

2.4. Codimension-1 Bifurcations from an equilibrium

Consider $\dot{x} = f(x, \mu)$, $x \in \mathbb{R}^n$, $\mu \in \mathbb{R}$ (*)

← smooth f throughout.

around an equilibrium $(x_0, \mu_0) = (0, 0)$ (i.e. $f(0, 0) = 0$). We say that

- (i) the system (*) undergoes a bifurcation of steady states, or a fold bifurcation around (x_0, μ_0) as μ crosses μ_0 , if there is new equilibrium appearing or existing equilibrium disappearing, as μ crosses μ_0 .

[HW] fold BF $\Rightarrow D_x f(x_0, \mu_0)$ has a zero eigenvalue.



The simplest / codimension-1 fold bifurcation is a bifurcation related to

$$D_x f(x_0, \mu_0) \simeq \begin{pmatrix} 0 & 0 \\ 0 & A_{n-1} \end{pmatrix} \quad \text{with } \sigma(A) \cap i\mathbb{R} = \emptyset \quad (**)$$

simple

We describe in detail several different codim-1 fold BFs:

- (i) Saddle-Node Bifurcation (generic scenario under (*)-(**))

Start by considering (*) in 1D:

$$\dot{x} = f(x, \mu), \quad x \in \mathbb{R}, \quad \mu \in \mathbb{R}. \quad (*_{1D})$$

Assume that $f(0, 0) = 0$ and $f_x(0, 0) = 0$ (codim-1 condition)

If the following non-degeneracy conditions hold:

$$(SN1) \quad f_{\mu}(0,0) \stackrel{=:a}{\neq} 0$$

$$(SN2) \quad \frac{1}{2} f_{xx}(0,0) \stackrel{=:b}{\neq} 0$$

then $(*)_{1D}$ is topo. equivalent (locally near $(0,0)$) to the normal form

$$\dot{x} = \text{sign}(a) \cdot \mu + \text{sign}(b) \cdot x^2.$$

Pf: Taylor expansion (f smooth) around $(0,0)$:

$$f(x, \mu) = \underbrace{f(0,0)}_{=0} + \underbrace{f_x(0,0)}_{=0} x + \underbrace{f_{\mu}(0,0)}_{\neq 0} \mu + \underbrace{\frac{1}{2} f_{xx}(0,0)}_{\neq 0} x^2 + \frac{1}{2} f_{x\mu}(0,0) x \mu + \frac{1}{2} f_{\mu\mu}(0,0) \mu^2 + O(|(x, \mu)|^3)$$

Moreover, $f_{\mu}(0,0) \neq 0 \xrightarrow{\text{I.F.T.}} f(x, \mu) = 0 \Leftrightarrow \mu = \mu(x) \text{ for } \mu(0) = 0$
 $\xrightarrow{\text{near } (0,0)}$

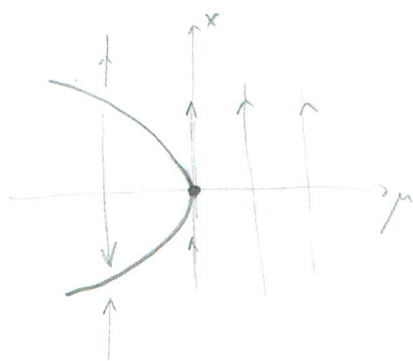
$$\Rightarrow f(x, \mu(x)) = 0 \Rightarrow f_x + f_{\mu} \cdot \mu' = 0 \Rightarrow \underbrace{f_x(0,0)}_{=0} + \underbrace{f_{\mu}(0,0)}_{\neq 0} \cdot \mu'(0) = 0 \Rightarrow \mu'(0) = 0.$$

Thus, $f(x, \mu) = \underbrace{f_{\mu}(0,0)}_{=:a \sim \pm 1} \mu + \underbrace{\frac{1}{2} f_{xx}(0,0)}_{=:b \sim \pm 1} x^2 + O(x^3).$

and $(*)_{1D}$ is topo. equiv. to

$$\dot{x} = \text{sign}(a) \cdot \mu + \text{sign}(b) \cdot x^2$$

e.g. $\dot{x} = \mu + x^2$



as μ crosses 0, the two equilibria $\pm \sqrt{\mu}$ (one stable, the other unstable) collide and disappear.

[HW] analyze other 3 cases: $\begin{matrix} \mu - x^2 \\ \mu + x^2 \\ -\mu - x^2 \end{matrix}$

Back to n-D.

$$\dot{x} = f(x, \mu), \quad x \in \mathbb{R}^n, \mu \in \mathbb{R}, \quad (*)$$

where $f(0,0)=0$ and $D_x f(0,0)$ has a simple eigenvalue 0.

