

1.8] Solutions of the PF equation:

1.8.1] Singular points.

For the Legendre family \mathcal{E}_2 we derived the following differential equation:

$$\lambda(\lambda-1)\pi'' + (2\lambda-1)\pi' + \frac{1}{4}\pi = 0 \quad \text{(PF)}$$

i.e.
$$\pi'' + \frac{(2\lambda-1)\pi'}{\lambda(\lambda-1)} + \frac{1}{4\lambda(\lambda-1)}\pi = 0$$

which makes it obvious that $\lambda=0,1$ are regular singular points, what about $\lambda=\infty$

we define $x = \frac{1}{\lambda}$, $\frac{\partial}{\partial \lambda} = \frac{\partial x}{\partial \lambda} \frac{\partial}{\partial x} = -x^2 \frac{\partial}{\partial x}$.

The equation (PF) becomes:

$$-x^2 \frac{d}{dx} \left(-x^2 \frac{d}{dx} \right) \pi(x) + \frac{\left(\frac{2}{x}-1\right)}{\frac{1}{x}\left(\frac{1}{x}-1\right)} \frac{-x^2 d\pi}{dx} + \frac{1}{4 \frac{1}{x}\left(\frac{1}{x}-1\right)} \pi(x) = 0$$

$$\Leftrightarrow x^4 \pi''(x) + 2x^3 \pi'(x) - \frac{(2x^2-x)x^2}{(1-x)} \pi'(x) + \frac{x^2}{4(1-x)} \pi(x) = 0$$

$$\Rightarrow \pi''(x) + \frac{2}{x} \pi'(x) - \frac{(2x^2-x)}{(1-x)} \frac{1}{x^2} \pi'(x) + \frac{1}{4x^2(1-x)} \pi(x) = 0$$

$$\pi''(x) + \frac{2x - 2x^2 - 2x^2 - x}{(1-x)x^2} \pi'(x) + \frac{1}{4x^2(1-x)} \pi(x) = 0$$

$$\pi''(x) + \frac{x(1-4x)}{x^2(1-x)} \pi'(x) + \frac{1}{4x^2(1-x)} \pi(x) = 0$$

hence $x=0$ is a regular singular point

$\Leftrightarrow \lambda=\infty$ is.

1.8.2 | Solutions at $\lambda = 0$

We rewrite (PF) as

$$\mathcal{L} = \theta^2 - \lambda (\theta + 1/2)^2, \quad \text{where } \theta = \lambda \frac{d}{d\lambda} \quad (\text{Exercise})$$

we make an ansatz $\pi(\lambda) = \sum_{n=0}^{\infty} a_n \lambda^{n+\delta}$

$$\mathcal{L} \pi = 0 \quad \text{gives} \quad \sum_{n=0}^{\infty} a_n (n+\delta)^2 \lambda^{n+\delta} - \lambda \sum_{n=0}^{\infty} a_n (n+\delta+1/2)^2 \lambda^{n+\delta}$$

the lowest term gives the indicial equation $\delta^2 = 0$

i.e. we have a little problem, we can only find one solution as power series.

$$\text{let } \pi(\lambda) = \sum_{n=0}^{\infty} a_n \lambda^n$$

$$\text{from } d\pi = 0 \quad \text{we get} \quad \sum_{n=0}^{\infty} n^2 a_n \lambda^n - \lambda \sum_{n=0}^{\infty} a_n (n+1/2)^2 \lambda^n$$

$$\Leftrightarrow \sum_{n=1}^{\infty} n^2 a_n \lambda^n - \sum_{n=0}^{\infty} a_n (n+1/2)^2 \lambda^{n+1}$$

$$\Leftrightarrow \sum_{n=1}^{\infty} (n^2 a_n - a_{n-1} (n-1/2)^2) \lambda^n = 0$$

$$\Leftrightarrow a_0 \text{ is free and } a_n = \frac{(n-1/2)^2}{n^2} a_{n-1} \quad (R)$$

Recall the Gamma function $\Gamma(z) = \int_0^{\infty} x^{z-1} e^{-x} dx$ (13)

$$\Gamma(n) = (n-1)!$$

it has $\text{Res}(\Gamma, -n) = \frac{(-1)^n}{n!}$, $n \in \mathbb{N}$, $\Psi(z) = \frac{\Gamma'(z)}{\Gamma(z)}$

The recurrence (R) gives (setting $a_0 = 1$)

$$\pi^0(\lambda) = \sum_{m=0}^{\infty} \frac{\Gamma(m+1/2)^2}{\Gamma(1/2)^2 \Gamma(m+1)^2} \lambda^m$$

What about the second solution?

Frobenius: for $\alpha = \alpha'$ a second linearly independent solution can be found from the first power series by

$$\pi^1(\lambda) = \frac{1}{2\pi i} \frac{d}{d\rho} \sum_{m=0}^{\infty} \frac{\Gamma(m+\rho+1/2)^2}{\Gamma(1/2)^2 \Gamma(m+\rho+1)^2} \lambda^{m+\rho} \Big|_{\rho=0}$$

$$(z^\rho = \exp(\rho \log \lambda))$$

$$\Leftrightarrow \pi^1(\lambda) = \frac{1}{2\pi i} \pi^0(\lambda) \log \lambda + \frac{1}{2\pi i} \sum_{m=0}^{\infty} \frac{\Gamma(m+1/2)^2}{\Gamma(1/2)^2 \Gamma(m+1)^2} \times (2\Psi(m+1/2) - 2\Psi(m+1))$$

$\log z!$ this is multi-valued \Rightarrow monodromy

$\lambda=0$

$$\lambda \rightarrow e^{2\pi i} \lambda, \quad \log \lambda \rightarrow \log e^{2\pi i} \lambda = \log \lambda + 2\pi i$$

$$\begin{pmatrix} \pi^0 \\ \pi^1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} \pi^0 \\ \pi^1 \end{pmatrix} =: M_0$$

