathbb{C} -vector space V the bimodule maps

$$\begin{aligned} s_\alpha^V : R \otimes_{\mathbb{C}} V \otimes_{\mathbb{C}} R &\longrightarrow {}_\alpha(R \otimes_{\mathbb{C}} V \otimes_{\mathbb{C}} R)_\alpha, \\ r \otimes_{\mathbb{C}} v \otimes_{\mathbb{C}} s &\longmapsto \sigma_\alpha(r) \otimes_{\mathbb{C}} v \otimes_{\mathbb{C}} \sigma_\alpha(s). \end{aligned} \quad (2.28)$$

Then for a matrix bi-factorisation D with $D_0 = R \otimes_{\mathbb{C}} \check{D}_0 \otimes_{\mathbb{C}} R$ and $D_1 = R \otimes_{\mathbb{C}} \check{D}_1 \otimes_{\mathbb{C}} R$ for some vector spaces \check{D}_0, \check{D}_1 , we set

$$U_{D_i}(\alpha) = (\varphi_{-\alpha}^{D_i})^{-1} \circ s_{-\alpha}^{\check{D}_i}, \quad U_D(\alpha) = \begin{pmatrix} U_{D_0}(\alpha) & 0 \\ 0 & U_{D_1}(\alpha) \end{pmatrix}. \quad (2.29)$$

Now let $\hat{\phi} : D \rightarrow D'$ be a morphism of graded matrix bi-factorisations. It follows from the above and equation (2.26) that the condition of $\hat{\phi}$ having R-charge q is equivalent to

$$U_{D'}(\alpha) [\phi(\sigma_\alpha(a), \sigma_\alpha(b))]^\wedge (U_D(\alpha))^{-1} = e^{+iq\alpha} [\phi(a, b)]^\wedge \quad (2.30)$$

for all $\alpha \in \mathbb{C}$, which is the standard R-charge condition.

Definition 2.14. The *category of graded matrix bi-factorisations* $\text{MF}_{\text{bi}}^{\text{R}}(W)$ has graded matrix bi-factorisations as objects and the same morphisms as $\text{MF}_{\text{bi}}(W)$.

The tensor product of $\mathrm{MF}_{\mathrm{bi}}^{\mathrm{R}}(W)$ is the same as that of $\mathrm{MF}_{\mathrm{bi}}(W)$ apart from the fact that we have to specify what \otimes does to the $\mathfrak{u}(1)$ -action: for the tensor product $(D_0, D_1, d_0, d_1, \varphi^{D_0}, \varphi^{D_1}) \otimes (D'_0, D'_1, d'_0, d'_1, \varphi^{D'_0}, \varphi^{D'_1})$ of two graded matrix bi-factorisations it is given by

$$((\varphi \otimes \varphi')_0, (\varphi \otimes \varphi')_1) \quad \text{with} \quad \varphi = \begin{pmatrix} \varphi^{D_0} & 0 \\ 0 & \varphi^{D_1} \end{pmatrix}, \quad \varphi' = \begin{pmatrix} \varphi^{D'_0} & 0 \\ 0 & \varphi^{D'_1} \end{pmatrix} \quad (2.31)$$

where we use (2.8) to compute $\varphi \otimes \varphi'$.

We also have to specify the $\mathfrak{u}(1)$ -action of the unit object. It is given by

$$U_I(\alpha) = \begin{pmatrix} 1 & 0 \\ 0 & e^{i\alpha(q_x-1)} \end{pmatrix}. \quad (2.32)$$

We will see in section 3 that this is a special case of the general R-matrices U_{P_S} which we will infer from comparison with the conformal field theory description.

It is easy to compute the charges of the spectrum of the identity defect:

Lemma 2.15. $\mathrm{Hom}_{\mathrm{MF}_{\mathrm{bi}}^{\mathrm{R}}(x^d)}(I, I) = \mathbb{C}\{\phi_I^i\}_{i \in \{0, \dots, d-2\}}$, where $\phi_I^i = \begin{pmatrix} \hat{a}^i & 0 \\ 0 & \hat{a}^i \end{pmatrix}$ has R-charge $2i/d$.

Proof. This follows as in lemma 2.9 and from a direct computation. \square

Similarly, the associator and the left and right unit isomorphisms have zero R-charge as expected:

Lemma 2.16. $\alpha_{D, D', D''}$, λ_D and ρ_D have R-charge 0 for all D, D', D'' in $\mathrm{MF}_{\mathrm{bi}}^{\mathrm{R}}(W)$.

Proof. Direct computation. \square

The arguments of the previous subsection immediately carry over to the graded situation and we have:

Theorem 2.17. $(\mathrm{MF}_{\mathrm{bi}}^{\mathrm{R}}(W), \otimes, \alpha, (I, U_I), \lambda, \rho)$ is a monoidal category.

2.4. Tensor structure for many variables

We shall now generalise the above results to the many-variable case, i. e. $W \in R = \mathbb{C}[x_1, \dots, x_N]$ for arbitrary $N \in \mathbb{Z}_{\geq 1}$. We will use the obvious generalisation of the “[\cdot]^-calculus” introduced in subsection 2.2 to the case of many formal variables a_i, b_i . Also, in the graded case the superpotential W has to be quasi-homogeneous of degree 2.

Matrix bi-factorisations and their morphisms are defined as before, and also the definitions of the tensor product \otimes and the associator α remain unchanged. The unit object is slightly more involved: as in [KhR] we set

$$\Delta_i(W) = \frac{W(\mathbf{b}_{[i-1]}\mathbf{a}) - W(\mathbf{b}_{[i]}\mathbf{a})}{a_i - b_i}, \quad \mathbf{b}_{[i]}\mathbf{a} = (b_1, \dots, b_i, a_{i+1}, \dots, a_N) \quad (2.33)$$

for all $i \in \{1, \dots, N\}$. Then the unit object is given by

$$I = \left((R \otimes_{\mathbb{C}} R)^{\oplus 2^{N-1}}, (R \otimes_{\mathbb{C}} R)^{\oplus 2^{N-1}}, \iota_0, \iota_1 \right) \quad (2.34)$$

where in matrix form we have

$$\begin{pmatrix} 0 & \iota_1 \\ \iota_0 & 0 \end{pmatrix} = \bigotimes_{i=1}^N \begin{pmatrix} 0 & [a_i - b_i]^\wedge \\ [\Delta_i(W)]^\wedge & 0 \end{pmatrix}. \quad (2.35)$$

Here, the operation \otimes' is defined like the tensor product in (2.7) apart from the fact that “ $\mathbb{1} \otimes_R$ ” and “ $\otimes_R \mathbb{1}$ ” do not appear in \otimes' . For example, we have

$$I = \begin{pmatrix} 0 & 0 & [a_1 - b_1]^\wedge & [a_2 - b_2]^\wedge \\ 0 & 0 & -[\Delta_2(W)]^\wedge & [\Delta_1(W)]^\wedge \\ [\Delta_1(W)]^\wedge & -[a_2 - b_2]^\wedge & 0 & 0 \\ [\Delta_2(W)]^\wedge & [a_1 - b_1]^\wedge & 0 & 0 \end{pmatrix} \quad \text{for } N = 2. \quad (2.36)$$

For $N > 2$, the N -fold product is to be bracketed from the left, i. e., denoting the i -th factor in (2.35) by M_i , the expression stands for $(\dots ((M_1 \otimes' M_2) \otimes' M_3) \otimes' \dots) \otimes' M_N$. The R-charge matrix U_I of I is given by the tensor product of the corresponding one-variable matrices.

The left and right unit isomorphisms

$$\lambda_D : I \otimes D \longrightarrow D, \quad \rho_D : D \otimes I \longrightarrow D \quad (2.37)$$

for a matrix bi-factorisation D are in general represented by $(2 \times 2^{N+1})$ -block matrices. Explicitly, they are given via

$$(\lambda_D)_{i,j} = (\mu \otimes_R \mathbb{1}_{D_0}) \delta_{i,1} \delta_{j,1} + (\mu \otimes_R \mathbb{1}_{D_1}) \delta_{i,2} \delta_{j,3 \cdot 2^{N-1} + 1}, \quad (2.38a)$$

$$(\rho_D)_{i,j} = (\mathbb{1}_{D_0} \otimes_R \mu) \delta_{i,1} \delta_{j,1} + (\mathbb{1}_{D_1} \otimes_R \mu) \delta_{i,2} \delta_{j,2^N + 1}. \quad (2.38b)$$

We recall from subsection 2.2 that $\text{MF}_{\text{bi}}(W)_f$ is the full subcategory of $\text{MF}_{\text{bi}}(W)$ consisting of all matrix bi-factorisations which are isomorphic (in $\text{MF}_{\text{bi}}(W)$) to matrix factorisations with finite-rank R -bimodules. The category $\text{MF}_{\text{bi}}^{\text{R}}(W)_f$ is defined analogously, and we have the following result.

Theorem 2.18. $(\text{MF}_{\text{bi}}(W)_f, \otimes, \alpha, I, \lambda, \rho)$ and $(\text{MF}_{\text{bi}}^{\text{R}}(W)_f, \otimes, \alpha, (I, U_I), \lambda, \rho)$ are monoidal categories.

The proof is given in appendix A.1 where we also discuss the inverses to λ_D and ρ_D as needed. Our proof requires the restriction to $\text{MF}_{\text{bi}}(W)_f$ and $\text{MF}_{\text{bi}}^{\text{R}}(W)_f$, but we suspect that this is not necessary.

3. Applications

In this section we examine matrix bi-factorisations of $W = x^d$ in more detail and compare them to the description of defects in the associated $\mathcal{N} = 2$ minimal conformal field theories. We do so by computing fusing matrices and find that both descriptions produce the same results. This is expected from the Landau-Ginzburg/CFT correspondence and adds strength to conjecture 1.1. As is familiar from other situations, the explicit matrix factorisation computations are more straightforward than the conformal field theory analysis.

3.1. Conformal field theory calculations

In this subsection we use the coset description of $\mathcal{N} = 2$ minimal models to compute a number of spectra of defect fields and fusing matrices in the defect category.

3.1.1. $\mathcal{N} = 2$ minimal models

Let \mathfrak{sVir}_d be the $\mathcal{N} = 2$ super-Virasoro vertex operator algebra with central charge $c_d = 3 - \frac{6}{d}$ for some integer $d \geq 3$, see e. g. [KS, EG, Ad, HM]. Its bosonic part $(\mathfrak{sVir}_d)_{\text{bos}}$ can be obtained via the coset construction, the relevant coset being $(\widehat{\mathfrak{su}}(2)_{d-2} \oplus \widehat{\mathfrak{u}}(1)_4) / \widehat{\mathfrak{u}}(1)_{2d}$. The vertex operator algebra $(\mathfrak{sVir}_d)_{\text{bos}}$ contains in particular the bosonic stress tensor T with zero mode L_0 , and a $\widehat{\mathfrak{u}}(1)$ -current J with zero mode J_0 .