Let v be an eigenvector of 0 s.t. $D_x f(0,0) \cdot v = 0$ and $\langle v, v \rangle = 1$, inner product in \mathbb{R}^n

let w be a left eigenvector of 0 s.t. $w^T D_x f(0,0) = 0$ and $\langle w, v \rangle = 1$.

⌈ If $w^T A = \lambda w^T$, $w \neq 0$, then w is a left eigenvector of λ for A .

Recall if v_1, \dots, v_n are eigenvectors of $\lambda_1, \dots, \lambda_n$, then

$$T^{-1} A T = J \leftarrow \text{Jordan form}, \quad T = (v_1, \dots, v_n)$$

Indeed, T^{-1} is composed of left eigenvectors as row vectors, i.e. ↑ ↑ ↑
columns of T

$$T^{-1} = \begin{pmatrix} w_1^T \\ \vdots \\ w_n^T \end{pmatrix} \leftarrow \begin{matrix} \leftarrow \\ \leftarrow \\ \leftarrow \end{matrix} \text{rows of } T^{-1}$$

$$\text{In particular, } \mathbb{1} = T^{-1} T \Rightarrow w_i^T v_j = \delta_{ij} = \begin{cases} 0 & i \neq j \\ 1 & i = j \end{cases}$$

The left and right eigenvectors chosen this way are called dual bases.

Based on Taylor expansion for multivariable functions:

$$f(x) = f(a) + Df(a)(x-a) + \underbrace{\frac{1}{2} (x-a)^T D^2 f(a) (x-a)}_{\text{bilinear form}} + O(|x-a|^3), \quad x, a \in \mathbb{R}^n$$

$$j\text{th component} = \frac{1}{2} \sum_{k,l=1}^n \frac{\partial^2 f}{\partial x_k \partial x_l} (x_k - a_k)(x_l - a_l)$$

We can replace (SN1)-(SN2) for 1D case with:

$$(SN_1) \quad \langle w, D_\mu f(0,0) \rangle \neq 0$$

$$(SN_2) \quad \langle w, \frac{1}{2} v^T D_{xx} f(0,0) v \rangle \neq 0$$

and state a similar result for (*):

Consider

$$\dot{x} = f(x, \mu), \quad x \in \mathbb{R}^n, \mu \in \mathbb{R} \quad (*)$$

Such that

$$f(0,0) = 0 \text{ and } D_x f(0,0) \text{ has a simple ev. } 0, \quad (**)$$

and

$$(SN_1) \quad \langle w, D_\mu f(0,0) \rangle \neq 0$$

$$(SN_2) \quad \langle w, \frac{1}{2} v^T D_{xx} f(0,0) v \rangle \neq 0.$$

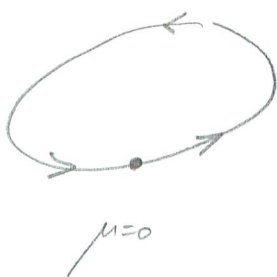
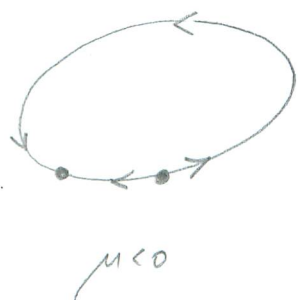
Then, (*) undergoes a saddle-node bifurcation around $x_0 = 0$ as μ crosses $\mu_0 = 0$, which is characterized by collision and disappearance of two equilibria with opposite stability of (*).

Remark: we also call the saddle-node bifurcation to be a generic phenomenon for systems of form (*) satisfying (**).

That is, if a system, say (*), under some condition given by finitely many equalities of derivatives, say (**), will definitely undergo certain phenomenon, saddle-node BF, unless further equalities of derivatives have to be satisfied, then we will call this phenomenon generic.

In other words, we can formulate the above result as:

"A system of form (*) satisfying condition (**) generically undergoes a saddle-node bifurcation."



Saddle-node (local) bifurcation on $W^c \simeq S^1$ leads to a new periodic orbit through homoclinic orbit. This is a global change of portraits. — global bifurcation.

Sometimes ^{different} setting of (*) may prevent (SN1) or (SN2) to be satisfied.