We need to find M_1 and M_∞ , what we should do. (14)

i) Solve $\mathcal{L}\pi = 0$ around $\lambda=1$ and $\lambda=\infty$

ii) find the monodromy in the basis of solutions

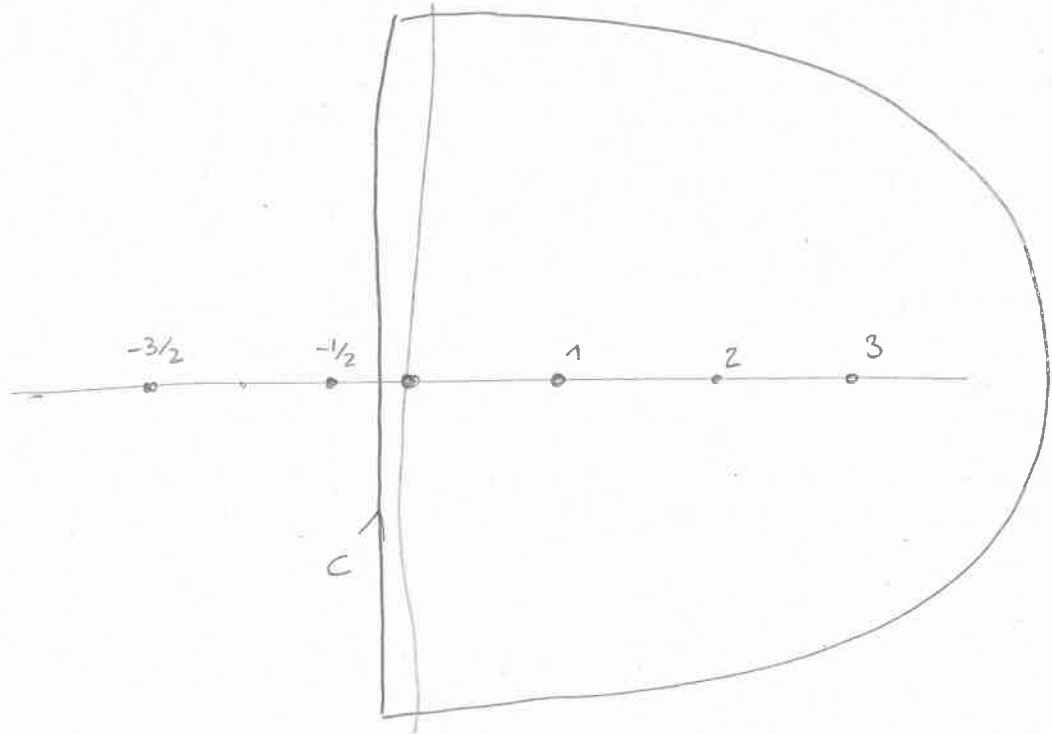
iii) Relate the different bases of solutions (analytic continuation)

For iii) it is useful to use analytic properties of the Γ function:

for example

$$\pi^0(\lambda) = \frac{1}{2\pi i \Gamma(1/2)} \int_C \frac{\Gamma(-s) \Gamma(s+1/2)^2}{\Gamma(s+1)} (-\lambda)^s ds$$

$|\arg(-\lambda)| < \pi$



1.8.3) Monodromy of ${}_2F_1$

(5)

We use a shortcut, namely the expressions for $\pi^\circ(\lambda)$ and $\pi'(\lambda)$ in terms of hypergeometric functions

recall, ${}_2F_1(a, b; c; z) := \mathcal{P} \left\{ \begin{matrix} 0 & \infty & 1 \\ 0 & a & 0 \\ 1-c & b & c-a-b \end{matrix} \right. z$

we get $\pi^\circ(\lambda) = {}_2F_1(\frac{1}{2}, \frac{1}{2}; 1; \lambda)$

$\pi'(\lambda) = \frac{-1}{2i} {}_2F_1(\frac{1}{2}, \frac{1}{2}; 1; 1-\lambda)$

with monodromy $M_0 = \begin{pmatrix} 1 & 0 \\ 1 & 1 \end{pmatrix}$

to find the monodromy around $A=1$ we write

$u = 1-\lambda$ and get

$\pi^\circ(u) = {}_2F_1(\frac{1}{2}, \frac{1}{2}; 1; 1-u)$

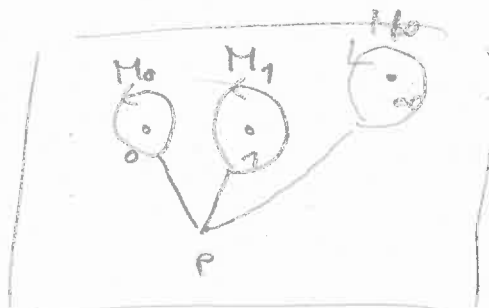
$\pi'(u) = \frac{-1}{2i} {}_2F_1(\frac{1}{2}, \frac{1}{2}; 1; u)$

and immediately obtain $M_1 = \begin{pmatrix} 1 & -4 \\ 0 & 1 \end{pmatrix}$

we know that

$M_\infty^{-1} \stackrel{!}{=} M_1 \cdot M_0$

$\Leftrightarrow M_\infty = \begin{pmatrix} -3 & -4 \\ 1 & 1 \end{pmatrix}$



$$M_0, M_1, M_\infty \in \Gamma_0(4) = \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \text{SL}(2, \mathbb{Z}) \mid b \equiv 0 \pmod{4} \right\} \quad (16)$$

the group is generated by M_0, M_1 .
