

Based on joint work with

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and earlier work with

(hep-th/
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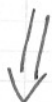
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Linear deformations of quaternionic-Kähler manifolds

① Intro & Motivation

$4n$ -dim. QK manifolds have holonomy groups contained in

$$Sp(n) \cdot Sp(1)$$



Metrics are Einstein

$$R_{ij}(g) = \Lambda g_{ij}$$

↑ "cosmological constant"

This implies

$R > 0$: compact

($R = 0$: Hyperkähler)

$R < 0$: Appears in "physics"
($N=2$ supergravity)

Typically non-compact.

Examples : non-compact versions
of Weyl spaces

$$\frac{Sp(m, 1)}{Sp(n) \times Sp(1)}$$

$$\frac{SO(m, 4)}{SO(n) \times SO(4)}$$

$$\frac{SU(m, 2)}{SU(n) \times SU(2) \times U(1)}$$

⋮

(Alekseevsky, de Wit & Van Proeyen, ...)

Question : what are the possible deformations of a QK metric preserving the quaternionic properties ?

For $R > 0$: there are none !
(complete)
(LeBrun, LeBrun & Salamon)

For $R < 0$: deformations are possible ! (Follows from string theory !)
↓
Mathematically not studied

This talk : study infinitesimal deformations

Method : study deformations of the twistor space of the associated Swann bundle
(hyperkähler)

dim	↑			<u>Ex</u>
4n+6		Z_S : Twistor space		$\mathbb{R}^8 \times \mathbb{C}P^1$
4n+4		S : Swann bundle		\mathbb{R}^8
4n+2		Z_{ell} : Twistor space	} $U(1) \times D$	$\mathbb{C}P^3$
4n		ell : Quaternion-Kähler	} $\frac{SU(2)}{U(1)}$	S^4

② Twistor spaces of HK

Swann bundle in Hyperkähler

→ 3 HK closed 2-forms

$$\omega^+ : (2,0)$$

$$\omega^3 : (1,1) \quad \text{Kähler-form}$$

$$\omega^- : (0,2)$$

Twistor space : $(2,0)$ form w.r.t any complex structure

$$\mathbb{F} \\ J(S, \bar{S})$$

$$S \in \mathbb{C}P^1$$

$$\Omega = \omega^+ - iS \omega^3 + S^2 \omega^- \quad (\text{closed})$$

(Hitchin, Karlhede, Lindström, Roček)
(HKLR)

Locally, \exists Darboux coordinates

$$\Omega = d\mu_I \wedge dv^I(s)$$

(Hypermultiplets:

∞ # of
auxiliary
fields)

$$\mu_I(s) = w_I + s w_{I,1} + \dots$$

$$v^I(s) = v^I + s v_{I,1}^I + \dots$$

$$\Omega = \underbrace{dw_I \wedge dv^I}_{\omega^+} + \underbrace{is \partial \bar{\partial} K}_{\omega^3} + \mathcal{F}^2 \omega^- + \mathcal{O}(\equiv) \text{ (truncates)}$$

"Solving the twistor lines" = finding Ω

\Rightarrow determine ω_1, ω_3

In general, solving the twistor lines is very difficult, unless there are extra symmetries. We will focus on

ell^{4n} : $n+1$ commuting isometries



$S^{4(n+1)}$: $n+1$ commuting triholomorphic isometries



$$v^I(s) = v^I + x^I s - \bar{v}^I s^2 \quad (\text{Lindstrom \& Roček})$$

($N=2$ tensor multiplet) $O(2)$

(*) Kähler potential

$$K(v, \bar{v}, w + \bar{w}) = \mathcal{L}(v, \bar{v}, x) - x^I (w + \bar{w})_I$$

with

$$(w + \bar{w})_I = \frac{\partial \mathcal{L}}{\partial x^I}$$

$$\mathcal{L} = \oint_C \frac{dS}{2\pi i S} H(\gamma, S)$$

(HKLR)

Isometries

$$W_I \rightarrow W_I + i \epsilon_I$$

Solution for the twistor line

$$\mu_I(s) = -\bar{\mu}_I\left(\frac{i}{s}\right) + \partial_{\eta^I} H(\eta, s)$$

with $\eta^I = \frac{1}{s} v^I = \frac{v^I}{s} + x^I - \bar{v}^I s$ (*)