(iz) Transcritical Bifurcation (generic scenario for (*)-(**) with $f(0, \mu) \equiv 0 \forall \mu$)

In classical bifurcation theory, (*) is assumed to have a trivial solution from which bifurcation occurs, i.e.

$$f(0, \mu) \equiv 0 \quad \forall \mu \text{ near } \mu_0 = 0 \quad (T)$$

$\Rightarrow D_\mu f(0, 0) = 0 \Rightarrow \text{~~(SN1)~~},$ Instead, we assume

$$(T1) \quad \langle w, \frac{1}{2} v^T D_{\mu x} f(0, 0) v \rangle \neq 0$$

$$\frac{1}{2} f_{\mu x}(0, 0) \neq 0$$

in 1D.

$$(SN2) = (T2) \quad \langle w, \frac{1}{2} v^T D_{xx} f(0, 0) v \rangle \neq 0$$

$$\frac{1}{2} f_{xx}(0, 0) \neq 0$$

Consider 1D again, let $g(x, \mu) = f_x(x, \mu)$, then $g(0, \mu) = 0$, $g_x = f_{xx}$, $g_\mu = f_{\mu x}$

$$g(x, \mu) = 0 + \underbrace{g_x(0, 0)}_{\neq 0} x + \underbrace{g_\mu(0, 0)}_{\neq 0} \mu + \dots$$

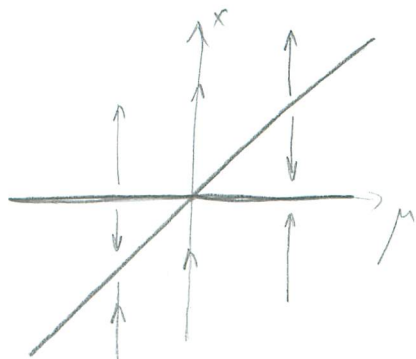
$$\left. \begin{array}{l} g_x(0, 0) \neq 0 \\ g_\mu(0, 0) \neq 0 \end{array} \right\} \Rightarrow \mu = O(x)$$

$$= \pm 2x \pm \mu$$

$$\Rightarrow f(x, \mu) = \pm x^2 \pm \mu x$$

e.g. $f(x, \mu) = x^2 + \mu x$

[HW] analyse the other 3 cases.



One shows in a similar way as for SN-BF, that:

"A system of form (*) satisfying conditions (**) and (T), generically undergoes a transcritical bifurcation, which is characterized by the exchange of stability of two equilibria as μ crosses $\mu_0=0$ " //

(i3) Pitchfork Bifurcation (generic scenario for (*)-(**) with $f(-x, \mu) = -f(x, \mu)$) \mathbb{Z}_2 -symmetry

Many systems possess symmetry (such as Duffing, Lorenz equations)

Assume

$$f(-x, \mu) = -f(x, \mu) \quad \forall x \in \mathbb{R}^n, \mu \in \mathbb{R} \quad (S)$$

$$\Rightarrow f \text{ is odd in } x \Rightarrow f_x \text{ is even in } x \Rightarrow f_{xx} \text{ is odd in } x \Rightarrow f_{xx}(0,0) = 0 \Rightarrow \text{SN2}$$

Also, (P) $\Rightarrow f(0, \mu) = 0 \quad \forall \mu \Rightarrow \text{SN1}$, Instead, we assume

$$(T1) = (P1) \quad \langle w, \frac{1}{2} D_{\mu\mu}^2 f(0,0) v \rangle \neq 0$$

$$(P2) \quad \langle w, \frac{1}{6} D_{xxx}^3 f(0,0) (v, v, v) \rangle \neq 0$$

$$\frac{1}{2} f_{\mu\mu} \neq 0$$

$$\frac{1}{6} f_{xxx} \neq 0$$

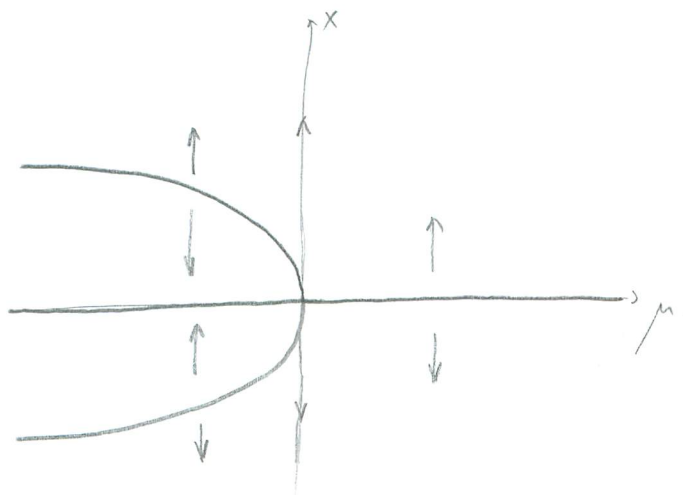
in 1D

Similarly, one shows in 1D, we have

$$f(x, \mu) = \pm \mu x \pm x^3$$

e.g.