Denote the category of representations of $(\mathfrak{sVir}_d)_{\text{bos}}$ by $\mathcal{C}_d^{\mathcal{N}=2}$; it is a semi-simple braided monoidal category. The simple objects in $\mathcal{C}_d^{\mathcal{N}=2}$, i. e. the irreducible representations of $(\mathfrak{sVir}_d)_{\text{bos}}$, are labelled by triples

$$[l, m, s] \quad \text{where} \quad l \in \{0, \dots, d-2\}, \quad m \in \mathbb{Z}_{2d}, \quad s \in \mathbb{Z}_4 \quad (3.1)$$

subject to the selection rule and field identification

$$l + m + s \in 2\mathbb{Z} \quad \text{and} \quad [l, m, s] = [d-2-l, d+m, s+2]. \quad (3.2)$$

Let $R_{[l,m,s]}$ be the corresponding representation. The decomposition of the fusion tensor product of two representations $R_{[l,m,s]}$ and $R_{[l',m',s']}$ is described by the fusion ring with

$$[l, m, s] \star [l', m', s'] = \sum_{u=|l-l'|}^{\min(l+l', 2d-4-l-l')}^{(+2)} [u, m+m', s+s'], \quad (3.3)$$

where the superscript $(+2)$ means that the sum is carried out in steps of two. The chiral and anti-chiral primaries are those states in the NS sector (i. e. in representations $R_{[l,m,s]}$ with s even) for which $q = \pm 2h$, with q the J_0 -eigenvalue

of the state and h its L_0 -eigenvalue. Chiral or anti-chiral primaries occur only in the representations $R_{[l,\pm l,0]}$ for $l \in \{0, \dots, d-2\}$, and there they are the highest weight states. The relevant eigenvalues are

$$q_{[l,l,0]} = 2h_{[l,l,0]} = -q_{[l,-l,0]} = 2h_{[l,-l,0]} = \frac{l}{d}. \quad (3.4)$$

We will be interested in the A-type modular invariant theory. The bosonic part of its space of bulk fields reads

$$\mathcal{H}_{\text{bos}} = \bigoplus_{[l,m,s]} R_{[l,m,s]} \otimes \bar{R}_{[l,m,-s]}, \quad (3.5)$$

where the sum runs over the set of labels for irreducible representations. The representations containing chiral primaries are thus precisely

$$\mathcal{H}_{\text{bos,chiral}} = \bigoplus_{l=0}^{d-2} R_{[l,l,0]} \otimes \bar{R}_{[l,l,0]} \subset \mathcal{H}_{\text{bos}}. \quad (3.6)$$

3.1.2. Topological defects in the A-type $\mathcal{N} = 2$ minimal models

Let us denote by $\mathcal{D}_d^{\mathcal{N}=2}$ the monoidal category describing the topological defects in the bosonic subsector of the A-type $\mathcal{N} = 2$ minimal model at level $d-2$ which preserve the entire bosonic chiral algebra $(\mathfrak{sVir}_d)_{\text{bos}} \otimes (\overline{\mathfrak{sVir}_d})_{\text{bos}}$. Recall from the introduction that

- the objects of $\mathcal{D}_d^{\mathcal{N}=2}$ label different defect conditions,
- the tensor product describes the fusion of defect lines,
- the morphism spaces $\text{Hom}_{\mathcal{D}_d^{\mathcal{N}=2}}(X_1 \otimes \dots \otimes X_m, Y_1 \otimes \dots \otimes Y_n)$ describe defect-junction fields of left/right conformal weight $(0,0)$ that sit on a junction point with m incoming defect lines labelled X_1, \dots, X_m and n outgoing defect lines labelled Y_1, \dots, Y_n .

The monoidal category $\mathcal{D}_d^{\mathcal{N}=2}$ can be computed using the methods of [FRS1, Fr1, Fr3]. This is sketched in appendix A.2, here we just summarise the results, namely one finds that

$$\mathcal{D}_d^{\mathcal{N}=2} \cong \mathcal{C}_d^{\mathcal{N}=2} \quad \text{as monoidal categories.} \quad (3.7)$$

Consequently, we can use the same set of labels $[l, m, s]$ for simple objects in $\mathcal{D}_d^{\mathcal{N}=2}$ as we did for $\mathcal{C}_d^{\mathcal{N}=2}$. To reduce confusion, let us denote the simple objects of $\mathcal{D}_d^{\mathcal{N}=2}$ by $D_{[l,m,s]}$ rather than by $R_{[l,m,s]}$. Because of the equivalence (3.7), the tensor product of two such irreducible objects decomposes as

$$D_{[l,m,s]} \otimes D_{[l',m',s']} \cong \bigoplus_{[u,n,t]} D_{[u,n,t]}, \quad (3.8)$$

where the sum runs over all labels $[u, n, t]$ that occur in the fusion product $[l, m, s] \star [l', m', s']$ as given in (3.3). Note that $\mathbf{1} = D_{[0,0,0]}$ is the unit object in $\mathcal{D}_d^{\mathcal{N}=2}$; it corresponds to the identity defect. The fusion ring of elementary defects in $\mathcal{N} = 2$ minimal models has also been computed in [BR].

We can now ask about the space of 2-fold junction fields, i.e. defect fields on a given defect, or defect changing fields between two distinct defects. For comparison with the matrix factorisation computation in the next subsection, it will be useful to know the chiral primaries in the spaces of bosonic defect changing fields from the defect labelled $[0, 2n, 0]$ to the defect labelled $[u, u + 2n, 0]$. These are given in [BG, ERR] (as boundary spectra of permutation branes), and we also compute them in the appendix, with the result that the space of bosonic defect changing fields $[0, 2n, 0] \rightarrow [u, u + 2n, 0]$ is

$$\mathcal{H}_{\text{bos}}^{[0,2n,0] \rightarrow [u,u+2n,0]} = \bigoplus_{[l,m,s],[l',m',s']} R_{[l,m,s]} \otimes \bar{R}_{[l',m',s']}, \quad (3.9)$$

where the sum runs over those pairs $[l, m, s], [l', m', s']$ which have the property that $[u, u, 0]$ appears in the fusion product $[l, m, s] \star [l', -m', s']$. In particular, the result is independent of n . It is easy to see that the subspace of representations containing chiral primaries is

$$\mathcal{H}_{\text{bos, chiral}}^{[0,2n,0] \rightarrow [u,u+2n,0]} = \bigoplus_{l=0}^{d-u-2} R_{[u+l,u+l,0]} \otimes \bar{R}_{[l,l,0]}. \quad (3.10)$$

In the special case $n = u = 0$, the expression (3.9) describes the defect fields on the identity defect, which is nothing but the space of bulk fields (3.5).

3.1.3. Some examples of fusing matrices in the defect category

As described in the introduction, we can use two-to-one defect junctions to build many-to-one defect junctions. In the semi-simple case, we can use the two-to-one junctions to build two different bases of the space of three-to-one junctions. These bases are related by the fusing matrices, or 6j-symbols.

For the comparison to matrix factorisations we will only use the subset of objects labelled by $[i, j, 0]$. Let $D_{[m,n,0]}$ be a simple object that appears in the decomposition of $D_{[i,j,0]} \otimes D_{[k,l,0]}$ as given in (3.8). It occurs with multiplicity one, and a basis of the corresponding morphism space is provided by choosing a non-zero element

$$\lambda_{(i,j)(k,l)}^{(m,n)} \in \text{Hom}_{\mathcal{D}_d^{\mathcal{N}=2}}(D_{[i,j,0]} \otimes D_{[k,l,0]}, D_{[m,n,0]}). \quad (3.11)$$

Fix a non-zero $\lambda_{(i,j)(k,l)}^{(m,n)}$ in each such non-zero Hom-space. We now list the fusing matrices in the category $\mathcal{D}_d^{\mathcal{N}=2}$ in three examples. The corresponding calculations are sketched in appendix A.2.

■ $(D_{[0,2i,0]} \otimes D_{[0,2j,0]}) \otimes D_{[0,2k,0]} \longrightarrow D_{[0,2i+2j+2k,0]}$ with $i, j, k \in \mathbb{Z}_d$:

From the fusion rules we see that this morphism space is one-dimensional. The two basis vectors obtained from combining the elements (3.11) are necessarily related by a non-zero constant,

$$\lambda_{(0,2i)(0,2j+2k)}^{(0,2i+2j+2k)} \circ (\text{id} \otimes \lambda_{(0,2j)(0,2k)}^{(0,2j+2k)}) \circ \alpha = \psi_{i,j,k} \cdot \lambda_{(0,2i+2j)(0,2k)}^{(0,2i+2j+2k)} \circ (\lambda_{(0,2i)(0,2j)}^{(0,2i+2j)} \otimes \text{id}), \quad (3.12)$$

where $\psi_{i,j,k} \in \mathbb{C}^\times$ and α denotes the associator in $\mathcal{D}_d^{\mathcal{N}=2}$. The value of the constants $\psi_{i,j,k}$ depends on the choice of basis vectors $\lambda_{(0,2i)(0,2j)}^{(0,2i+2j)}$. The pentagon identity implies that ψ is a 3-cocycle on \mathbb{Z}_d , and different basis choices change it by a coboundary, see e.g. [MS, FRS3]. The basis independent information contained in ψ is thus a class in $H^3(\mathbb{Z}_d, \mathbb{C}^\times) \cong \mathbb{Z}_d$. For the specific choices made in the calculation in appendix A.2 one simply finds

$$\psi_{i,j,k} = 1 \quad \text{for all } i, j, k \in \mathbb{Z}_d, \quad (3.13)$$

so that ψ represents the trivial class in $H^3(\mathbb{Z}_d, \mathbb{C}^\times)$.

■ $(D_{[d-2,0,0]} \otimes D_{[f,f,0]}) \otimes D_{[d-2,0,0]} \longrightarrow D_{[f,f,0]}$, $d \in 2\mathbb{Z}$, $f = \frac{d}{2} - 1$:

The fusion rules are $D_{[f,f,0]} \otimes D_{[d-2,0,0]} \cong D_{[f,f,0]}$, so that this morphism space is again one-dimensional, and we find

$$\lambda_{(d-2,0)(f,f)}^{(f,f)} \circ (\text{id} \otimes \lambda_{(f,f)(d-2,0)}^{(f,f)}) \circ \alpha = (-1)^{(d-2)/2} \cdot \lambda_{(f,f)(d-2,0)}^{(f,f)} \circ (\lambda_{(d-2,0)(f,f)}^{(f,f)} \otimes \text{id}). \quad (3.14)$$

Because the same basis morphisms λ appear on either side of this equation, the factor $(-1)^{(d-2)/2}$ is independent of the choice of basis in (3.11).

