

3. Special Geometry

Lecture 8

3.1.1 Middle dim cohomology of threefolds.

We will consider a family of Calabi-Yau threefolds

$\pi: X \rightarrow B$ over a small complex domain B

with fibers $\pi^{-1}(z) = X_z$.

* We will assume that the local family extends to

a toric variety M .

* We will further assume that $\dim M = h^{2,1}(X_{z_0}) = n$

for a smooth X_{z_0}

We fix a smooth X_{z_0}

3.1.1 Proposition: $H^3(X_{z_0}, \mathbb{C})$ carries a polarized Hodge structure of weight 3.

We have $H^3(X_{z_0}, \mathbb{C}) = H^{3,0} \oplus H^{2,1} \oplus H^{1,2} \oplus H^{0,3}$

The lattice is given by $H^3(X_{z_0}, \mathbb{Z})$, which

has a symplectic form given by

$$Q(u, v) = \int_{X_{z_0}} u \wedge v$$

We choose a symplectic basis $\{\alpha_I, \beta^J\}_{0 \leq I, J \leq n}$

satisfying $Q(\alpha_I, \beta^J) = \delta_I^J$, $Q(\alpha_I, \alpha_J) = 0 = Q(\beta^I, \beta^J)$

with respect to this basis the symplectic form

$$\text{is } Q = \begin{pmatrix} 0 & \mathbb{1}_{n+1} \\ -\mathbb{1}_{n+1} & 0 \end{pmatrix}$$

denote the dual homology basis by $\{A^I, B_J\}_{0 \leq I, B \leq n}$

3.2 | The period map:

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We define the period domain

$$\mathcal{D} = \{ [w] \in \mathbb{P}(H^3(X_{z_0}, \mathbb{C})) \mid \Re(w, w) = 0, i\Re(w, \bar{w}) > 0 \}$$

We choose a holomorphic three form of the fiber X_{z_0} and denote it by $\Omega(z)$, then the period map

$\mathcal{P}_0: B \rightarrow \mathcal{D}$ is defined by

$$\mathcal{P}_0(z) = \left[\sum_I \left(\int_{A^I} \Omega(x) \right) \alpha_I + \sum_J \left(\int_{B^J} \Omega(x) \right) \beta^J \right]$$

Using path independent identification $H_3(X_z, \mathbb{Z}) \cong H_3(X_{z_0}, \mathbb{Z})$

we may globalize this period map on B to M

by introducing a covering space \tilde{M} . We denote

the resulting period map $\mathcal{P}: \tilde{M} \rightarrow \mathcal{D}$

and assume $\Gamma \subset \mathrm{Sp}(2r+2, \mathbb{Z})$ as the covering group.

We denote by \mathcal{U} the restriction to \mathcal{D} of

the tautological line bundle $\mathcal{O}(-1)$ over

$\mathbb{P}(H^3(X_{z_0}, \mathbb{C}))$. We set $\mathcal{L} = \mathcal{P}^* \mathcal{U}$, i.e. the

pullback to \tilde{M} . The complex conjugate of \mathcal{L}

will be denoted by $\bar{\mathcal{L}}$.

The period integral may be considered as a

section of \mathcal{L} .

3.3 The prepotential.

recall that $\dim H^3(X_{2n}, \mathbb{C}) = h^{3,0} + h^{2,1} + h^{1,2} + h^{0,3}$
 $= 2 + 2h^{2,1} = 2 + 2n$

\mathcal{D} is a projectivization, hence $\dim \mathcal{D} = 2n+1$.

The symplectic form ω on $H^3(X_{2n}, \mathbb{C})$ induces

a one form on \mathcal{D} by $\theta := \langle d\omega, \omega \rangle$.

With this one form (\mathcal{D}, θ) becomes a $2n+1$ dim contact manifold.

Def. A contact manifold P is a smooth manifold P of odd dimension $2n+1$ which is equipped with a differential 1-form A that is non-degenerate in the sense that the wedge product $A \wedge (dA)^n$ does not vanish.