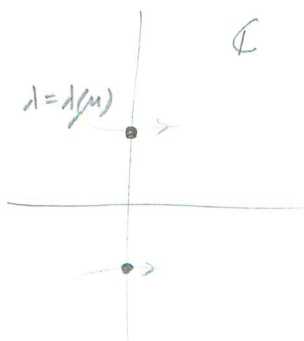
$$f(x, \mu) = \mu x + x^3 = x(\mu + x^2)$$



"A system of form (*) satisfying (**) and (S), generically undergoes a pitchfork bifurcation, which is characterized by appearance or disappearance of two equilibria with the same stability of (*) as μ crosses $\mu_0 = 0$ "

(ii) the system (*) undergoes a bifurcation of oscillating states, or a Hopf bifurcation around (x_0, μ_0) as μ crosses μ_0 , if there are non-constant periodic solutions appearing or disappearing, as μ crosses μ_0 . We assume (codim-1 BF)

$$D_x f(x_0, \mu_0) \simeq \left(\begin{array}{cc|c} 0 & -\omega & 0 \\ \omega & 0 & 0 \\ \hline 0 & 0 & A_{n-2} \end{array} \right), \quad \sigma(A) \cap i\mathbb{R} = \emptyset. \quad (**_H)$$



- $D_x f(x_0, \mu_0)$ remains invertible $\xrightarrow{\text{I.F.T.}} x = x(\mu) \ (x(0) = 0)$
unique equilibrium
for every μ near μ_0 .
- stability of the unique equilibrium
does change, as $\dim W^u$ and $\dim W^s$ change.

Aufgabe 11.2 \Rightarrow (*) satisfying $(**_H)$ in 2D can be brought to

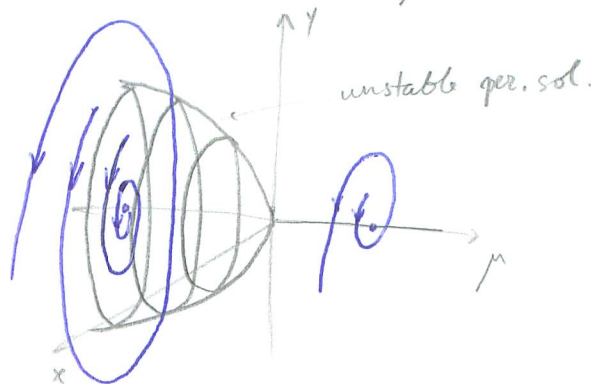
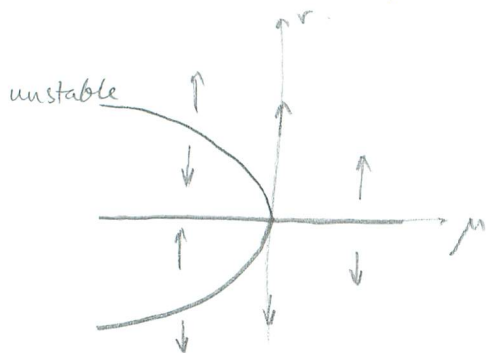
the normal form: (skipping the μ^2 and μ^2 terms in Aufgabe 11.2, why can it be skipped?)

$$\begin{cases} \dot{x} = (-\omega - b\mu)y + a\mu x + (cx - dy)(x^2 + y^2) \\ \dot{y} = a\mu y + (\omega + b\mu)x + (dx + cy)(x^2 + y^2) \end{cases} \quad a, b, c, d \in \mathbb{R}.$$

Polar- $\begin{cases} \dot{r} = (a\mu + cr^2)r \\ \dot{\theta} = \omega + b\mu + d r^2 \end{cases}$

generic conditions $a \neq 0$
 $c \neq 0 \Rightarrow \dot{r} = (\pm\mu \pm r^2)r$

e.g. $\dot{r} = (\mu + r^2)r$



For n-D, using center manifold + normal form, one shows the

generic conditions are:

(H1) $a = \frac{d}{d\mu} \operatorname{Re}(\lambda(\mu)) \Big|_{\mu=\mu_0=0} \neq 0$

(H2) $c = \frac{1}{16} (f_{xxx} + f_{xyy} + g_{xxy} + g_{yyx}) + \frac{1}{16\omega} [f_{xy}(f_{xx} + f_{yy}) - g_{xy}(g_{xx} + g_{yy}) - f_{xx}g_{xx} + f_{yy}g_{yy}]$
for the normal form on W^c by $\begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} = \begin{pmatrix} 0 & -\omega \\ \omega & 0 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} + \begin{pmatrix} f(x,y) \\ g(x,y) \end{pmatrix}, (x,y) \in E^c$

" A system of form (*) with $(**_H)$ generically undergoes a Hopf

bifurcation, which is characterized by a birth or disappearance of

a surface of per. solutions which lie on the paraboloid $\mu = -\frac{c}{a}(x^2 + y^2)$

up to the 3rd order of $|x,y|$. If $c > 0$, they are repelling; if $c < 0$, they are attracting.