■ $(D_{[1,1,0]} \otimes D_{[1,1,0]}) \otimes D_{[1,1,0]} \longrightarrow D_{[1,3,0]}$, $d \geq 4$:

This morphism space is two-dimensional, and accordingly the two bases are related by a (2×2) -matrix F , namely, for $p \in \{0, 2\}$,

$$\lambda_{(1,1)(p,2)}^{(1,3)} \circ (\text{id} \otimes \lambda_{(1,1)(1,1)}^{(p,2)}) \circ \alpha = \sum_{q \in \{0,2\}} F_{pq} \cdot \lambda_{(q,2)(1,1)}^{(1,3)} \circ (\lambda_{(1,1)(1,1)}^{(q,2)} \otimes \text{id}). \quad (3.15)$$

Under a rescaling $\lambda_{(i,j)(k,l)}^{(m,n)} \mapsto \eta_{(i,j)(k,l)}^{(m,n)} \cdot \lambda_{(i,j)(k,l)}^{(m,n)}$ for some $\eta_{(i,j)(k,l)}^{(m,n)} \in \mathbb{C}^\times$, the entries of F change as

$$\begin{aligned} F_{00} &\longmapsto \frac{\eta_{(0,2)(1,1)}^{(1,3)}}{\eta_{(1,1)(0,2)}^{(1,3)}} F_{00}, & F_{02} &\longmapsto \frac{\eta_{(2,2)(1,1)}^{(1,3)} \eta_{(1,1)(1,1)}^{(2,2)}}{\eta_{(1,1)(0,2)}^{(1,3)} \eta_{(1,1)(1,1)}^{(0,2)}} F_{02}, \\ F_{20} &\longmapsto \frac{\eta_{(0,2)(1,1)}^{(1,3)} \eta_{(1,1)(1,1)}^{(0,2)}}{\eta_{(1,1)(2,2)}^{(1,3)} \eta_{(1,1)(1,1)}^{(2,2)}} F_{20}, & F_{22} &\longmapsto \frac{\eta_{(2,2)(1,1)}^{(1,3)}}{\eta_{(1,1)(2,2)}^{(1,3)}} F_{22}. \end{aligned} \quad (3.16)$$

The six independent constants can be used to adjust three of the four entries of F at will. The fourth one is then fixed, because the constants $\eta_{(i,j)(k,l)}^{(m,n)}$ cancel from the ratio

$$\frac{F_{00}F_{22}}{F_{02}F_{20}}. \quad (3.17)$$

Provided that all entries of F are non-zero, this ratio is thus the only basis-independent information contained in F . For the choice for (3.11) implicit in appendix A.2 one has

$$F_{00} = -F_{22} = \frac{-1}{2 \cos \frac{\pi}{d}}, \quad F_{02} = F_{20} = \frac{\sqrt{\sin \frac{\pi}{d} \sin \frac{3\pi}{d}}}{\sin \frac{2\pi}{d}} \quad (3.18)$$

so that

$$\frac{F_{00}F_{22}}{F_{02}F_{20}} = \frac{-1}{1 + 2 \cos \frac{2\pi}{d}}. \quad (3.19)$$

3.2. Matrix bi-factorisation calculations

We will now reproduce the conformal field theory results of the previous subsection in the technically less pronounced matrix bi-factorisation formalism. To do so, we make use of the fact [BR] that the defect

$$D_{[l,l+2m,0]} \text{ corresponds to } P_{\{m,\dots,m+l\}} \quad (3.20)$$

in $\text{MF}_{\text{bi}}(x^d)$. However, in order to compare with conformal field theory the object $P_{\{m,\dots,m+l\}}$ needs to be lifted to the graded category $\text{MF}_{\text{bi}}^{\text{R}}(x^d)$. This is easily done by constructing R-charge matrices $U_{m,\dots,m+l} \equiv U_{P_{\{m,\dots,m+l\}}}$ with respect to which the morphisms in $\text{MF}_{\text{bi}}^{\text{R}}(x^d)$ have the same charges as the defect fields in the CFT description.

More precisely, we can check that

$$\text{Hom}_{\text{MF}_{\text{bi}}(x^d)}(P_{\{0\}}, P_{\{0,\dots,j\}}) \cong \text{Hom}_{\text{MF}_{\text{bi}}(x^d)}(P_{\{i\}}, P_{\{i,\dots,i+j\}}) \quad (3.21)$$

for all $i, j \in \mathbb{N}$, and an explicit basis for $\text{Hom}_{\text{MF}_{\text{bi}}(x^d)}(P_{\{0\}}, P_{\{0,\dots,j\}})$ is given by

$$\phi_k = \begin{pmatrix} [a^k p_{\{1,\dots,j\}}(a, b)]^\wedge & 0 \\ 0 & \hat{a}^k \end{pmatrix}, \quad k \in \{0, \dots, d-j-2\}. \quad (3.22)$$

But from (3.4) and (3.10) we know that also the charges of the spectra $P_{\{0\}} \rightarrow P_{\{0,\dots,j\}}$ and $P_{\{i\}} \rightarrow P_{\{i,\dots,i+j\}}$ coincide, which means that we must have $U_{0,\dots,j}(\alpha) = U_{i,\dots,i+j}(\alpha)$. Then with the general ansatz $U_{0,\dots,j}(\alpha) = \begin{pmatrix} e^{iqx\alpha c_1(j)} & 0 \\ 0 & e^{iqx\alpha c_2(j)} \end{pmatrix}$ we find that the condition of $P_{\{0,\dots,j\}}$ having R-charge 1 translates into $c_2(j) = c_1(j) + j + \frac{2-d}{2}$. The constants $c_1(j)$ can be fixed by computing the R-charges of ϕ_k to be

$q_k = 2(c_1(j) + j + k)/d$. Comparing this with (3.4) and (3.10) we get $c_1(j) = -\frac{j}{2}$ and thus

$$U_{0,\dots,j}(\alpha) = U_{i,\dots,i+j}(\alpha) = \begin{pmatrix} e^{-iq_x\alpha j/2} & 0 \\ 0 & e^{iq_x\alpha(j-d+2)/2} \end{pmatrix}. \quad (3.23)$$

Now we can treat the examples of the previous subsection in the Landau-Ginzburg description.

■ $(P_{\{i\}} \otimes P_{\{j\}}) \otimes P_{\{k\}} \longrightarrow P_{\{i+j+k\}}$:

A possible choice of non-zero morphisms in $\text{MF}_{\text{bi}}^{\text{R}}(x^d)$ corresponding to the defect fields $\lambda_{(0,2i)(0,2j)}^{(0,2i+2j)}$ of charge 0 between group-like defects is

$$\Lambda_{i,j} = \begin{pmatrix} M_j & 0 & 0 & 0 \\ 0 & 0 & M_j & 0 \end{pmatrix} : P_{\{i\}} \otimes P_{\{j\}} \longrightarrow P_{\{i+j\}} \quad (3.24)$$

where M_j is the bimodule map

$$R \otimes_{\mathbb{C}} R \otimes_{\mathbb{C}} R \longrightarrow R \otimes_{\mathbb{C}} R, \quad 1 \otimes_{\mathbb{C}} x^m \otimes_{\mathbb{C}} 1 \longmapsto \eta^{jm} \cdot 1 \otimes_{\mathbb{C}} x^m. \quad (3.25)$$

It is straightforward to check that

$$\Lambda_{i,j+k} \circ (\text{id} \otimes \Lambda_{j,k}) \circ \alpha = 1 \cdot \Lambda_{i+j,k} \circ (\Lambda_{i,j} \otimes \text{id}) \quad (3.26)$$

already in $\text{DG}_{\text{bi}}^{\text{R}}(x^d)$. Hence we recover the same fusing ‘‘matrix’’ 1 as in (3.13).

■ $(P_{\{\frac{d}{2}+1,\dots,\frac{3d}{2}-1\}} \otimes P_{\{0,\dots,\frac{d}{2}-1\}}) \otimes P_{\{\frac{d}{2}+1,\dots,\frac{3d}{2}-1\}} \longrightarrow P_{\{0,\dots,\frac{d}{2}-1\}}$, $d \in 2\mathbb{Z}$:

Let us next consider the twice as interesting example where the fusing ‘‘matrix’’ can be +1 or -1. According to the dictionary (3.20) there should be morphisms

$$\begin{aligned} F_l &: P_{\{\frac{d}{2}+1,\dots,\frac{3d}{2}-1\}} \otimes P_{\{0,\dots,\frac{d}{2}-1\}} \longrightarrow P_{\{0,\dots,\frac{d}{2}-1\}}, \\ F_r &: P_{\{0,\dots,\frac{d}{2}-1\}} \otimes P_{\{\frac{d}{2}+1,\dots,\frac{3d}{2}-1\}} \longrightarrow P_{\{0,\dots,\frac{d}{2}-1\}} \end{aligned}$$

in $\text{MF}_{\text{bi}}^{\text{R}}(x^d)$ that correspond to the defect fields $\lambda_{(f,f)(d-2,0)}^{(f,f)}$, $\lambda_{(d-2,0)(f,f)}^{(f,f)}$ in (3.14). A possible choice is

$$F_l = \begin{pmatrix} 0 & M'_{\frac{d}{2}} & 0 & 0 \\ 0 & 0 & (-1)^{\frac{d}{2}-1} M'_{\frac{d}{2}} & 0 \end{pmatrix}, \quad F_r = \begin{pmatrix} 0 & M_{\frac{d}{2}} & 0 & 0 \\ 0 & 0 & 0 & M_{\frac{d}{2}} \end{pmatrix} \quad (3.27)$$

of R-charge 0, where M'_j is the bimodule map

$$R \otimes_{\mathbb{C}} R \otimes_{\mathbb{C}} R \longrightarrow R \otimes_{\mathbb{C}} R, \quad 1 \otimes_{\mathbb{C}} x^m \otimes_{\mathbb{C}} 1 \longmapsto \eta^{-jm} \cdot x^m \otimes_{\mathbb{C}} 1. \quad (3.28)$$

Again, one easily computes that

$$F_l \circ (\text{id} \otimes F_r) \circ \alpha = (-1)^{(d-2)/2} \cdot F_r \circ (F_l \otimes \text{id}) \quad (3.29)$$

holds already on the DG level, and we find the same fusing matrix $(-1)^{(d-2)/2}$ as in (3.14).

■ $P_{\{1,2\}} \longrightarrow (P_{\{0,1\}} \otimes P_{\{0,1\}}) \otimes P_{\{0,1\}}$:

Finally, we now treat the last example of subsection 3.1 involving a (2×2) -fusing matrix. In contrast to the previous simple examples it will not be sufficient to work in the DG category $\text{DG}_{\text{bi}}^{\text{R}}(x^d)$, but we will find that the analogue of equations (3.15) only holds in $\text{MF}_{\text{bi}}^{\text{R}}(x^d)$. Also, it turns out that the calculations here are easier in the dual, i.e. one-to-many, direction. As we recalled in the introduction, this will produce the same fusing matrix entries.