We have the following:

i) Locally the period map $\mathcal{P}_0: B \rightarrow \mathcal{D}$ is an embedding
 (Bogomolov, Tian, Todorov)

ii) $\theta|_{\mathcal{P}_0(B)} = 0$ (due to Griffiths' transversality)

remember $\Omega \in F^3 = H^{3,0}$
 $\nabla \Omega \in F^2 = H^{3,0} \oplus H^{2,1}$

A result of Bryant and Griffiths states that the image of the period map \mathcal{P}_0 can be recovered by half of the period integrals

$$\chi^I(z) = \int_{A^I} \Omega(z)$$

3.3.1.1 Proposition/Definition

i) The map $z \mapsto [X^0(z), \dots, X^n(z)] \in \mathbb{P}^n$ is a local isomorphism $B \rightarrow \mathbb{P}^n$

ii) Integrating $\Theta|_{P_0(B)} = 0$ on B we can write the other half of the period integrals

$$P_J(z) = \frac{\partial F(X^I)}{\partial X^J} \quad (J = 0, 1, \dots, r)$$

in terms of a holomorphic function $F(X)$ called prepotential

The function $F(X)$ is a holomorphic function of $X^0(z), \dots, X^r(z)$ and has the following homogeneity property

$$\sum_{I=0}^r X^I \frac{\partial F}{\partial X^I} = 2F(X)$$

This potential function exists locally for the small domain B . When globalizing to \tilde{M} the monodromy group plays a role (The group action is referred to as duality transformation in physics)

3.4 | The Griffiths-Yukawa cubic:

3.4.1 | Definition:

We define the Griffiths-Yukawa cubic to be a section of $\mathcal{L}^2 \otimes \text{Sym}^3 T^* \tilde{M}$, whose components are given by (fixing a section Ω of \mathcal{L})

$$C_{ijk}(z) = - \int_{X_2} \Omega(z) \wedge \partial_i \partial_j \partial_k \Omega(z), \quad \partial_i = \frac{\partial}{\partial z^i}$$

$i = 1, \dots, \dim \tilde{M} = n$

Proposition 3.4.2.

The components of the cubic are related to the prepotential by

$$C_{ijk} = \sum_{I,J,K=0}^n \frac{\partial x^I}{\partial z^i} \frac{\partial x^J}{\partial z^j} \frac{\partial x^K}{\partial z^k} \frac{\partial^3 F}{\partial x^I \partial x^J \partial x^K}$$

we denote by $C_{IJK} = \frac{\partial^3 F}{\partial x^I \partial x^J \partial x^K}$

3.5.1 The period matrix

Let $\Omega(z_0) \in H^{3,0} = F^3$,

$F^2 = H^{3,0} \oplus H^{2,1}$ $\dim F^2 = n+1$, we fix a

basis $\Omega_0, \Omega_1, \dots, \Omega_n$ of F^2 st $\Omega_0 \in F^3$

and define

$$\Pi = \begin{pmatrix} \int_{A^1} \Omega_0 & \int_{B^1} \Omega_0 \\ \int_{A^1} \Omega_1 & \vdots \\ \vdots & \int_{B^1} \Omega_n \\ \int_{A^1} \Omega_n & \int_{B^1} \Omega_n \end{pmatrix} \text{ as an } (n+1, 2n+2)$$

matrix which we call the period matrix, the first row coincides with the period integral

$$\Omega = \sum_I x^I \alpha_I + \sum_J \beta_J \beta_J$$

3.5.11 Proposition:

The following are properties of the Hodge filtration and the Hodge-Riemann bilinear relations

1) $\Pi Q \Pi^tr = 0$

2) $i \Pi Q \overline{\Pi^tr} > 0$

here z_j means that we decompose the

$(n+1) \times (n+1)$ hermitian matrix $i\pi \circlearrowleft \overline{\pi}^{\text{tr}}$

into a block form compatible with the filtration $F^3 C F^2$

then the first diagonal block (1×1) is positive definite

and the whole $(n+1) \times (n+1)$ hermitian matrix

has 1 positive and n negative eigenvalues.