Ex: $H = \frac{i}{2} \eta^2$

$$\mu = w + \bar{w} s$$

check:

$$w + \bar{v} s = -\left(\bar{w} + \frac{v}{s}\right) - \left(\frac{v}{s} + x - \bar{v} s\right)$$

$$w + \bar{w} = -x$$

$$\omega_3^{\mathbb{R}} = i(d\bar{w} \wedge dw + d\bar{v} \wedge dv) \quad \mathbb{R}^4$$

HK-geometry fixed by specifying

$$H(\eta^I, s)$$

For Swann-bundles: indep. of s and

$H(\eta^I)$ homogeneous of degree one

Ex

$$H(\eta^I) = \frac{F(\eta^A)}{\eta^b}$$

$$\eta^I = \{\eta^b, \eta^A\}$$

(Roček - Vafa - Vandoren) // C-map:
F: hom. degree 2

Deformations can be studied by deforming the twistor lines.

Deformations preserving the $n+1$ commuting isometries are of the type

$$H(\eta, s) = H^0(\eta, s) + H^1(\eta, s)$$

EX

$$H = \frac{F(\eta^a)}{\eta^b} + c \eta^b \log \eta^b \quad c \in \mathbb{R}$$

(string one-loop correction)

$$\eta^I \rightarrow \eta^I + 0$$

$$\mu_a^{\#} \rightarrow \mu_a^{\#} + 0$$

$$\mu_b \rightarrow \mu_b + c \ln \left[\frac{s - s_+}{s - s_-} \right]$$

$$s_{\pm} = \frac{x^b \mp z^b}{2\bar{v}^b} \quad z^2 = x^2 + 4v\bar{v}$$

(roots of v_{\bullet}^b)

$$K = \langle \text{Leg } \mathcal{L} \rangle_x$$

$$K(v, \bar{v}, w + \bar{w})$$

$$\mathcal{L}(v, \bar{v}, x) = \oint_{C_{+,-}} \frac{ds}{2\pi i s} H(\eta)$$

③ General infinitesimal deformations

Perturb twistor lines

$$V^I(s) = V_0^I(s) + \delta V^I(s)$$

$$\mu_I(s) = \mu_I^\circ(s) + \delta \mu_I(s)$$

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unperturbed
(known)

↑ "small"

Allow the function H to depend on

$$\rho_I \equiv \text{Im}[\mu_I^\circ(s)] = -i(w - \bar{w})_I + \dots$$

Presence of ρ_I will break isometry

$$w_I \rightarrow w_I + i\varepsilon_I$$

$$H(\eta, \rho, \varepsilon) = H^\circ(\eta^\circ, \rho) + \delta H^\bullet(\eta^\circ, \rho, \varepsilon)$$

Solve for the perturbations $\delta V^I, \delta \mu_I$

in terms of unperturbed data and δH .

Solution :

$$\delta v^I = i \oint \frac{ds'}{2\pi i s'} \frac{s^3 + s'^3}{s'(s'-s)} \delta H^I(\eta^0, s, s')$$

$$\delta \mu_I = \oint \frac{ds'}{2\pi i s'} \frac{s+s'}{2(s'-s)} G_I(\eta^0, s, s')$$

with

$$\delta H^I = \frac{\partial}{\partial s_I} (\delta H)$$

$$G_I = \frac{\partial}{\partial \eta^I} (\delta H) + H_{IJ}^0 \left(i \delta H^J + \frac{1}{s} \delta v^J \right)$$

Further (main) result : Kähler potential

$$K(u, \bar{u}, w, \bar{w}) = K_0(u, \bar{u}, w + \bar{w})$$

$$+ \delta K_{\#}(u, \bar{u}, w - \bar{w})$$

with

$$\delta K_{\#} = \oint \frac{ds}{2\pi i s} \delta H(\eta, s, \cancel{s})$$

Swann-
bundles

H: hom.
degree 1
in η .

(Penrose₁ transform)
twistor

Example :

Atiyah-Hitchin metric (HK)
asymptotically ($r \rightarrow \infty$) approaches
(negative mass) Taub-Nut.

TN:

$$\mathcal{L} = -\frac{1}{2} \oint_{C_\infty} \frac{dS}{2\pi i S} \eta \ln \eta + \frac{1}{m} \oint_{C_0} \frac{dS}{2\pi i S} \eta^2$$

$$\delta H = 4\eta \cos(4\varphi)$$

$$= 2\eta (e^{4i\varphi} + e^{-4i\varphi})$$

" \uparrow
instanton "

\uparrow
anti-instanton

(Coulomb-branch of
3D gauge theory)