A possible choice of zero-charge morphisms corresponding to those which are dual to the six defect fields in (3.15) is

$$A_1 = A_{\{1\}}^{\{0,1\}\{0,1\}} = \begin{pmatrix} [a - (\eta + 1)x + \eta b]^{\wedge} \circ J & 0 \\ (A_1)_{21} & 0 \\ 0 & J \\ 0 & -\eta J \end{pmatrix}, \quad (3.30a)$$

$$A_2 = A_{\{0,1,2\}}^{\{0,1\}\{0,1\}} = \begin{pmatrix} J & 0 \\ (A_2)_{21} & 0 \\ 0 & [a + (\eta + 1)x - (\eta^2 + \eta + 1)b]^{\wedge} \circ J \\ 0 & [(\eta^2 + \eta + 1)a - (\eta^2 + \eta)x - \eta^2 b]^{\wedge} \circ J \end{pmatrix}, \quad (3.30b)$$

$$A_3 = A_{\{1,2\}}^{\{0,1\}\{1\}} = \begin{pmatrix} J & 0 \\ (A_3)_{21} & 0 \\ 0 & J \\ 0 & [(\eta + 1)a - \eta x - \eta^2 b]^{\wedge} \circ J \end{pmatrix}, \quad (3.30c)$$

$$A_4 = A_{\{1,2\}}^{\{0,1\}\{0,1\}} = \begin{pmatrix} (A_4)_{11} & 0 \\ (A_4)_{21} & 0 \\ 0 & (A_4)_{32} \\ 0 & -J \end{pmatrix}, \quad (3.30d)$$

$$A_5 = A_{\{1,2\}}^{\{1\}\{0,1\}} = \begin{pmatrix} J & 0 \\ (A_5)_{21} & 0 \\ 0 & [a + \eta x - (\eta^2 + \eta)b]^{\wedge} \circ J \\ 0 & \eta^2 J \end{pmatrix}, \quad (3.30e)$$

$$A_6 = A_{\{1,2\}}^{\{0,1,2\}\{0,1\}} = \begin{pmatrix} [a - (\eta^2 + \eta + 1)x + (\eta^2 + \eta)b]^{\wedge} \circ J & 0 \\ (A_6)_{21} & 0 \\ 0 & J \\ 0 & (A_6)_{42} \end{pmatrix} \quad (3.30f)$$

where the map $J : R \otimes_{\mathbb{C}} R \rightarrow R \otimes_{\mathbb{C}} R \otimes_{\mathbb{C}} R$ is given by $r \otimes_{\mathbb{C}} s \mapsto r \otimes_{\mathbb{C}} 1 \otimes_{\mathbb{C}} s$ and we have

$$(A_4)_{11} = [(\eta^{-1} + \eta^{-2})a - (1 + \eta^{-1} + \eta^{-2})x - b]^{\wedge} \circ J$$

$$(A_4)_{32} = [(\eta^{-1} + \eta^{-2})a + \frac{1 + \eta^{-1}}{\eta + 1}x - (\eta + 1 + \eta^{-1})b]^\wedge \circ J ,$$

$$(A_6)_{42} = [-(\eta^3 + \eta^2 + \eta)a + \eta^3x + (\eta^4 + \eta^3)b]^\wedge \circ J ,$$

$$(A_K^{IJ})_{21} = ((A_K^{IJ})_{42}\hat{p}_{\{0,\dots,d-1\}\setminus I}(a, x) - (A_K^{IJ})_{32}\hat{p}_{\{0,\dots,d-1\}\setminus J}(x, b)) / \hat{p}_K(a, b) .$$

One may check that the entries $(A_K^{IJ})_{21}$ are indeed polynomials in a, x, b . The fusing matrix $\tilde{F} = \begin{pmatrix} \tilde{F}_{00} & \tilde{F}_{02} \\ \tilde{F}_{20} & \tilde{F}_{22} \end{pmatrix}$ is to be determined from the equations

$$\alpha \circ (A_1 \otimes \text{id}) \circ A_5 = \tilde{F}_{00} \cdot (\text{id} \otimes A_1) \circ A_3 + \tilde{F}_{02} \cdot (\text{id} \otimes A_2) \circ A_4 , \quad (3.31a)$$

$$\alpha \circ (A_2 \otimes \text{id}) \circ A_6 = \tilde{F}_{20} \cdot (\text{id} \otimes A_1) \circ A_3 + \tilde{F}_{22} \cdot (\text{id} \otimes A_2) \circ A_4 \quad (3.31b)$$

of (8×2) -matrices. It turns out that only the constraints on \tilde{F} coming from the entries $(1, 1)$, $(5, 2)$, $(6, 2)$ and $(7, 2)$ can be solved without descending to cohomology, the unique solution being

$$\tilde{F} = \frac{1}{\eta + 1} \begin{pmatrix} -\eta & \eta^2 \\ \eta^2 + \eta + 1 & \eta^2 \end{pmatrix} , \quad (3.32)$$

while the whole of equations (3.31) can only hold in $\text{MF}_{\text{bi}}^{\text{R}}(x^d)$. Instead of constructing a general homotopy up to which (3.31) is satisfied, we have contended ourselves with checking that (3.32) solves (3.31) in $\text{MF}_{\text{bi}}^{\text{R}}(x^d)$ for many values of d with the help of the computer algebra system Singular [GPS].

We note that (3.32) is not the same fusing matrix as the result (3.18) but that the normalisation-invariant quantity (3.17) indeed produces the same result as (3.19):

$$\frac{\tilde{F}_{00}\tilde{F}_{22}}{\tilde{F}_{02}\tilde{F}_{20}} = \frac{-\eta}{\eta^2 + \eta + 1} = \frac{-1}{1 + 2 \cos \frac{2\pi}{d}} . \quad (3.33)$$

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A. Appendix

A.1. Inverses for unit isomorphisms

One-variable case

To give the inverse morphisms and the homotopies we need some more notation. Let U, U', V, V' be vector spaces. If $\phi(a, b) \in \text{Hom}(U, V[a, b])$ and $\phi'(a, b) \in \text{Hom}(U', V'[a, b])$, we define the map

$$\phi(a, x) \otimes_{\mathbb{C}} \mathbb{1}_R \otimes_{\mathbb{C}} \phi'(x, b) : U \otimes_{\mathbb{C}} R \otimes_{\mathbb{C}} U' \longrightarrow (V \otimes_{\mathbb{C}} R \otimes_{\mathbb{C}} V')[a, b] \quad (\text{A.1})$$

to be

$$u \otimes_{\mathbb{C}} r \otimes_{\mathbb{C}} u' \longmapsto \sum_{k,l,m,n} \phi_{kl}(u) \otimes_{\mathbb{C}} x^{l+m} r \otimes_{\mathbb{C}} \phi'_{mn}(u') a^k b^n. \quad (\text{A.2})$$

That is, the formal variable x acts by multiplication on the middle factor R . With this definition we have

$$\hat{\phi} \otimes_R \hat{\phi}' = [\phi(a, x) \otimes_{\mathbb{C}} \mathbb{1}_R \otimes_{\mathbb{C}} \phi'(x, b)]^\wedge. \quad (\text{A.3})$$

More generally we set, for $\alpha, \beta, \gamma, \delta$ each standing for one of the formal variables a, x, b ,

$$\phi(\alpha, \beta) \otimes_{\mathbb{C}} \mathbb{1}_R \otimes_{\mathbb{C}} \phi'(\gamma, \delta) = \sum_{k,l,m} \xi_{k,l,m} \circ (\mathbb{1}_U \otimes_{\mathbb{C}} (x^l \cdot) \otimes_{\mathbb{C}} \mathbb{1}_{U'}) a^k b^m, \quad (\text{A.4})$$

where $\xi_{k,l,m}$ are defined by $\phi(\alpha, \beta) \otimes_{\mathbb{C}} \mathbb{1}_R \otimes_{\mathbb{C}} \phi'(\gamma, \delta) = \sum_{k,l,m} \xi_{k,l,m} a^k x^l b^m$. Finally, we define the maps

$$\begin{aligned} \mathbb{1} \otimes_{\mathbb{C}} \mathbb{1}_U : U &\longrightarrow R \otimes_{\mathbb{C}} U, & \mathbb{1}_U \otimes_{\mathbb{C}} \mathbb{1} : U &\longrightarrow U \otimes_{\mathbb{C}} R, \\ u &\longmapsto \mathbb{1} \otimes_{\mathbb{C}} u, & u &\longmapsto u \otimes_{\mathbb{C}} \mathbb{1}. \end{aligned} \quad (\text{A.5})$$

By slight abuse of notation we will abbreviate, for $\alpha, \beta \in \{a, x, b\}$,

$$(1 \otimes_{\mathbb{C}} \phi(\alpha, \beta)) \equiv (\mathbb{1}_R \otimes_{\mathbb{C}} \phi(\alpha, \beta)) \circ (1 \otimes_{\mathbb{C}} \mathbb{1}_U), \quad (\text{A.6a})$$

$$(\phi(\alpha, \beta) \otimes_{\mathbb{C}} 1) \equiv (\phi(\alpha, \beta) \otimes_{\mathbb{C}} \mathbb{1}_R) \circ (\mathbb{1}_U \otimes_{\mathbb{C}} 1). \quad (\text{A.6b})$$

Note that then also

$$1 \otimes_{\mathbb{C}} \phi(a, b) = (1 \otimes_{\mathbb{C}} \mathbb{1}_V) \circ \phi(a, b), \quad \phi(a, b) \otimes_{\mathbb{C}} 1 = (\mathbb{1}_V \otimes_{\mathbb{C}} 1) \circ \phi(a, b). \quad (\text{A.7})$$

However, the corresponding equalities for $1 \otimes_{\mathbb{C}} \phi(x, b)$ or $\phi(a, x) \otimes_{\mathbb{C}} 1$ do not hold (or even make sense).

It is enough to show that λ_D and ρ_D are invertible for $D = (D_0, D_1, d_0, d_1)$ where $D_0 = R \otimes_{\mathbb{C}} \check{D}_0 \otimes_{\mathbb{C}} R$ and $D_1 = R \otimes_{\mathbb{C}} \check{D}_1 \otimes_{\mathbb{C}} R$ for some vector spaces \check{D}_0 ,

\check{D}_1 . The general statement follows since every free R -bimodule is isomorphic to one of the form $R \otimes_{\mathbb{C}} U \otimes_{\mathbb{C}} R$ and since λ_D and ρ_D are natural in D . We will show that the inverse of λ_D is given by

$$\lambda_D^{-1} = \begin{pmatrix} [1 \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_0}]^{\wedge} & 0 \\ [\frac{1 \otimes_{\mathbb{C}} \check{d}_0(a,b) - 1 \otimes_{\mathbb{C}} \check{d}_0(x,b)}{a-x}]^{\wedge} & 0 \\ 0 & [\frac{1 \otimes_{\mathbb{C}} \check{d}_1(a,b) - 1 \otimes_{\mathbb{C}} \check{d}_1(x,b)}{a-x}]^{\wedge} \\ 0 & [1 \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_1}]^{\wedge} \end{pmatrix} : D \longrightarrow I \otimes D. \quad (\text{A.8})$$

First we need to check that λ_D^{-1} is a morphism in $\text{MF}_{\text{bi}}(W)$, i. e.

$$(I \otimes D) \circ \lambda_D^{-1} = \lambda_D^{-1} \circ D. \quad (\text{A.9})$$

Let us denote the left-hand side of this equation by $\mathcal{L} = \begin{pmatrix} 0 & \mathcal{L}_1 \\ \mathcal{L}_0 & 0 \end{pmatrix}$ and the right-hand side by $\mathcal{R} = \begin{pmatrix} 0 & \mathcal{R}_1 \\ \mathcal{R}_0 & 0 \end{pmatrix}$. Using the composition rule below (2.14), as well as (A.3), (A.6) and (A.7) we find

$$\begin{aligned} (\mathcal{L}_0)_2 &= (\mathbb{1}_{I_0} \otimes_R d_0) \circ [1 \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_0}]^{\wedge} + [1 \otimes_{\mathbb{C}} \check{d}_0(a,b) - 1 \otimes_{\mathbb{C}} \check{d}_0(x,b)]^{\wedge} \\ &= [1 \otimes_{\mathbb{C}} \check{d}_0(a,b)]^{\wedge} = [(1 \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_1}) \circ \check{d}_0(a,b)]^{\wedge} = (\mathcal{R}_0)_2. \end{aligned} \quad (\text{A.10})$$

Similarly one gets $(\mathcal{L}_1)_1 = (\mathcal{R}_1)_1$. Next we compute

$$\begin{aligned} (\mathcal{L}_0)_1 - (\mathcal{R}_0)_1 &= (\iota_0 \otimes_R \mathbb{1}_{D_0}) \circ [1 \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_0}]^{\wedge} - [\frac{1 \otimes_{\mathbb{C}} \check{d}_1(a,b) - \check{d}_1(x,b)}{a-x} \circ \check{d}_0(a,b)]^{\wedge} \\ &\quad - [(1 \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_1}) \circ \frac{1 \otimes_{\mathbb{C}} \check{d}_0(a,b) - \check{d}_0(x,b)}{a-x}]^{\wedge} \\ &= (\iota_0 \otimes_R \mathbb{1}_{D_0}) \circ [1 \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_0}]^{\wedge} - [\frac{1 \otimes_{\mathbb{C}} W(a) - 1 \otimes_{\mathbb{C}} W(x)}{a-x}]^{\wedge} \end{aligned}$$

where we used (A.3) in the first step. If we now expand $\frac{W(a)-W(b)}{a-b} = \sum_{m,n} V_{mn} a^m b^n$ we find that the last two terms in the above computation cancel each other:

$$\begin{aligned} ((\iota_0 \otimes_R \mathbb{1}_{D_0}) \circ [1 \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_0}]^{\wedge})(r \otimes_{\mathbb{C}} v \otimes_{\mathbb{C}} s) &= \sum_{m,n} V_{mn} r x^m \otimes_{\mathbb{C}} x^n \otimes_{\mathbb{C}} v \otimes_{\mathbb{C}} s \\ &= [\frac{1 \otimes_{\mathbb{C}} W(a) - 1 \otimes_{\mathbb{C}} W(x)}{a-x}]^{\wedge}(r \otimes_{\mathbb{C}} v \otimes_{\mathbb{C}} s). \end{aligned}$$

Similarly one shows that $(\mathcal{L}_1)_2 = (\mathcal{R}_1)_2$. This means that the condition (A.9) is satisfied and λ_D^{-1} is indeed a morphism.

It remains to be seen that λ_D^{-1} is inverse to λ_D . Because of $(\mu \otimes_R \mathbb{1}_{D_i}) \circ [1 \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_i}]^{\wedge} = \mathbb{1}_{D_i}$ it immediately follows that $\lambda_D \circ \lambda_D^{-1} = \mathbb{1}_D$. Conversely, however, $[1 \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_0}]^{\wedge} \circ (\mu \otimes_R \mathbb{1}_{D_0})$ takes $r \otimes_{\mathbb{C}} s \otimes_{\mathbb{C}} v \otimes_{\mathbb{C}} t \in R \otimes_{\mathbb{C}} R \otimes_{\mathbb{C}} D_0 \otimes_{\mathbb{C}} R$ to $rs \otimes_{\mathbb{C}} 1 \otimes_{\mathbb{C}} v \otimes_{\mathbb{C}} t$, and we are lead to introduce a map $r_{xa} \in \text{Hom}(R, R[a,b])$ given by $r_{xa}(x^n) = a^n$ (independent of b). This can be written as a power series

$$r_{xa} = \sum_{n=0}^{\infty} a^n \delta_{x^n} \quad \text{where} \quad \delta_{x^n} : R \longrightarrow R, \quad \delta_{x^n}(x^m) = \delta_{m,n}. \quad (\text{A.11})$$

The infinite sum does not cause a problem because only finitely many terms are non-zero for a given element of R . We can write

$$[1 \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_0}]^{\wedge} \circ (\mu \otimes_R \mathbb{1}_{D_0}) = [r_{xa} \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_0}]^{\wedge}. \quad (\text{A.12})$$

Next we note that the map $\frac{r_{xa}-1}{a-x}$ from R to $R[a, b]$ is well-defined since $(r_{xa} - 1)(x^n) = a^n - x^n$, and

$$\frac{r_{xa} - 1}{a - x} = \sum_{n=1}^{\infty} \frac{a^n - x^n}{a - x} \delta_{x^n} = \sum_{n=1}^{\infty} \sum_{i=0}^{n-1} a^i x^{n-i-1} \delta_{x^n}. \quad (\text{A.13})$$

Lemma A.1. We have

- (i) $[\frac{r_{xa}-1}{a-x} \otimes_{\mathbb{C}} \mathbb{1}_U]^{\wedge} \circ [(a-x) \otimes_{\mathbb{C}} \mathbb{1}_U]^{\wedge} = -[\mathbb{1}_R \otimes_{\mathbb{C}} \mathbb{1}_U]^{\wedge}$,
- (ii) $[(a-x) \otimes_{\mathbb{C}} \mathbb{1}_U]^{\wedge} \circ [\frac{r_{xa}-1}{a-x} \otimes_{\mathbb{C}} \mathbb{1}_U]^{\wedge} = [1 \otimes_{\mathbb{C}} \mathbb{1}_U]^{\wedge} \circ (\mu \otimes_R \mathbb{1}_{R \otimes_{\mathbb{C}} U \otimes_{\mathbb{C}} R}) - [\mathbb{1}_R \otimes_{\mathbb{C}} \mathbb{1}_U]^{\wedge}$,
- (iii) for $\phi \in \text{Hom}(U, V[a, b])$,
$$\begin{aligned} & [\frac{r_{xa}-1}{a-x} \otimes_{\mathbb{C}} \mathbb{1}_U]^{\wedge} \circ [\mathbb{1}_R \otimes_{\mathbb{C}} \phi(x, b)]^{\wedge} - [\mathbb{1}_R \otimes_{\mathbb{C}} \phi(x, b)]^{\wedge} \circ [\frac{r_{xa}-1}{a-x} \otimes_{\mathbb{C}} \mathbb{1}_U]^{\wedge} \\ &= [\frac{1 \otimes_{\mathbb{C}} \phi(a, b) - 1 \otimes_{\mathbb{C}} \phi(x, b)}{a-x}]^{\wedge} \circ (\mu \otimes_R \mathbb{1}_{R \otimes_{\mathbb{C}} U \otimes_{\mathbb{C}} R}). \end{aligned}$$

Proof. To see (i)–(iii), apply the left hand side to $r \otimes_{\mathbb{C}} x^n \otimes_{\mathbb{C}} u \otimes_{\mathbb{C}} s$. For (i) this gives

$$\begin{aligned} & [\frac{r_{xa}-1}{a-x} \otimes_{\mathbb{C}} \mathbb{1}_U]^{\wedge} ((rx \otimes_{\mathbb{C}} x^n - r \otimes_{\mathbb{C}} x^{n+1}) \otimes_{\mathbb{C}} u \otimes_{\mathbb{C}} s) \\ &= \sum_{i=0}^{n-1} rx^{i+1} \otimes_{\mathbb{C}} x^{n-i-1} \otimes_{\mathbb{C}} u \otimes_{\mathbb{C}} s - \sum_{i=0}^n rx^i \otimes_{\mathbb{C}} x^{n-i} \otimes_{\mathbb{C}} u \otimes_{\mathbb{C}} s \\ &= -r \otimes_{\mathbb{C}} x^n \otimes_{\mathbb{C}} u \otimes_{\mathbb{C}} s, \end{aligned}$$

for (ii) one finds

$$\begin{aligned} & \sum_{i=0}^{n-1} rx^{i+1} \otimes_{\mathbb{C}} x^{n-i-1} \otimes_{\mathbb{C}} u \otimes_{\mathbb{C}} s - \sum_{i=0}^{n-1} rx^i \otimes_{\mathbb{C}} x^{n-i} \otimes_{\mathbb{C}} u \otimes_{\mathbb{C}} s \\ &= ([1 \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_0}]^{\wedge} \circ (\mu \otimes_R \mathbb{1}_{D_0}) - \mathbb{1}_{I_0 \otimes_R D_0})(r \otimes_{\mathbb{C}} x^n \otimes_{\mathbb{C}} u \otimes_{\mathbb{C}} s), \end{aligned}$$

while for (iii) one has

$$\begin{aligned} & \sum_{k,l} \left(\sum_{i=0}^{n+k-1} rx^i \otimes_{\mathbb{C}} x^{n+k-i-1} - \sum_{i=0}^{n-1} rx^i \otimes_{\mathbb{C}} x^{n+k-i-1} \right) \otimes_{\mathbb{C}} \phi_{kl}(u) \otimes_{\mathbb{C}} x^l s \\ &= \sum_{k,l} \sum_{i=0}^{k-1} rx^{n+i} \otimes_{\mathbb{C}} x^{k-i-1} \otimes_{\mathbb{C}} \phi_{kl}(u) \otimes_{\mathbb{C}} x^l s \\ &= [\frac{1 \otimes_{\mathbb{C}} \phi(a, b) - 1 \otimes_{\mathbb{C}} \phi(x, b)}{a-x}]^{\wedge} (rx^n \otimes_{\mathbb{C}} u \otimes_{\mathbb{C}} s). \end{aligned}$$

□

We define

$$\psi_D^l = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & [\frac{r_{xa}-1}{a-x} \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_1}]^\wedge \\ [\frac{r_{xa}-1}{a-x} \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_0}]^\wedge & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \quad (\text{A.14})$$

and claim that this is a homotopy between $\lambda_D^{-1} \circ \lambda_D$ and $\text{id}_{I \otimes D}$, i. e.

$$\lambda_D^{-1} \circ \lambda_D - \text{id}_{I \otimes D} = (I \otimes D) \circ \psi_D^l + \psi_D^l \circ (I \otimes D), \quad (\text{A.15})$$

from which it follows that λ_D really is an isomorphism in $\text{MF}_{\text{bi}}(W)$. To prove the claim, we denote the left-hand side of (A.15) by $\mathcal{L} = \begin{pmatrix} \mathcal{L}_0 & 0 \\ 0 & \mathcal{L}_1 \end{pmatrix}$ and the right-hand side by $\mathcal{R} = \begin{pmatrix} \mathcal{R}_0 & 0 \\ 0 & \mathcal{R}_1 \end{pmatrix}$. The equality $\mathcal{L} = \mathcal{R}$ follows in a straightforward way from lemma A.1. For example

$$\begin{aligned} (\mathcal{R}_0)_{11} &= (\iota_1 \otimes_R \mathbb{1}_{D_0}) \circ [\frac{r_{xa}-1}{a-x} \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_0}]^\wedge = [1 \otimes_{\mathbb{C}} \mathbb{1}_U]^\wedge \circ (\mu \otimes_R \mathbb{1}) - \mathbb{1} = (\mathcal{L}_0)_{11}, \\ (\mathcal{R}_0)_{12} &= 0 = (\mathcal{L}_0)_{12} \\ (\mathcal{R}_0)_{22} &= [\frac{r_{xa}-1}{a-x} \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_1}]^\wedge \circ (\iota_1 \otimes_R \mathbb{1}_{D_1}) = -\mathbb{1} = (\mathcal{L}_0)_{22}, \\ (\mathcal{R}_0)_{21} &= [\frac{r_{xa}-1}{a-x} \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_1}]^\wedge \circ [\mathbb{1}_R \otimes_{\mathbb{C}} \check{d}_0(x, b)]^\wedge - [\mathbb{1}_R \otimes_{\mathbb{C}} \check{d}_0(x, b)]^\wedge \circ [\frac{r_{xa}-1}{a-x} \otimes_{\mathbb{C}} \mathbb{1}_{\check{D}_0}]^\wedge \\ &= [\frac{1 \otimes_{\mathbb{C}} \check{d}_0(a, b) - 1 \otimes_{\mathbb{C}} \check{d}_0(x, b)}{a-x}]^\wedge \circ (\mu \otimes_R \mathbb{1}_{D_0}) = (\mathcal{L}_0)_{21}. \end{aligned}$$

and similarly for the entires of \mathcal{R}_1 and \mathcal{L}_1 .

The inverse of ρ_D is given by

$$\rho_D^{-1} = \begin{pmatrix} [\mathbb{1}_{\check{D}_0} \otimes_{\mathbb{C}} 1]^\wedge & 0 \\ [\frac{\check{d}_0(a, x) \otimes_{\mathbb{C}} 1 - \check{d}_0(a, b) \otimes_{\mathbb{C}} 1}{x-b}]^\wedge & 0 \\ 0 & [\mathbb{1}_{\check{D}_1} \otimes_{\mathbb{C}} 1]^\wedge \\ 0 & [\frac{\check{d}_1(a, b) \otimes_{\mathbb{C}} 1 - \check{d}_1(a, x) \otimes_{\mathbb{C}} 1}{x-b}]^\wedge \end{pmatrix} : D \longrightarrow D \otimes I, \quad (\text{A.16})$$

and the relevant homotopy is

$$\psi_D^r = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & [\mathbb{1}_{\check{D}_1} \otimes_{\mathbb{C}} \frac{r_{xb}-1}{b-x}]^\wedge & 0 \\ 0 & 0 & 0 & 0 \\ -[\mathbb{1}_{\check{D}_0} \otimes_{\mathbb{C}} \frac{r_{xb}-1}{b-x}]^\wedge & 0 & 0 & 0 \end{pmatrix}. \quad (\text{A.17})$$

From the above it also immediately follows that $\lambda_I = \rho_I$ as we have

$$\rho_I \circ \lambda_I^{-1} = \begin{pmatrix} (\mathbb{1}_{I_0} \otimes_R \mu) \circ [1 \otimes_{\mathbb{C}} \mathbb{1}_{\check{I}_0}]^\wedge & 0 \\ 0 & (\mathbb{1}_{I_1} \otimes_R \mu) \circ [\frac{1 \otimes_{\mathbb{C}} \check{t}_1(a, b) - 1 \otimes_{\mathbb{C}} \check{t}_1(x, b)}{a-x}]^\wedge \end{pmatrix}, \quad (\text{A.18})$$

and it is straightforward to check that the right-hand side acts as the identity. Thus we have $\rho_I \circ \lambda_I^{-1} = \text{id}_I$, and hence $\rho_I = \lambda_I$.

Many-variable case

To check that the left and right unit maps (2.38) are indeed morphisms in $\text{DG}_{\text{bi}}(W)$, that they are natural and that they satisfy the triangle axiom is straightforward. Hence to prove theorem 2.18 we only need to show that λ_D and ρ_D are isomorphisms in $\text{MF}_{\text{bi}}(W)$. Instead of constructing explicit homotopies as in the one-variable case, in the following we will argue more abstractly and compactly. The price to pay is that the argument only works in $\text{MF}_{\text{bi}}(W)_{\text{f}}$.

Let the \mathbb{C} -vector space $H(D)$ be the cohomology of a matrix (bi-)factorisation D as defined in [KhR, sec. 3]. $H(D)$ is finite-dimensional iff $D \in \text{MF}_{\text{bi}}(W)_{\text{f}}$ [KhR, cor. 5].

Two morphisms F and G in $\text{Hom}_{\text{MF}_{\text{bi}}(W)}(D, D)$ induce linear maps $H(F)$ and $H(G)$ from $H(D)$ to itself. If we additionally have $F \circ G = \text{id}_D$, then we also have $H(F) \circ H(G) = \text{id}_{H(D)}$. If $D \in \text{MF}_{\text{bi}}(W)_{\text{f}}$ then $H(G)$ must be an isomorphism of vector spaces. But then according to [KhR, prop. 8] also G is an isomorphism in $\text{MF}_{\text{bi}}(W)$, and because of $F \circ G = \text{id}_D$ we find $F = G^{-1}$.

Let us now apply this reasoning to the case at hand. The inverses $\lambda_D^{-1} : D \rightarrow I \otimes D$ and $\rho_D^{-1} : D \rightarrow D \otimes I$ are most easily presented with the help of exterior algebras. However, we need not introduce this language here since for our purposes it is sufficient to state the result that these morphisms are of the form

$$\lambda_D^{-1} = \begin{pmatrix} [1 \otimes_{\mathbb{C}} \mathbb{1}_{D_0}]^{\wedge} & 0 \\ * & 0 \\ \vdots & \vdots \\ * & 0 \\ \hline * & 0 \\ \vdots & \vdots \\ * & 0 \\ \hline 0 & * \\ \vdots & \vdots \\ 0 & * \\ \hline 0 & [1 \otimes_{\mathbb{C}} \mathbb{1}_{D_1}]^{\wedge} \\ 0 & * \\ \vdots & \vdots \\ 0 & * \end{pmatrix}, \quad \rho_D^{-1} = \begin{pmatrix} [1 \otimes_{\mathbb{C}} \mathbb{1}_{D_0}]^{\wedge} & 0 \\ * & 0 \\ \vdots & \vdots \\ * & 0 \\ \hline * & 0 \\ \vdots & \vdots \\ * & 0 \\ \hline 0 & [1 \otimes_{\mathbb{C}} \mathbb{1}_{D_1}]^{\wedge} \\ 0 & * \\ \vdots & \vdots \\ 0 & * \\ \hline 0 & * \\ \vdots & \vdots \\ 0 & * \end{pmatrix}.$$

One easily checks that $\lambda_D \circ \lambda_D^{-1} = \text{id}_D$ and $\rho_D \circ \rho_D^{-1} = \text{id}_D$.

In [KhR, prop. 23] and [BR, sec. 5] the existence of isomorphisms $\varphi_D^l : I \otimes D \rightarrow D$ and $\varphi_D^r : D \otimes I \rightarrow D$ was shown. Thus, if we set $F = \lambda_D \circ (\varphi_D^l)^{-1}$ and $G = \varphi_D^l \circ \lambda_D^{-1}$ (or $F = \rho_D \circ (\varphi_D^r)^{-1}$ and $G = \varphi_D^r \circ \rho_D^{-1}$) we find by the above argument that λ_D and ρ_D are indeed isomorphisms.

A.2. Some categories for $\mathcal{N} = 2$ minimal models

In this appendix we collect the ingredients necessary to compute the fusing matrices for the category $\mathcal{D}_d^{\mathcal{N}=2}$ describing the defects of $\mathcal{N} = 2$ minimal models and sketch the computation of these fusing matrices. We do not provide all details, but we have tried to include sufficient references to places in the literature where more details can be found.

Fusing and braiding matrices

Let \mathcal{C} be a braided fusion category over \mathbb{C} (see e. g. [BK]), that is, a \mathbb{C} -linear semi-simple rigid braided monoidal category with finitely many isomorphism classes of simple objects, finite-dimensional morphism spaces, and such that the unit object $\mathbf{1}$ of \mathcal{C} is simple. Denote the monoidal structure by $(\otimes, \alpha, \mathbf{1}, \lambda, \rho)$ as in definition 2.4 and the braiding isomorphisms by $c_{U,V} : U \otimes V \rightarrow V \otimes U$. The categories presented below are also balanced, and we denote the twist isomorphisms by $\theta_U : U \rightarrow U$. (In fact, these categories are even modular, but we will not use this.)

Select a set $\{U_i \mid i \in \mathcal{I}\}$ of representatives of the isomorphism classes of simple objects in \mathcal{C} such that $U_0 = \mathbf{1}$. For notational simplicity we will assume that $\text{Hom}(U_i \otimes U_j, U_k)$ is either zero- or one-dimensional. This is the case for the categories we will consider below. For every triple $i, j, k \in \mathcal{I}$ such that $\text{Hom}(U_i \otimes U_j, U_k) \neq 0$, pick non-zero vectors

$$\lambda_{ij}^k \in \text{Hom}(U_i \otimes U_j, U_k), \quad \bar{\lambda}_k^{ij} \in \text{Hom}(U_k, U_i \otimes U_j), \quad (\text{A.19})$$

such that $\lambda_{ij}^k \circ \bar{\lambda}_k^{ij} = \text{id}_{U_k}$. For $k \neq l$ we have $\text{Hom}(U_k, U_l) = 0$ and so automatically $\lambda_{ij}^k \circ \bar{\lambda}_l^{ij} = 0$. If either i or j are zero, we impose the normalisation conditions

$$\lambda_{0i}^i = \lambda_{U_i} : \mathbf{1} \otimes U_i \longrightarrow U_i, \quad \lambda_{i0}^i = \rho_{U_i} : U_i \otimes \mathbf{1} \longrightarrow U_i. \quad (\text{A.20})$$

■ **Fusing matrices:** The fusing matrices implement a basis transformation between two sets of basis vectors $(U_i \otimes U_j) \otimes U_k \rightarrow U_l$ constructed from the λ_{ij}^k ,

$$\lambda_{ip}^l \circ (\text{id}_{U_i} \otimes \lambda_{jk}^p) \circ \alpha_{U_i, U_j, U_k} = \sum_{q \in \mathcal{I}} F_{pq}^{(ijk)l} \cdot \lambda_{qk}^l \circ (\lambda_{ij}^q \otimes \text{id}_{U_k}). \quad (\text{A.21})$$

Similarly, one can use the $\bar{\lambda}_k^{ij}$ to construct two bases of $U_l \rightarrow U_i \otimes (U_j \otimes U_k)$. This gives

$$\alpha_{U_i, U_j, U_k} \circ (\bar{\lambda}_q^{ij} \otimes \text{id}_{U_k}) \circ \bar{\lambda}_l^{qk} = \sum_{p \in \mathcal{I}} F_{pq}^{(ijk)l} \cdot (\text{id}_{U_i} \otimes \bar{\lambda}_p^{jk}) \circ \bar{\lambda}_l^{ip} \quad (\text{A.22})$$

for the same numbers $F_{pq}^{(ijk)l}$ as in (A.21).

■ **Braiding matrices:** The braiding matrices $R^{(ij)k} \in \mathbb{C}^\times$ encode the action of the braiding isomorphisms on simple objects. There are the following two equivalent expressions defining the $R^{(ij)k}$:

$$\lambda_{ji}^k \circ c_{U_i, U_j} = R^{(ij)k} \cdot \lambda_{ij}^k, \quad c_{U_i, U_j} \circ \bar{\lambda}_k^{ij} = R^{(ij)k} \cdot \bar{\lambda}_k^{ji}. \quad (\text{A.23})$$

■ **Twist eigenvalues:** The twist eigenvalues are simply the constants $\theta_k \in \mathbb{C}^\times$ determined by $\theta_{U_k} = \theta_k \cdot \text{id}_{U_k}$. They are well-defined because U_k is simple, which under the present assumptions implies that it is also absolutely simple, $\text{Hom}(U_k, U_k) = \mathbb{C} \cdot \text{id}_{U_k}$.

Rational free boson

Here we describe the category $\mathcal{C}_{2N}^{\mathfrak{u}(1)}$. It arises as the category of representations of the rational vertex operator algebra obtained when extending the Heisenberg vertex operator algebra (the $\widehat{\mathfrak{u}}(1)$ -current algebra) by two fields of weight N . Its fusing matrices were computed in [BS], here we use the conventions of [FGRS, app. B].

The index set for the simple objects U_k is $\mathcal{I} = \{0, \dots, 2N-1\}$. The fusion rules are $U_m \otimes U_n \cong U_{[m+n]}$, where $m, n \in \mathcal{I}$ and $[m+n] = m+n \pmod{2N} \in \mathcal{I}$. The twist eigenvalues, braiding matrices and fusing matrices are given by

$$\theta_k = e^{-\pi i k^2 / (2N)}, \quad R^{(kl)[k+l]} = e^{-\pi i k l / (2N)}, \quad F_{[s+t][r+s]}^{(rst)[r+s+t]} = (-1)^{r \sigma(s+t)}, \quad (\text{A.24})$$

where $k, l, r, s, t \in \mathcal{I}$, and $\sigma(k+l)$ is 0 if $k+l < 2N$ and 1 otherwise.

Affine $\mathfrak{su}(2)$

Here we describe the category $\mathcal{C}_k^{\mathfrak{su}(2)}$. It arises as the category of integrable highest weight representations of $\widehat{\mathfrak{su}}(2)_k$, or as a semi-simple quotient of the category of representations of $U_q(\mathfrak{sl}(2))$ with $q = e^{\pi i / (k+2)}$. Its data was given in [HSWY, KiR]; we use the conventions of [FRS1, sec. 2.5.2].

The index set for the simple objects U_i is $\mathcal{I} = \{0, \dots, k\}$. The fusion rules are

$$U_i \otimes U_j \cong \bigoplus_{l=|i-j|}^{\min(i+j, 2k-i-j)} \overset{(+2)}{U_l}. \quad (\text{A.25})$$

where the superscript $(+2)$ means that the direct sum is carried out in steps of two. The twist eigenvalues and braiding matrices are given by

$$\theta_r = e^{-2\pi i \Delta_r}, \quad R^{(rs)t} = (-1)^{(r+s-t)/2} e^{-i\pi(\Delta_t - \Delta_r - \Delta_s)}, \quad (\text{A.26})$$

where $\Delta_n = n(n+2)/(4k+8)$. The fusing matrices are

$$F_{pq}^{(rst)u} = \left\{ \begin{array}{ccc} t/2 & s/2 & p/2 \\ r/2 & u/2 & q/2 \end{array} \right\} \quad (\text{A.27})$$

with

$$\begin{aligned}
\left\{ \begin{array}{ccc} a & b & e \\ d & c & f \end{array} \right\} &= (-1)^{a+b-c-d-2e} \sqrt{[2e+1][2f+1]} \Delta(a, b, e) \Delta(a, c, f) \\
&\cdot \Delta(c, e, d) \Delta(d, b, f) \sum_z (-1)^z [z+1]! \left([z-a-b-e]! \right. \\
&\cdot [z-a-c-f]! [z-b-d-f]! [z-d-c-e]! \\
&\cdot \left. [a+b+c+d-z]! [a+d+e+f-z]! [b+c+e+f-z]! \right)^{-1} \quad (\text{A.28})
\end{aligned}$$

and

$$\Delta(a, b, c) := \sqrt{[-a+b+c]! [a-b+c]! [a+b-c]! / [a+b+c+1]!}. \quad (\text{A.29})$$

The symbols $[n]$ and $[n]!$ stand for q -numbers and q -factorials, respectively, i. e.

$$[n] = \frac{\sin\left(\frac{\pi n}{k+2}\right)}{\sin\left(\frac{\pi}{k+2}\right)}, \quad [n]! = \prod_{m=1}^n [m], \quad [0]! = 1. \quad (\text{A.30})$$

The range of the z -summation is such that the arguments are non-negative, i. e. z runs over all integers (in steps of 1) from $\max(a+b+e, a+c+f, b+d+f, c+d+e)$ to $\min(a+b+c+d, a+d+e+f, b+c+e+f)$.

$\mathcal{N} = 2$ minimal models

The bosonic part $(\mathfrak{sVir}_d)_{\text{bos}}$ of the $\mathcal{N} = 2$ super-Virasoro vertex operator algebra with central charge $c_d = 3 - \frac{6}{d}$ can be obtained from the coset construction for $(\widehat{\mathfrak{su}}(2)_{d-2} \oplus \widehat{\mathfrak{u}}(1)_4) / \widehat{\mathfrak{u}}(1)_{2d}$ [KS]. Using the results of [Fr1, Fr2] to obtain the representation category of the coset theory from those of the component theories, one finds that the category of $(\mathfrak{sVir}_d)_{\text{bos}}$ -representations can be described as

$$\mathcal{C}_d^{\mathcal{N}=2} = \left(\mathcal{C}_{d-2}^{\widehat{\mathfrak{su}}(2)} \boxtimes \bar{\mathcal{C}}_{2d}^{\widehat{\mathfrak{u}}(1)} \boxtimes \mathcal{C}_4^{\widehat{\mathfrak{u}}(1)} \right)_C^{\text{loc}}. \quad (\text{A.31})$$

Here, the overline in $\bar{\mathcal{C}}_{2d}^{\widehat{\mathfrak{u}}(1)}$ means that we replace the braiding and twist by their inverses, and \boxtimes is the product of \mathbb{C} -linear categories. That is, one takes triples $U_l \times U_m \times U_s$ of objects and tensor products (over \mathbb{C}) of morphism spaces, and completes the resulting category with respect to direct sums of objects. Let us abbreviate the simple objects as

$$U_{l,m,s} \equiv U_l \times U_m \times U_s \in \mathcal{C}_{d-2}^{\widehat{\mathfrak{su}}(2)} \boxtimes \bar{\mathcal{C}}_{2d}^{\widehat{\mathfrak{u}}(1)} \boxtimes \mathcal{C}_4^{\widehat{\mathfrak{u}}(1)}. \quad (\text{A.32})$$

The subscript C in (A.31) stands for the object

$$C = \mathbf{1} \oplus J \quad \text{with} \quad J = U_{d-2,d,2}. \quad (\text{A.33})$$

The object J obeys $J \otimes J \cong \mathbf{1}$ and $\theta_J = \text{id}_J$. It follows from [FRS3, def. 3.17 & prop. 3.22] that C carries a unique (up to isomorphism) structure of a commutative simple special symmetric Frobenius algebra (see [Fr1, sec. 2] for definitions). Finally, the notation $(\cdot)_C^{\text{loc}}$ in (A.31) denotes the category of local C -modules (see [Fr1, sec. 3.4] for a definition of local modules). By [KO] (see also [Fr1, prop. 3.21]), this is again a braided monoidal category.

We proceed to compute the category $\mathcal{C}_d^{\mathcal{N}=2}$. Let $\text{Ind}_C(V) = C \otimes V$ denote the C -module induced by V . The C -action on $C \otimes V$ is given by the multiplication of C . As there is no object $U_{l,m,s}$ for which $U_{l,m,s} \otimes J \cong U_{l,m,s}$, the induced modules $\text{Ind}_C(U_{l,m,s})$ are all simple, and we have

$$\begin{aligned} \text{Ind}_C(U_{l,m,s}) &\cong \text{Ind}_C(U_{l',m',s'}) \\ \Leftrightarrow (l, m, s) &= (l', m', s') \text{ or } (l, m, s) = (d-2-l', d+m', 2+s'). \end{aligned} \quad (\text{A.34})$$

Every C -module is a submodule of an induced module [FS, lem. 4.15], so that we have found all simple C -modules. All other C -modules are isomorphic to direct sums of these, as the category of C -modules is semi-simple [FS, prop. 5.24]. A simple C -module M is local iff $\theta_M = \xi_M \cdot \text{id}_M$ for some $\xi_M \in \mathbb{C}^\times$ [Fr1, cor. 3.18]. In the present case, this means that $\text{Ind}_C(U_{l,m,s})$ is local iff $\theta_{l,m,s} = \theta_{d-2-l, d+m, 2+s}$. One checks that

$$\text{Ind}_C(U_{l,m,s}) \text{ local} \quad \Leftrightarrow \quad l + m + s \in 2\mathbb{Z}. \quad (\text{A.35})$$

Altogether, we see that the isomorphism classes of simple objects in $\mathcal{C}_d^{\mathcal{N}=2}$ are labelled by \mathbb{Z}_2 -orbits in the set of labels (l, m, s) for which $l + m + s \in 2\mathbb{Z}$, with \mathbb{Z}_2 -action given by $(l, m, s) \mapsto (d-2-l, d+m, 2+s)$. Let $[l, m, s]$ denote a fixed choice of representative on each orbit, i. e. if $[l, m, s] = (l', m', s')$ then also $[d-2-l, d+m, 2+s] = (l', m', s')$. We denote the simple objects of $\mathcal{C}_d^{\mathcal{N}=2}$ by

$$R_{l,m,s} := \text{Ind}_C(U_{l,m,s}). \quad (\text{A.36})$$

Let \mathcal{E}_d be the full subcategory of $\mathcal{C}_{d-2}^{\text{su}(2)} \boxtimes \bar{\mathcal{C}}_{2d}^{\text{u}(1)} \boxtimes \mathcal{C}_4^{\text{u}(1)}$ consisting of all objects that are isomorphic to direct sums of simple objects $U_{l,m,s}$ with $l+m+s \in 2\mathbb{Z}$. It is easy to see that \mathcal{E}_d is a monoidal subcategory. By definition of the tensor product and the braiding of C -modules [Fr1, sec. 3.4], the induced modules provide a braided monoidal functor

$$\text{Ind}_C(\cdot) : \mathcal{E}_d \longrightarrow \mathcal{C}_d^{\mathcal{N}=2}. \quad (\text{A.37})$$

We have already seen that this functor is essentially surjective, however by (A.34) it is not full. (It is, however, injective on morphisms.)

In making the connection to matrix factorisations, we are primarily interested in the simple objects $U_{l,m,0}$ and their direct sums. Let $\mathcal{E}_{d,s=0}$ be the full subcategory of \mathcal{E}_d consisting of all objects isomorphic to direct sums of simple objects of the form $U_{l,m,0}$. This is again a monoidal subcategory, in fact we can canonically identify

$$\mathcal{E}_{d,s=0} = \mathcal{C}_{d-2}^{\text{su}(2)} \boxtimes \bar{\mathcal{C}}_{2d}^{\text{u}(1)}. \quad (\text{A.38})$$

Denote by $\mathcal{C}_{d,s=0}^{\mathcal{N}=2}$ the full subcategory of $\mathcal{C}_d^{\mathcal{N}=2}$ consisting of objects isomorphic to $\text{Ind}_C(V)$ for some $V \in \mathcal{E}_{d,s=0}$. Because (A.34) never relates two modules in $\mathcal{C}_{d,s=0}^{\mathcal{N}=2}$, we now have a braided monoidal equivalence

$$\text{Ind}_C(\cdot) : \mathcal{E}_{d,s=0} \xrightarrow{\sim} \mathcal{C}_{d,s=0}^{\mathcal{N}=2}. \quad (\text{A.39})$$

Defects for $\mathcal{N} = 2$ minimal models

Let G be the object $G = R_{0,0,2} \in \mathcal{C}_d^{\mathcal{N}=2}$ and set $S = \mathbf{1} \oplus G$. There is a unique (up to isomorphism) structure of a simple special symmetric Frobenius algebra on S [FRS3, def. 3.17 & prop. 3.22]. This Frobenius algebra is non-commutative. The defect category $\mathcal{D}_d^{\mathcal{N}=2}$ is given by the category of S -bimodules in $\mathcal{C}_d^{\mathcal{N}=2}$,

$$\mathcal{D}_d^{\mathcal{N}=2} = S\text{-mod-}S. \quad (\text{A.40})$$

The category $\mathcal{D}_d^{\mathcal{N}=2}$ is again semi-simple (apply [FS, prop. 5.24] to $(S \otimes S^{\text{op}})$ -modules). Given an object $V \in \mathcal{C}_d^{\mathcal{N}=2}$ and an algebra automorphism $\psi \in \text{Aut}(S)$, denote by $\alpha_S^+(V)_\psi$ the S -bimodule obtained as follows (see [FRS1, sec. 5.4] for definitions and references). The underlying object is $S \otimes V$, the left action is by multiplication, $\rho^l = m \otimes \text{id}_V$ (we omit the coherence isomorphisms), and for the right action one uses the braiding and precomposes the action by ψ ,

$$\rho^r = (m \otimes \text{id}_V) \circ (\text{id}_S \otimes c_{V,S}) \circ (\text{id}_S \otimes \text{id}_V \otimes \psi). \quad (\text{A.41})$$

The group $\text{Aut}(S)$ is isomorphic to \mathbb{Z}_2 , its non-trivial element is the isomorphism $\omega = \text{id}_1 - \text{id}_G$. We have, for $\psi \in \text{Aut}(S)$,

$$\alpha_S^+(V)_\psi \cong \alpha_S^+(G \otimes V)_{\omega \circ \psi}. \quad (\text{A.42})$$

This follows from [Fr3, def. 5.2 & prop. 5.10] and the fact that $\theta_G = -\text{id}_G$. If V is simple, then $\alpha_S^+(V)_\psi$ is simple as an S -bimodule and $\alpha_S^+(V)_\psi \cong \alpha_S^+(V)_\xi$ iff $\psi = \xi$, see [Fr3, eq. (5.19) & prop. 5.9]. In particular, $\alpha_S^+(V)_{\text{id}} \not\cong \alpha_S^+(G \otimes V)_{\text{id}}$, and altogether, for V and V' simple,

$$\alpha_S^+(V)_{\text{id}} \cong \alpha_S^+(V')_{\text{id}} \Leftrightarrow V \cong V'. \quad (\text{A.43})$$

Braided induction (with trivial automorphism) provides a monoidal functor

$$\alpha_S^+(\cdot)_{\text{id}} : \mathcal{C}_d^{\mathcal{N}=2} \longrightarrow \mathcal{D}_d^{\mathcal{N}=2}. \quad (\text{A.44})$$

Every S -bimodule is a subbimodule of an $\alpha_S^+(W)_\psi$ for some $W \in \mathcal{C}_d^{\mathcal{N}=2}$, $\psi \in \text{Aut}(S)$ [Fr3, prop. 5.7]. In fact, by (A.42), every S -bimodule is a subbimodule of an $\alpha_S^+(W)_{\text{id}}$ for some $W \in \mathcal{C}_d^{\mathcal{N}=2}$, so that $\alpha_S^+(\cdot)_{\text{id}}$ is essentially surjective. Finally, (A.43) implies that $\alpha_S^+(\cdot)_{\text{id}}$ is fully faithful, and hence an equivalence.

Let us write $D_{l,m,s}$ for $\alpha_S^+(R_{l,m,s})_{\text{id}}$. Let $\mathcal{D}_{d,s=0}^{\mathcal{N}=2}$ denote the full subcategory of $\mathcal{D}_d^{\mathcal{N}=2}$ consisting of objects isomorphic to direct sums of simple objects of the form

$D_{l,m,0}$. Combining the equivalences (A.38), (A.39) and (A.44) we get a monoidal equivalence

$$\mathcal{F} : \mathcal{C}_{d-2}^{\mathfrak{su}(2)} \boxtimes \bar{\mathcal{C}}_{2d}^{\mathfrak{u}(1)} \longrightarrow \mathcal{D}_{d,s=0}^{\mathcal{N}=2}. \quad (\text{A.45})$$

This makes it easy to compute the fusing matrices of $\mathcal{D}_{d,s=0}^{\mathcal{N}=2}$, namely they are just the product of those of $\mathcal{C}_{d-2}^{\mathfrak{su}(2)}$ and $\mathcal{C}_{2d}^{\mathfrak{u}(1)}$ (the overline does not affect the fusing matrices). In more detail, define the non-zero morphisms

$$\lambda_{(i,a)(j,b)}^{(k,c)} \in \text{Hom}_{\mathcal{D}_d^{\mathcal{N}=2}}(D_{i,a,0} \otimes D_{j,b,0}, D_{k,c,0}), \quad \lambda_{(i,a)(j,b)}^{(k,c)} = \mathcal{F}(\lambda[\mathfrak{su}(2)]_{ij}^k \times \lambda[\mathfrak{u}(1)]_{ab}^c) \quad (\text{A.46})$$

where $\lambda[\mathfrak{su}(2)]_{ij}^k$ refers to the basis (A.19) chosen for $\mathcal{C}_{d-2}^{\mathfrak{su}(2)}$ and $\lambda[\mathfrak{u}(1)]_{ab}^c$ to that chosen for $\mathcal{C}_{2d}^{\mathfrak{u}(1)}$. Then

$$\begin{aligned} & \lambda_{(i,a)(p,e)}^{(l,d)} \circ (\text{id} \otimes \lambda_{(j,b)(k,c)}^{(p,e)}) \circ \alpha \\ &= \sum_{q=0}^{d-2} \sum_{f=0}^{2d-1} F[\mathfrak{su}(2)]_{pq}^{(ijk)l} \cdot F[\mathfrak{u}(1)]_{ef}^{(abc)d} \cdot \lambda_{(q,f)(k,c)}^{(l,d)} \circ (\lambda_{(i,a)(j,b)}^{(q,f)} \otimes \text{id}), \end{aligned} \quad (\text{A.47})$$

where $F[\mathfrak{su}(2)]$ stands for the fusing matrix (A.27), and $F[\mathfrak{u}(1)]$ for the fusing matrix (A.24).

Spectra of defects fields

For two objects $D, E \in \mathcal{D}_d^{\mathcal{N}=2}$, let $H_{\text{bos}}^{D \rightarrow E}$ be the space of bosonic defect changing fields that change the defect labelled D to the defect labelled E . It decomposes as

$$H_{\text{bos}}^{D \rightarrow E} = \bigoplus_{[l,m,s],[l',m',s']} (R_{[l,m,s]} \otimes \bar{R}_{[l',m',s']})^{M_{l,m,s;l',m',s'}}, \quad (\text{A.48})$$

where the multiplicities are given by the dimension of a space of S -bimodule isomorphisms [Fr3, sec. 2.2]

$$M_{l,m,s;l',m',s'} = \dim \text{Hom}_{S\text{-mod}S}(R_{l,m,s} \otimes^+ D \otimes^- R_{l',m',s'}, E). \quad (\text{A.49})$$

The definition of \otimes^\pm is given in [Fr3, sec. 2.1]. From (A.24) and (A.37) it follows that the double braiding of G with a simple object $R_{l,m,s} \in \mathcal{C}_d^{\mathcal{N}=2}$ is given by

$$c_{R_{l,m,s},G} \circ c_{G,R_{l,m,s}} = (-1)^s \cdot \text{id}_G \otimes \text{id}_{R_{l,m,s}}. \quad (\text{A.50})$$

If $D = \alpha_S^+(U)_{\text{id}}$ for some U , then from [Fr3, prop. 5.9] together with (A.42) and (A.50) we get

$$R_{l,m,s} \otimes^+ \alpha_S^+(U)_{\text{id}} \otimes^- R_{l',m',s'} \cong \alpha_S^+(R_{l,m,s} \otimes U \otimes R_{l',m',s'} \otimes G^{\otimes s'})_{\text{id}}. \quad (\text{A.51})$$

If now $D = \alpha_S^+(U)_{\text{id}}$ and $E = \alpha_S^+(V)_{\text{id}}$ for simple $U, V \in \mathcal{C}_d^{\mathcal{N}=2}$, then together with (A.43) this implies

$$M_{l,m,s;l',m',s'} = \dim \text{Hom}_{S\text{-mod}S}(\alpha_S^+(R_{l,m,s} \otimes U \otimes R_{l',m',s'} \otimes G^{\otimes s'})_{\text{id}}, \alpha_S^+(V)_{\text{id}})$$

$$= \dim \text{Hom}(R_{l,m,s} \otimes R_{l',m',-s'}, U^\vee \otimes V), \quad (\text{A.52})$$

where U^\vee is the dual of U . In the special case that $U = R_{0,2n,0}$ and $V = R_{u,2n+u,0}$ we have $U^\vee \otimes V \cong R_{u,u,0}$.

The above calculation leads to a space of bulk fields (or defect fields on the identity defect) which is a direct sum of terms $R_{l,m,s} \otimes_{\mathbb{C}} R_{l,-m,s}$. This is not quite what we want for the application in the main text, but it is related to that by reversing the sign of the anti-holomorphic $U(1)$ -current, which amounts to replacing

$$R_{l,m,s} \otimes_{\mathbb{C}} R_{l',-m',s'} \longmapsto R_{l,m,s} \otimes_{\mathbb{C}} R_{l',m',-s'}. \quad (\text{A.53})$$

This replacement has to be applied to (A.48) before comparing to the statements in section 3. (Alternatively, we could have started with a more complicated simple special symmetric Frobenius algebra S , which directly describes the correct modular invariant.)

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